

# CLAS Excited Baryon Program



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M. Dugger, NSTAR, October 2022

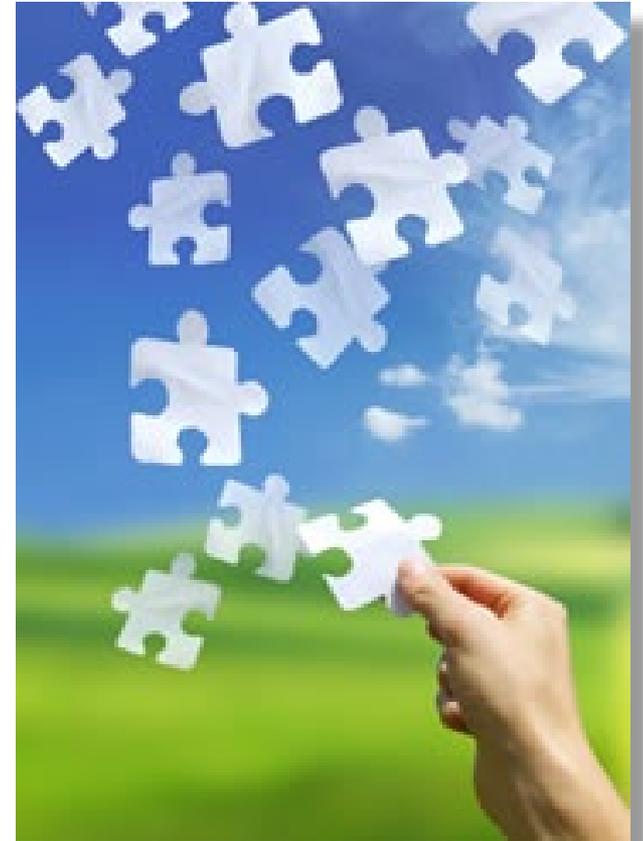


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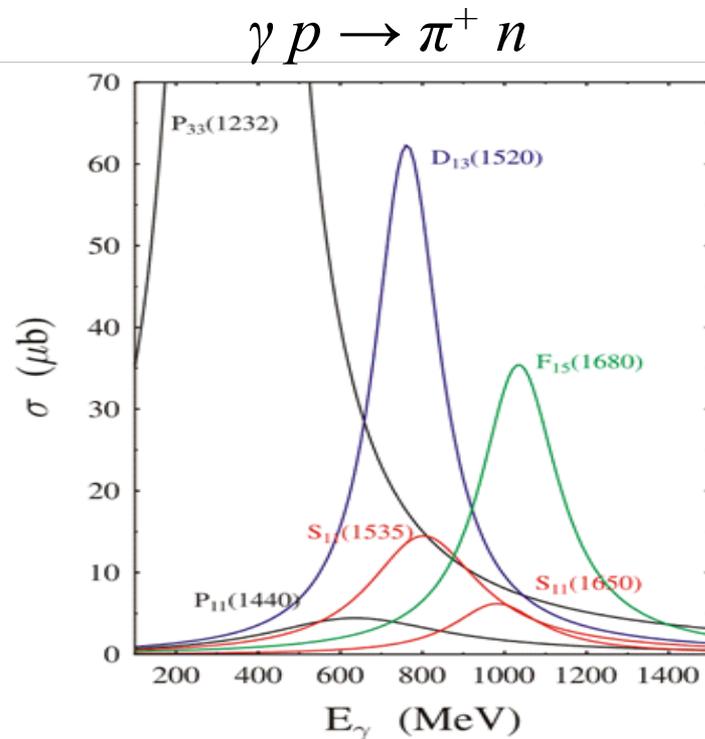
# Outline

- **Motivations**
- Helicity amplitudes
- Experimental facilities
- Reactions and results



# Nucleon resonances

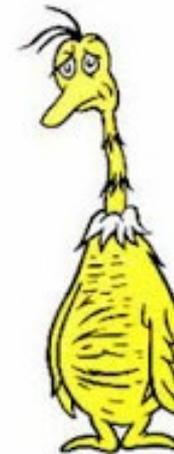
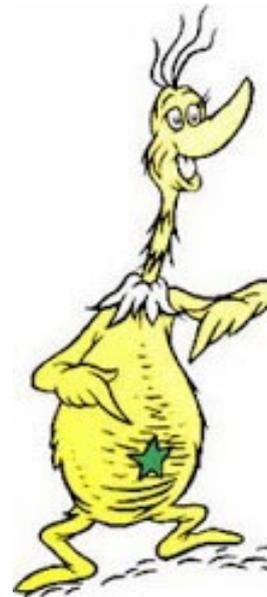
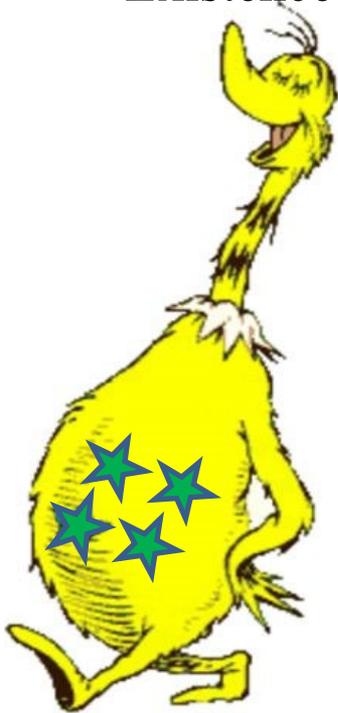
- As a three-quark system, the nucleon has a specific excitation spectrum comprised of nucleon resonances.
- This nucleon resonance spectrum has been found to have many broad overlapping states, making disentangling the spectrum difficult. ☹



# How well do we know the nucleon resonance spectrum?

Nucleon resonances are rated using the “star” system:

- \* Poor evidence of existence
- \*\* Fair evidence of existence
- \*\*\* Likely evidence of existence, or certain and properties need work
- \*\*\*\* Existence is certain and properties well explored



# Resonance status for $N^*$ and $\Delta^*$

Nucleon	Particle	$J^P$	Status as seen in																		
			overall	$N\gamma$	$N\pi$	$\Delta\pi$	$N\sigma$	$N\eta$	$\Delta K$	$\Sigma K$	$N\rho$	$N\omega$	$N\eta'$								
$N$	$N$	$1/2^+$	****																		
	$N(1440)$	$1/2^+$	****	****	****	****	***														
	$N(1520)$	$3/2^-$	****	****	****	****	**	****													
	$N(1535)$	$1/2^-$	****	****	****	***	*	****													
	$N(1650)$	$1/2^-$	****	****	****	***	*	****	*												
	$N(1675)$	$5/2^-$	****	****	****	****	***	*	*	*											
	$N(1680)$	$5/2^+$	****	****	****	****	***	*	*	*											
	$N(1700)$	$3/2^-$	***	**	***	***	*	*				*									
	$N(1710)$	$1/2^+$	****	****	****	*		***	**	*	*	*									
	$N(1720)$	$3/2^+$	****	****	****	***	*	*	****	*	*	*									
	$N(1860)$	$5/2^+$	**	*	**		*	*													
	$N(1875)$	$3/2^-$	***	**	**	*	**	*	*	*	*	*									
	$N(1880)$	$1/2^+$	***	**	*	**	*	*	**	**		**									
	$N(1895)$	$1/2^-$	****	****	*	*	*	****	**	**	*	*	****								
	$N(1900)$	$3/2^+$	****	****	**	**	*	*	**	**		*	**								
	$N(1990)$	$7/2^+$	**	**	**		*	*	*												
	$N(2000)$	$5/2^+$	**	**	*	**	*	*				*									
	$N(2040)$	$3/2^+$	*		*																
	$N(2060)$	$5/2^-$	***	***	**	*	*	*	*	*	*	*									
	$N(2100)$	$1/2^+$	***	**	***	**	**	*	*		*	*	**								
	$N(2120)$	$3/2^-$	***	***	**	**	**		**	*	*	*	*								
	$N(2190)$	$7/2^-$	****	****	****	****	**	*	**	*	*	*	*								
	$N(2220)$	$9/2^+$	****	**	****		*	*	*												
	$N(2250)$	$9/2^-$	****	**	****		*	*	*												
	$N(2300)$	$1/2^+$	**		**																
	$N(2570)$	$5/2^-$	**		**																
	$N(2600)$	$11/2^-$	***		***																
	$N(2700)$	$13/2^+$	**		**																

27  $N^*$  states:

- 13 with \*\*\*\*
- 7 with \*\*\*
- 6 with \*\*
- 1 with \*

Particle	$J^P$	Status as seen in						
		overall	$N\gamma$	$N\pi$	$\Delta\pi$	$\Sigma K$	$N\rho$	$\Delta\eta$
$\Delta(1232)$	$3/2^+$	****	****	****				
$\Delta(1600)$	$3/2^+$	****	****	***	****			
$\Delta(1620)$	$1/2^-$	****	****	****	****			
$\Delta(1700)$	$3/2^-$	****	****	****	****	*	*	
$\Delta(1750)$	$1/2^+$	*	*	*		*		
$\Delta(1900)$	$1/2^-$	***	***	***	*	**	*	
$\Delta(1905)$	$5/2^+$	****	****	****	**	*	*	**
$\Delta(1910)$	$1/2^+$	****	***	****	**	**		*
$\Delta(1920)$	$3/2^+$	***	***	***	***	**	**	**
$\Delta(1930)$	$5/2^-$	***	*	***	*	*		
$\Delta(1940)$	$3/2^-$	**	*	**	*			*
$\Delta(1950)$	$7/2^+$	****	****	****	**	***		
$\Delta(2000)$	$5/2^+$	**	*	**	*			*
$\Delta(2150)$	$1/2^-$	*		*				
$\Delta(2200)$	$7/2^-$	***	***	**	***	**		
$\Delta(2300)$	$9/2^+$	**		**				
$\Delta(2350)$	$5/2^-$	*		*				
$\Delta(2390)$	$7/2^+$	*		*				
$\Delta(2400)$	$9/2^-$	**	**	**				
$\Delta(2420)$	$11/2^+$	****	*	****				
$\Delta(2750)$	$13/2^-$	**		**				
$\Delta(2950)$	$15/2^+$	**		**				

22  $\Delta^*$  states:

- 8 with \*\*\*\*
- 4 with \*\*\*
- 6 with \*\*
- 4 with \*

# Resonance status for $N^*$ and $\Delta^*$

Nucleon $\rightarrow$	Particle	$J^P$	Status as seen in																		
			overall	$N\gamma$	$N\pi$	$\Delta\pi$	$N\sigma$	$N\eta$	$\Delta K$	$\Sigma K$	$N\rho$	$N\omega$	$N\eta'$								
$N$	$N$	$1/2^+$	****																		
	$N(1440)$	$1/2^+$	****	****	****	****	***														
	$N(1520)$	$3/2^-$	****	****	****	****	**	****													
	$N(1535)$	$1/2^-$	****	****	****	***	*	****													
	$N(1650)$	$1/2^-$	****	****	****	***	*	****	*												
	$N(1675)$	$5/2^-$	****	****	****	****	***	*	*	*											
	$N(1680)$	$5/2^+$	****	****	****	****	***	*	*	*											
	$N(1700)$	$3/2^-$	***	**	***	***	*	*			*										
	$N(1710)$	$1/2^+$	****	****	****	*		***	**	*	*	*									
	$N(1720)$	$3/2^+$	****	****	****	***	*	*	****	*	*	*									
	$N(1860)$	$5/2^+$	**	*	**		*	*													
	$N(1875)$	$3/2^-$	***	**	**	*	**	*	*	*	*	*									
	$N(1880)$	$1/2^+$	***	**	*	**	*	*	**	**		**									
	$N(1895)$	$1/2^-$	****	****	*	*	*	****	**	**	*	*	****								
	$N(1900)$	$3/2^+$	****	****	**	**	*	*	**	**		*	**								
	$N(1990)$	$7/2^+$	**	**	**		*	*	*												
	$N(2000)$	$5/2^+$	**	**	*	**	*	*			*										
	$N(2040)$	$3/2^+$	*		*																
	$N(2060)$	$5/2^-$	***	***	**	*	*	*	*	*	*	*									
	$N(2100)$	$1/2^+$	***	**	***	**	**	*	*		*	*	**								
	$N(2120)$	$3/2^-$	***	***	**	**	**		**	*	*	*	*								
	$N(2190)$	$7/2^-$	****	****	****	****	**	*	**	*	*	*	*								
	$N(2220)$	$9/2^+$	****	**	****		*	*	*												
	$N(2250)$	$9/2^-$	****	**	****		*	*	*												
	$N(2300)$	$1/2^+$	**		**																
	$N(2570)$	$5/2^-$	**		**																
	$N(2600)$	$11/2^-$	***		***																
	$N(2700)$	$13/2^+$	**		**																

- 27  $N^*$  states:
- 13 with \*\*\*\*
  - 7 with \*\*\*
  - 6 with \*\*
  - 1 with \*

In 2013

- 26  $N^*$  states:
- 10 with \*\*\*\*
  - 5 with \*\*\*
  - 8 with \*\*
  - 3 with \*

Particle	$J^P$	Status as seen in						
		overall	$N\gamma$	$N\pi$	$\Delta\pi$	$\Sigma K$	$N\rho$	$\Delta\eta$
$\Delta(1232)$	$3/2^+$	****	****	****				
$\Delta(1600)$	$3/2^+$	****	****	***	****			
$\Delta(1620)$	$1/2^-$	****	****	****	****			
$\Delta(1700)$	$3/2^-$	****	****	****	****	*	*	
$\Delta(1750)$	$1/2^+$	*	*	*		*		
$\Delta(1900)$	$1/2^-$	***	***	***	*	**	*	
$\Delta(1905)$	$5/2^+$	****	****	****	**	*	*	**
$\Delta(1910)$	$1/2^+$	****	***	****	**	**		*
$\Delta(1920)$	$3/2^+$	***	***	***	***	**	**	**
$\Delta(1930)$	$5/2^-$	***	*	***	*	*		
$\Delta(1940)$	$3/2^-$	**	*	**	*			*
$\Delta(1950)$	$7/2^+$	****	****	****	**	***		
$\Delta(2000)$	$5/2^+$	**	*	**	*		*	
$\Delta(2150)$	$1/2^-$	*		*				
$\Delta(2200)$	$7/2^-$	***	***	**	***	**		
$\Delta(2300)$	$9/2^+$	**		**				
$\Delta(2350)$	$5/2^-$	*		*				
$\Delta(2390)$	$7/2^+$	*		*				
$\Delta(2400)$	$9/2^-$	**	**	**				
$\Delta(2420)$	$11/2^+$	****	*	****				
$\Delta(2750)$	$13/2^-$	**		**				
$\Delta(2950)$	$15/2^+$	**		**				

- 22  $\Delta^*$  states:
- 8 with \*\*\*\*
  - 4 with \*\*\*
  - 6 with \*\*
  - 4 with \*

In 2013

- 22  $\Delta^*$  states:
- 7 with \*\*\*\*
  - 3 with \*\*\*
  - 7 with \*\*
  - 5 with \*



# Resonance status for $\Xi^*$

State, $J^P$	Predicted masses (MeV)							
$\Xi \frac{1}{2}^+$	1305							
$\Xi \frac{3}{2}^+$	1505							
$\Xi^* \frac{1}{2}^-$	1755	1810	1835	2225	2285	2300	2320	2380
$\Xi^* \frac{3}{2}^-$	1785	1880	1895	2240	2305	2330	2340	2385
$\Xi^* \frac{5}{2}^-$	1900	2345	2350	2385				
$\Xi^* \frac{7}{2}^-$	2355							
$\Xi^* \frac{1}{2}^+$	1840	2040	2100	2130	2150	2230	2345	
$\Xi^* \frac{3}{2}^+$	2045	2065	2115	2165	2170	2210	2230	2275
$\Xi^* \frac{5}{2}^+$	2045	2165	2230	2230	2240			
$\Xi^* \frac{7}{2}^+$	2180	2240						

+

- List of Cascade Baryons predicted by Capstick and Isgur with mass less than  $2.4 \text{ GeV}/c^2$

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$\Xi \frac{1}{2}^+$	1305							
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$\Xi^* \frac{1}{2}^-$	1755	1810	1835	2225	2285	2300	2320	2380
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$\Xi^* \frac{7}{2}^-$	2355							
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$\Xi^* \frac{5}{2}^+$	2045	2165	2230	2230	2240			
$\Xi^* \frac{7}{2}^+$	2180	2240						

PDG		Overall
Particle	$J^P$	Status
$\Xi(1318)$	$1/2^+$	****
$\Xi(1530)$	$3/2^+$	****
$\Xi(1620)$		*
$\Xi(1690)$		***
$\Xi(1820)$	$3/2^-$	***
$\Xi(1950)$		***
$\Xi(2030)$	$5/2^?$	***
$\Xi(2120)$		*
$\Xi(2250)$		**
$\Xi(2370)$		**
$\Xi(2500)$		*

- List of Cascade Baryons predicted by Capstick and Isgur with mass less than  $2.4 \text{ GeV}/c^2$

# Resonance status for $\Xi^*$

State, $J^P$	Predicted masses (MeV)							
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$\Xi^* \frac{3}{2}^-$	1785	1880	1895	2240	2305	2330	2340	2385
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$\Xi^* \frac{7}{2}^-$	2355							
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$\Xi^* \frac{3}{2}^+$	2045	2065	2115	2165	2170	2210	2230	2275
$\Xi^* \frac{5}{2}^+$	2045	2165	2230	2230	2240			
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Particle	$J^P$	Status
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$\Xi(1950)$		***
$\Xi(2030)$	$5/2^?$	***
$\Xi(2120)$		*
$\Xi(2250)$		**
$\Xi(2370)$		**
$\Xi(2500)$		*

- List of Cascade Baryons predicted by Capstick and Isgur with mass less than  $2.4 \text{ GeV}/c^2$

State	$\Lambda K$	$\Sigma K$	$\Xi\pi$
$\Xi(1530)$			100 %
$\Xi(1690)$	seen	seen	seen
$\Xi(1820)$	large	small	small
$\Xi(1950)$	seen	seen?	seen
$\Xi(2030)$	20%	80%	small



# Resonance status for $\Xi^*$

State, $J^P$	Predicted masses (MeV)							
$\Xi \frac{1}{2}^+$	1305							
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$\Xi^* \frac{1}{2}^-$	1755	1810	1835	2225	2285	2300	2320	2380
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- List of Cascade Baryons predicted by Capstick and Isgur with mass less than  $2.4 \text{ GeV}/c^2$

PDG		
Particle	$J^P$	Overall Status
$\Xi(1318)$	$1/2^+$	****
$\Xi(1530)$	$3/2^+$	****
$\Xi(1620)$		*
$\Xi(1690)$		***
$\Xi(1820)$	$3/2^-$	***
$\Xi(1950)$		***
$\Xi(2030)$	$5/2^?$	***
$\Xi(2120)$		*
$\Xi(2250)$		**
$\Xi(2370)$		**
$\Xi(2500)$		*

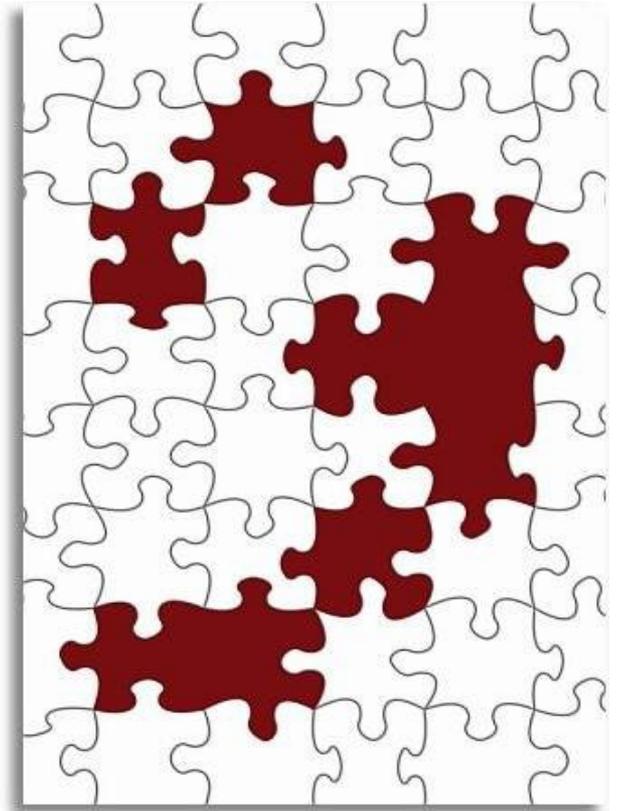
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$\Xi(1690)$	seen	seen	seen
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$\Xi(1950)$	seen	seen?	seen
$\Xi(2030)$	20%	80%	small



- Are there missing states?

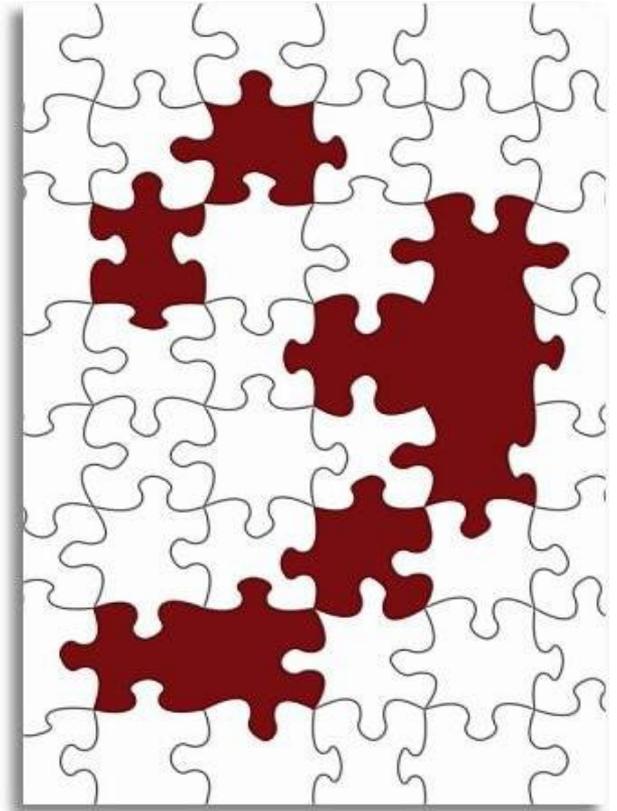
# So, where are the resonances?

- Masses, widths, and coupling constants not well known for many resonances



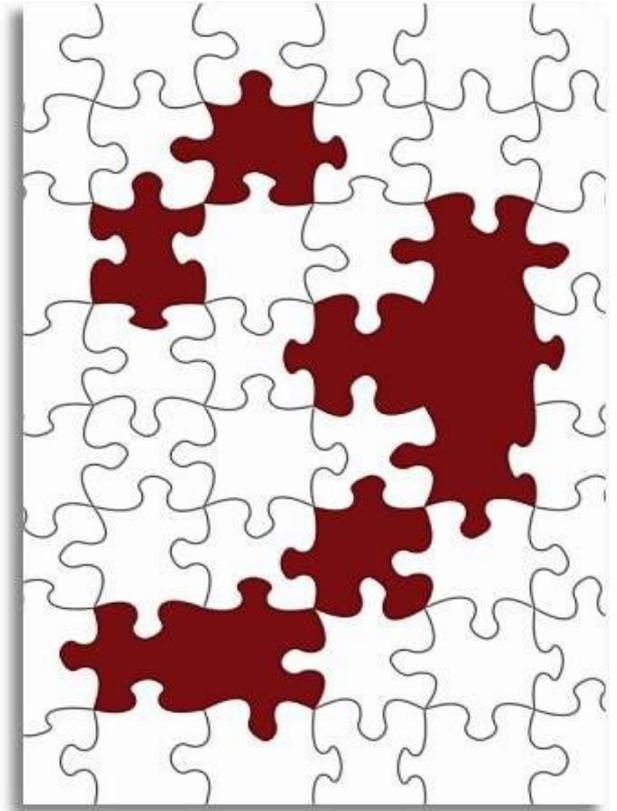
# So, where are the resonances?

- Masses, widths, and coupling constants not well known for many resonances
- Many models exist to “predict” the nucleon resonance spectrum - quark model, Goldstone-boson exchange, diquark and collective models, instanton-induced interactions, flux-tube models, lattice QCD - **BUT...**



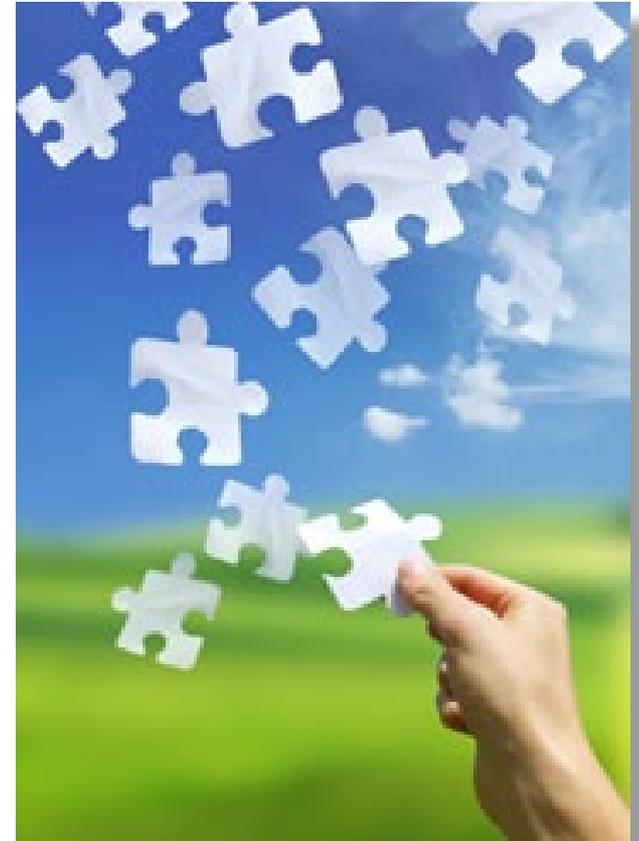
# So, where are the resonances?

- Masses, widths, and coupling constants not well known for many resonances
- Many models exist to “predict” the nucleon resonance spectrum - quark model, Goldstone-boson exchange, diquark and collective models, instanton-induced interactions, flux-tube models, lattice QCD - **BUT...**
- **THE BIG PUZZLE: Most models predict many more resonance states than have been observed.**



# Outline

- Motivations
- **Helicity amplitudes**
- Experimental facilities
- Reactions and results



# Helicity amplitudes for $\gamma + p \rightarrow p + \text{pseudoscalar}$

- 8 helicity states: 4 initial, 2 final  $\rightarrow 4 \cdot 2 = 8$
- Amplitudes are complex but parity symmetry reduces independent numbers to 8
- Overall phase unobservable  $\rightarrow 7$  independent numbers
- **HOWEVER**, not all possible observables are linearly independent and it turns out that there must be a minimum of 8 observables / experiments

$$A = \begin{array}{c} \text{Initial helicity} \\ \left[ \begin{array}{cc} A_{11} & A_{12} \\ A_{21} & A_{22} \end{array} \right] \text{Final helicity} \end{array}$$

helicity +1 photons ( $\varepsilon_+$ ):

$$A_{\varepsilon_+} = \frac{1}{2} \begin{bmatrix} \frac{3}{2} & \frac{1}{2} \\ H_1 & H_2 \\ -\frac{1}{2} & H_3 \\ & H_4 \end{bmatrix}$$

$$(A_{-\mu, -\lambda} = -e^{(\lambda-\mu)\pi} A_{\mu, \lambda})$$

Parity symmetry  $\rightarrow$

helicity -1 photons ( $\varepsilon_-$ ):

$$A_{\varepsilon_-} = \frac{1}{2} \begin{bmatrix} \frac{-1}{2} & \frac{-3}{2} \\ H_4 & -H_3 \\ -\frac{1}{2} & H_2 \\ & H_1 \end{bmatrix}$$

# Linkage between helicity amplitudes and the observables for single pseudoscalar photoproduction

Spin observable	Helicity representation
$\check{\Omega}^1 \equiv \mathcal{I}(\theta)$	$\frac{1}{2}( H_1 ^2 +  H_2 ^2 +  H_3 ^2 +  H_4 ^2)$ ← Differential cross section
$\check{\Omega}^4 \equiv \check{\Sigma}$	$\text{Re}(-H_1H_4^* + H_2H_3^*)$
$\check{\Omega}^{10} \equiv -\check{T}$	$\text{Im}(H_1H_2^* + H_3H_4^*)$
$\check{\Omega}^{12} \equiv \check{P}$	$\text{Im}(-H_1H_3^* - H_2H_4^*)$
$\check{\Omega}^3 \equiv \check{G}$	$\text{Im}(H_1H_4^* - H_3H_2^*)$
$\check{\Omega}^5 \equiv \check{H}$	$\text{Im}(-H_2H_4^* + H_1H_3^*)$
$\check{\Omega}^9 \equiv \check{E}$	$\frac{1}{2}( H_1 ^2 -  H_2 ^2 +  H_3 ^2 -  H_4 ^2)$
$\check{\Omega}^{11} \equiv \check{F}$	$\text{Re}(-H_2H_1^* - H_4H_3^*)$
$\check{\Omega}^{14} \equiv \check{O}_x$	$\text{Im}(-H_2H_1^* + H_4H_3^*)$
$\check{\Omega}^7 \equiv -\check{O}_z$	$\text{Im}(H_1H_4^* - H_2H_3^*)$
$\check{\Omega}^{16} \equiv -\check{C}_x$	$\text{Re}(H_2H_4^* + H_1H_3^*)$
$\check{\Omega}^2 \equiv -\check{C}_z$	$\frac{1}{2}( H_1 ^2 +  H_2 ^2 -  H_3 ^2 -  H_4 ^2)$
$\check{\Omega}^6 \equiv -\check{T}_x$	$\text{Re}(-H_1H_4^* - H_2H_3^*)$
$\check{\Omega}^{13} \equiv -\check{T}_z$	$\text{Re}(-H_1H_2^* + H_4H_3^*)$
$\check{\Omega}^8 \equiv \check{L}_x$	$\text{Re}(H_2H_4^* - H_1H_3^*)$
$\check{\Omega}^{15} \equiv \check{L}_z$	$\frac{1}{2}(- H_1 ^2 +  H_2 ^2 +  H_3 ^2 -  H_4 ^2)$

# Linkage between helicity amplitudes and the observables for single pseudoscalar photoproduction

Spin observable	Helicity representation	
$\check{\Omega}^1 \equiv \mathcal{I}(\theta)$	$\frac{1}{2}( H_1 ^2 +  H_2 ^2 +  H_3 ^2 +  H_4 ^2)$	Differential cross section
$\check{\Omega}^4 \equiv \check{\Sigma}$	$\text{Re}(-H_1H_4^* + H_2H_3^*)$	Beam polarization $\Sigma$
$\check{\Omega}^{10} \equiv -\check{T}$	$\text{Im}(H_1H_2^* + H_3H_4^*)$	
$\check{\Omega}^{12} \equiv \check{P}$	$\text{Im}(-H_1H_3^* - H_2H_4^*)$	
$\check{\Omega}^3 \equiv \check{G}$	$\text{Im}(H_1H_4^* - H_3H_2^*)$	
$\check{\Omega}^5 \equiv \check{H}$	$\text{Im}(-H_2H_4^* + H_1H_3^*)$	
$\check{\Omega}^9 \equiv \check{E}$	$\frac{1}{2}( H_1 ^2 -  H_2 ^2 +  H_3 ^2 -  H_4 ^2)$	
$\check{\Omega}^{11} \equiv \check{F}$	$\text{Re}(-H_2H_1^* - H_4H_3^*)$	
$\check{\Omega}^{14} \equiv \check{O}_x$	$\text{Im}(-H_2H_1^* + H_4H_3^*)$	
$\check{\Omega}^7 \equiv -\check{O}_z$	$\text{Im}(H_1H_4^* - H_2H_3^*)$	
$\check{\Omega}^{16} \equiv -\check{C}_x$	$\text{Re}(H_2H_4^* + H_1H_3^*)$	
$\check{\Omega}^2 \equiv -\check{C}_z$	$\frac{1}{2}( H_1 ^2 +  H_2 ^2 -  H_3 ^2 -  H_4 ^2)$	
$\check{\Omega}^6 \equiv -\check{T}_x$	$\text{Re}(-H_1H_4^* - H_2H_3^*)$	
$\check{\Omega}^{13} \equiv -\check{T}_z$	$\text{Re}(-H_1H_2^* + H_4H_3^*)$	
$\check{\Omega}^8 \equiv \check{L}_x$	$\text{Re}(H_2H_4^* - H_1H_3^*)$	
$\check{\Omega}^{15} \equiv \check{L}_z$	$\frac{1}{2}(- H_1 ^2 +  H_2 ^2 +  H_3 ^2 -  H_4 ^2)$	

# Linkage between helicity amplitudes and the observables for single pseudoscalar photoproduction

Spin observable	Helicity representation	
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$\check{\Omega}^{10} \equiv -\check{T}$	$\text{Im}(H_1H_2^* + H_3H_4^*)$	Target asymmetry $T$
$\check{\Omega}^{12} \equiv \check{P}$	$\text{Im}(-H_1H_3^* - H_2H_4^*)$	
$\check{\Omega}^3 \equiv \check{G}$	$\text{Im}(H_1H_4^* - H_3H_2^*)$	
$\check{\Omega}^5 \equiv \check{H}$	$\text{Im}(-H_2H_4^* + H_1H_3^*)$	
$\check{\Omega}^9 \equiv \check{E}$	$\frac{1}{2}( H_1 ^2 -  H_2 ^2 +  H_3 ^2 -  H_4 ^2)$	
$\check{\Omega}^{11} \equiv \check{F}$	$\text{Re}(-H_2H_1^* - H_4H_3^*)$	
$\check{\Omega}^{14} \equiv \check{O}_x$	$\text{Im}(-H_2H_1^* + H_4H_3^*)$	
$\check{\Omega}^7 \equiv -\check{O}_z$	$\text{Im}(H_1H_4^* - H_2H_3^*)$	
$\check{\Omega}^{16} \equiv -\check{C}_x$	$\text{Re}(H_2H_4^* + H_1H_3^*)$	
$\check{\Omega}^2 \equiv -\check{C}_z$	$\frac{1}{2}( H_1 ^2 +  H_2 ^2 -  H_3 ^2 -  H_4 ^2)$	
$\check{\Omega}^6 \equiv -\check{T}_x$	$\text{Re}(-H_1H_4^* - H_2H_3^*)$	
$\check{\Omega}^{13} \equiv -\check{T}_z$	$\text{Re}(-H_1H_2^* + H_4H_3^*)$	
$\check{\Omega}^8 \equiv \check{L}_x$	$\text{Re}(H_2H_4^* - H_1H_3^*)$	
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$\check{\Omega}^3 \equiv \check{G}$	$\text{Im}(H_1H_4^* - H_3H_2^*)$	Recoil polarization $P$
$\check{\Omega}^5 \equiv \check{H}$	$\text{Im}(-H_2H_4^* + H_1H_3^*)$	
$\check{\Omega}^9 \equiv \check{E}$	$\frac{1}{2}( H_1 ^2 -  H_2 ^2 +  H_3 ^2 -  H_4 ^2)$	
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$\check{\Omega}^{16} \equiv -\check{C}_x$	$\text{Re}(H_2H_4^* + H_1H_3^*)$	
$\check{\Omega}^2 \equiv -\check{C}_z$	$\frac{1}{2}( H_1 ^2 +  H_2 ^2 -  H_3 ^2 -  H_4 ^2)$	
$\check{\Omega}^6 \equiv -\check{T}_x$	$\text{Re}(-H_1H_4^* - H_2H_3^*)$	
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# Linkage between helicity amplitudes and the observables for single pseudoscalar photoproduction

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$\check{\Omega}^3 \equiv \check{G}$	$\text{Im}(H_1H_4^* - H_3H_2^*)$	Double polarization observables
$\check{\Omega}^5 \equiv \check{H}$	$\text{Im}(-H_2H_4^* + H_1H_3^*)$	
$\check{\Omega}^9 \equiv \check{E}$	$\frac{1}{2}( H_1 ^2 -  H_2 ^2 +  H_3 ^2 -  H_4 ^2)$	
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$\check{\Omega}^{16} \equiv -\check{C}_x$	$\text{Re}(H_2H_4^* + H_1H_3^*)$	
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$\check{\Omega}^6 \equiv -\check{T}_x$	$\text{Re}(-H_1H_4^* - H_2H_3^*)$	
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Transverse target  
+  
Longitudinal target

Polarized photons

# Linkage between helicity amplitudes and the observables for single pseudoscalar photoproduction

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Transverse target  
+  
Longitudinal target  
+  
Polarized photons

Differential cross section

Beam polarization  $\Sigma$

Target asymmetry  $T$

Recoil polarization  $P$

**Double polarization observables**

- Need **at least** 4 of the double observables from at least 2 groups for a “complete experiment”

# Linkage between helicity amplitudes and the observables for single pseudoscalar photoproduction

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**Double polarization observables**

• Need **at least** 4 of the double observables from at least 2 groups for a “complete experiment”

•  $\pi^0 p$ ,  $\pi^+ n$ , and  $\eta p$  will be nearly complete

Transverse target  
+  
Longitudinal target

Polarized photons

# Linkage between helicity amplitudes and the observables for single pseudoscalar photoproduction

Spin observable	Helicity representation
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**Double polarization observables**

• Need **at least** 4 of the double observables from at least 2 groups for a “complete experiment”

•  $\pi^0 p$ ,  $\pi^+ n$ , and  $\eta p$  will be nearly complete

•  $K^+ \Lambda$  will be complete!

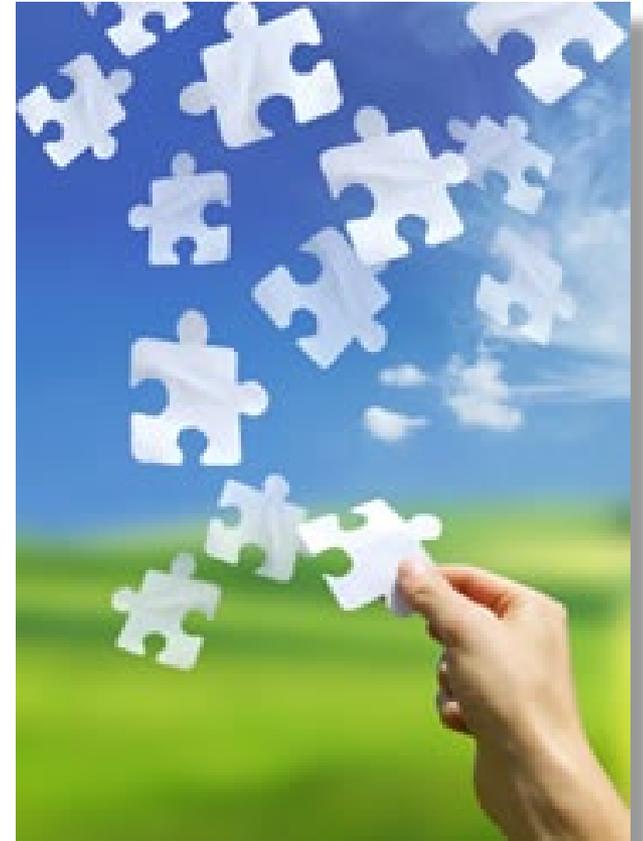
**So, finding missing resonances  
requires lots of different  
observables.**

**Cross sections are not enough!**



# Outline

- Motivations
- Helicity amplitudes
- **Experimental facilities**
- Reactions and results

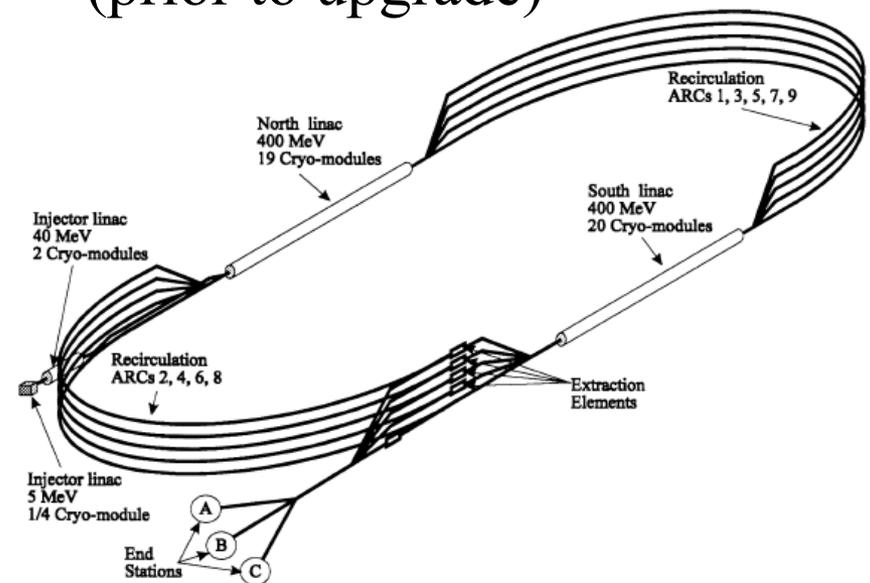


# Experimental facilities:

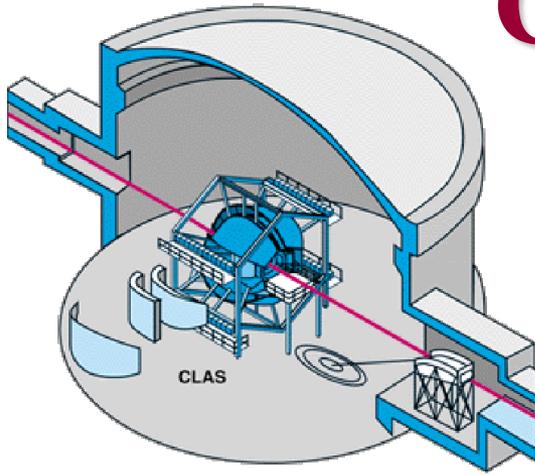
- The Thomas Jefferson National Accelerator Facility (Jefferson Laboratory = JLab).
- Continuous Electron Beam Accelerator Facility (CEBAF)



- Racetrack design
- Energies up to 6 GeV (prior to upgrade)

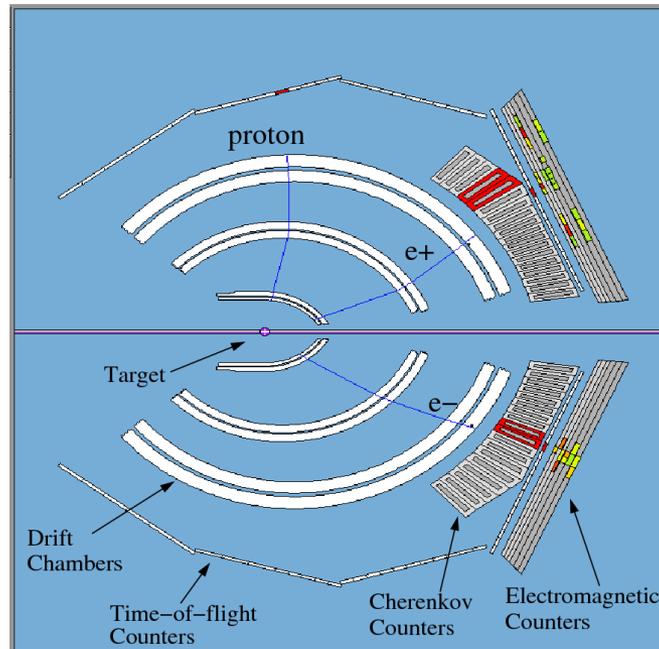


# CLAS (1997-2012)

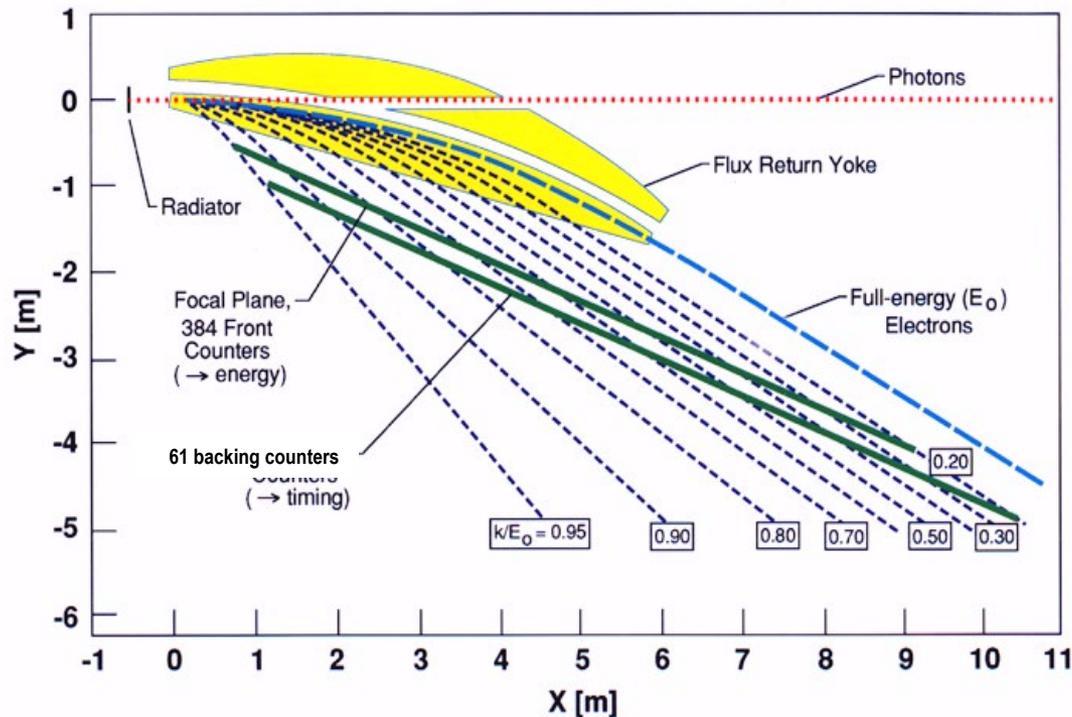


Lest we forget:

- CLAS was very good for detecting charged particles
- CLAS had a rather large acceptance

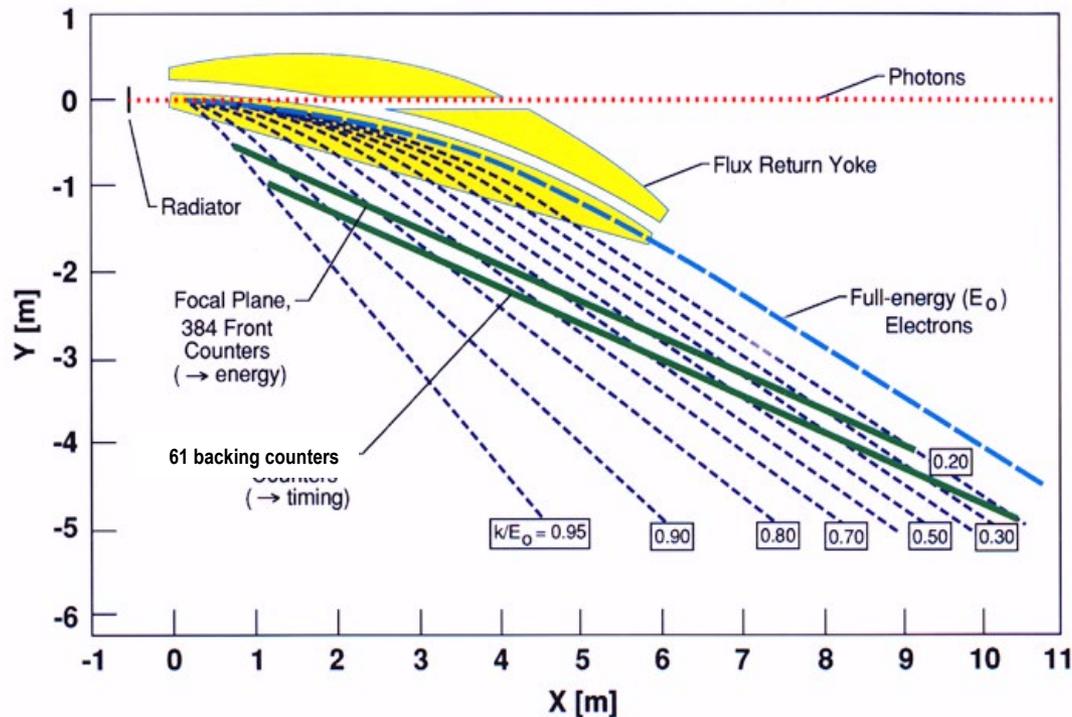


# Bremsstrahlung photon tagger (also deceased)



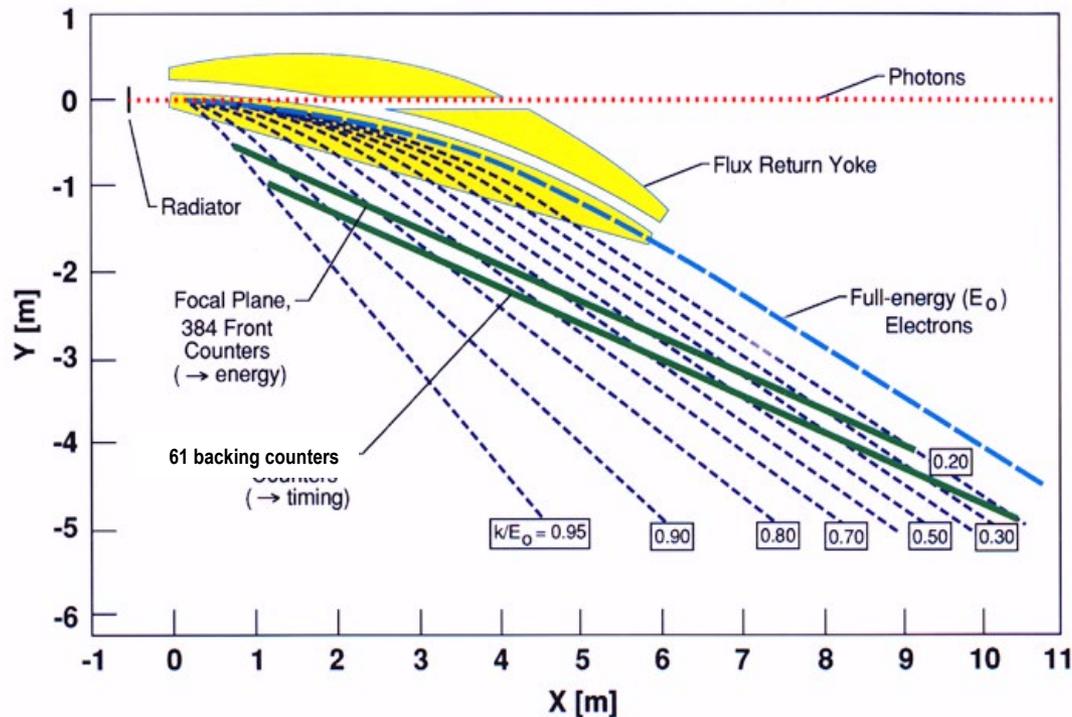
- Jefferson Lab Hall B bremsstrahlung photon tagger had:
  - $E_\gamma = 20\text{-}95\%$  of  $E_0$
  - $E_\gamma$  up to  $\sim 5.5$  GeV

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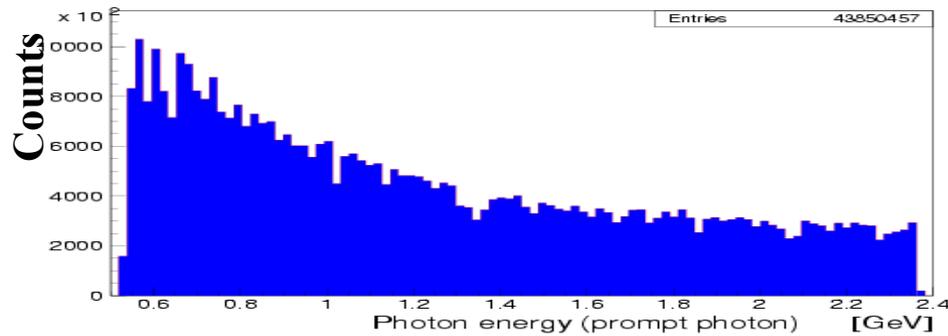
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  - $E_\gamma = 20\text{-}95\%$  of  $E_0$
  - $E_\gamma$  up to  $\sim 5.5$  GeV
  - **Circular polarized photons with longitudinally polarized electrons**

# Bremsstrahlung photon tagger (also deceased)

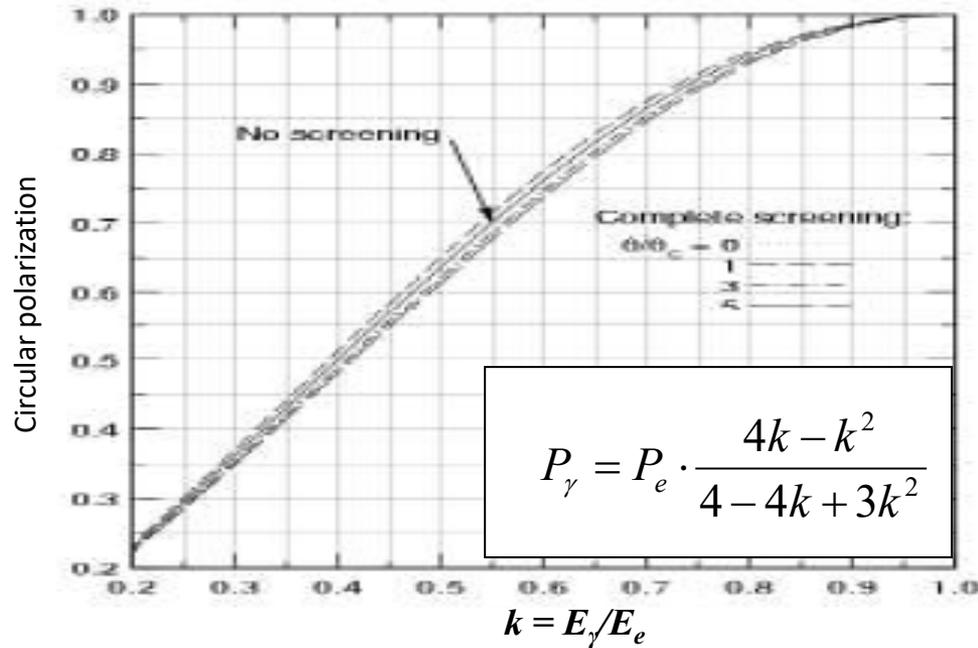


- Jefferson Lab Hall B bremsstrahlung photon tagger had:
  - $E_\gamma = 20\text{-}95\%$  of  $E_0$
  - $E_\gamma$  up to  $\sim 5.5$  GeV
  - **Circular polarized photons with longitudinally polarized electrons**
  - **Oriented diamond crystal for linearly polarized photons**

# Circular beam polarization



Circular polarization from 100% polarized electron beam



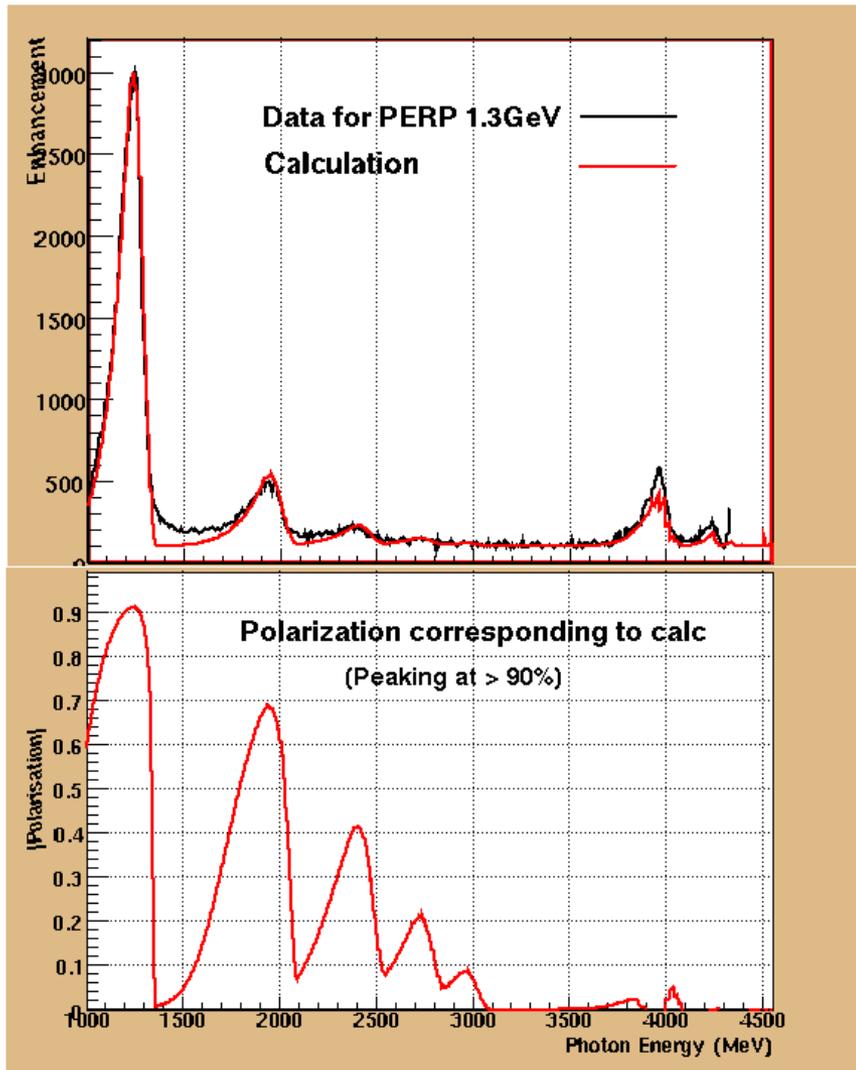
- Circular photon beam from longitudinally-polarized electrons

- Incident electron beam polarization > 85%



H. Olsen and L.C. Maximon, Phys. Rev. 114, 887 (1959)

# Linearly polarized photons



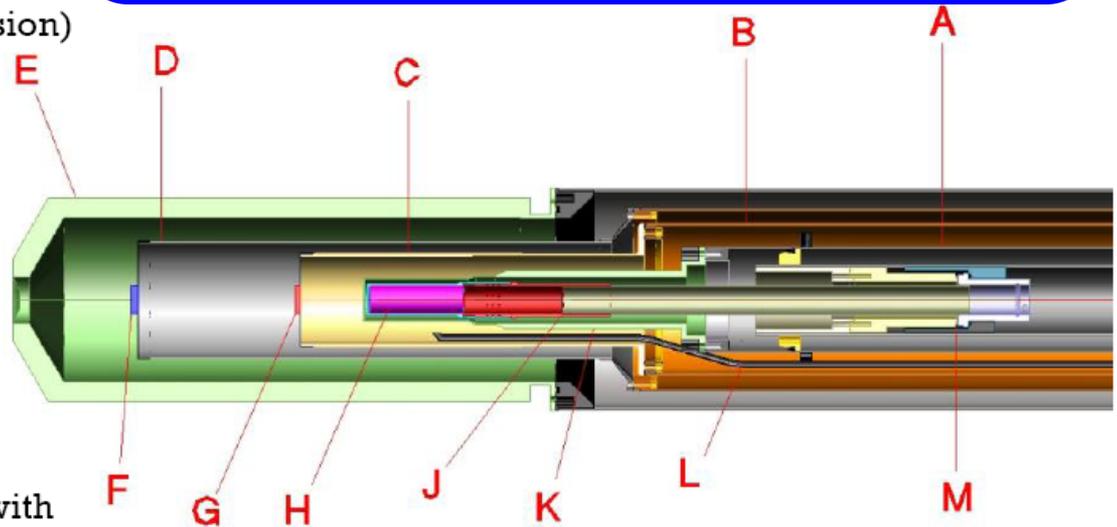
- Coherent bremsstrahlung from 50- $\mu$  oriented diamond
- Two linear polarization states (vertical & horizontal)
- Analytical QED coherent bremsstrahlung calculation fit to actual spectrum (Livingston/Glasgow)
- Vertical 1.3 GeV edge shown

# FROST target

- Butanol composition:  $C_4H_9OH$
- C and O are even-even nuclei → No polarization of the bound nucleons

## The FroST target and its components:

- A: Primary heat exchanger
- B: 1 K heat shield
- C: Holding coil
- D: 20 K heat shield
- E: Outer vacuum can (Rohacell extension)
- F:  $CH_2$  target
- G: Carbon target
- H: Butanol target
- J: Target insert
- K: Mixing chamber
- L: Microwave waveguide
- M: Kapton coldseal



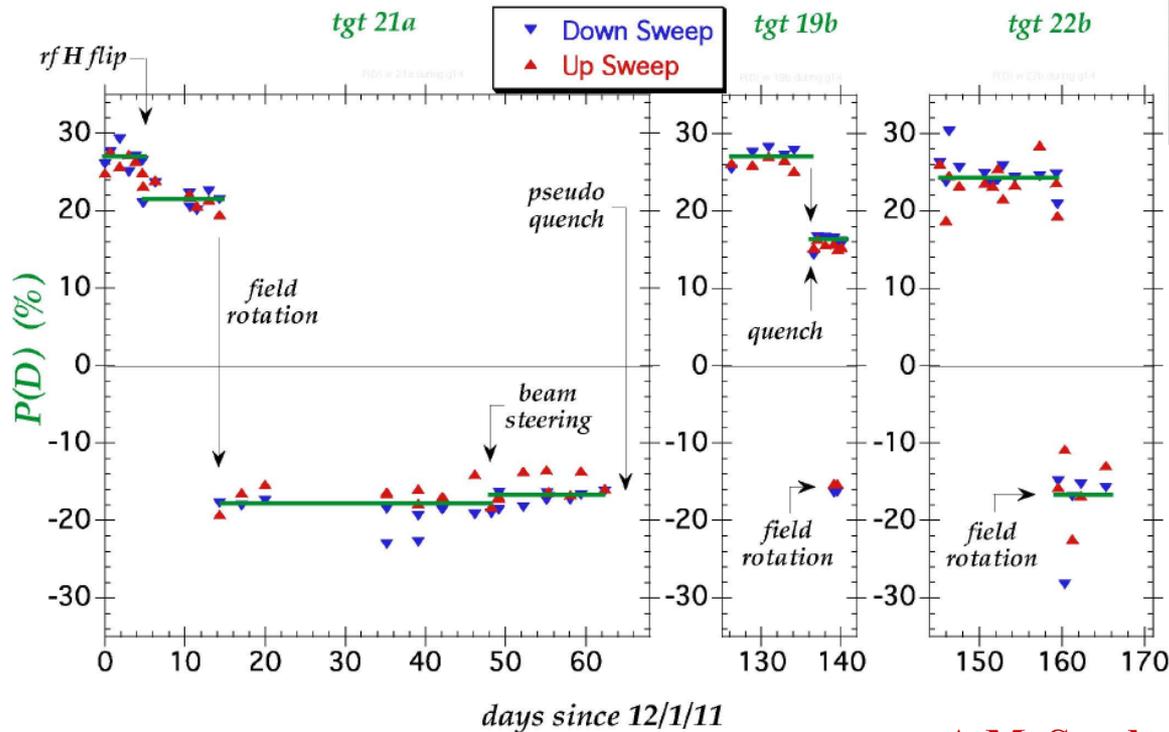
## Performance Specs:

- Base Temp: 28 mK w/o beam, 30 mK with
- Cooling Power: 800  $\mu W$  @ 50 mK, 10 mW @ 100 mK, and 60 mW @ 300 mK
- Polarization: +82%, -90%
- 1/e Relaxation Time: 2800 hours (+Pol), 1600 hours (-Pol)
- Roughly 1% polarization loss per day.

- Carbon target used to represent bound nucleon contribution of butanol

# HD-ICE target

## D polarization during g14/E06-101



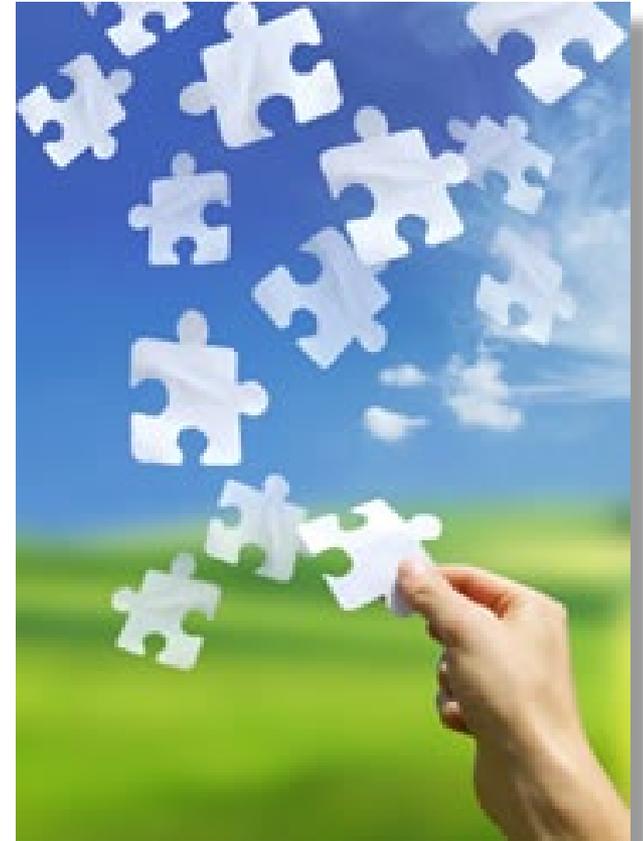
- Deuteron target

A.M. Sandorfi



# Outline

- Motivations
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- **Reactions and results**



# Pion photoproduction

# Isospin combinations for reactions involving $\pi^0$ and $\pi^+$

- Differing isospin compositions for  $N^*$  and  $\Delta^+$  for the  $\pi^0 p$  and  $\pi^+ n$  final states
- The  $\pi^0 p$  and  $\pi^+ n$  final states can help distinguish between the  $\Delta$  and  $N^*$

$$\begin{array}{ccc} \Delta^+ & & N^* \\ \downarrow & & \downarrow \\ \pi^0 + p : \sqrt{2/3} \left| I = \frac{3}{2}, I_3 = \frac{1}{2} \right\rangle - \sqrt{1/3} \left| I = \frac{1}{2}, I_3 = \frac{1}{2} \right\rangle & & \\ \pi^+ + n : \sqrt{1/3} \left| I = \frac{3}{2}, I_3 = \frac{1}{2} \right\rangle + \sqrt{2/3} \left| I = \frac{1}{2}, I_3 = \frac{1}{2} \right\rangle & & \end{array}$$

# Isospin photo-couplings

- Using both proton and neutron targets allows decomposition of iso-singlet and iso-vector photo-couplings  $C^0$ ,  $C^1$

Example:

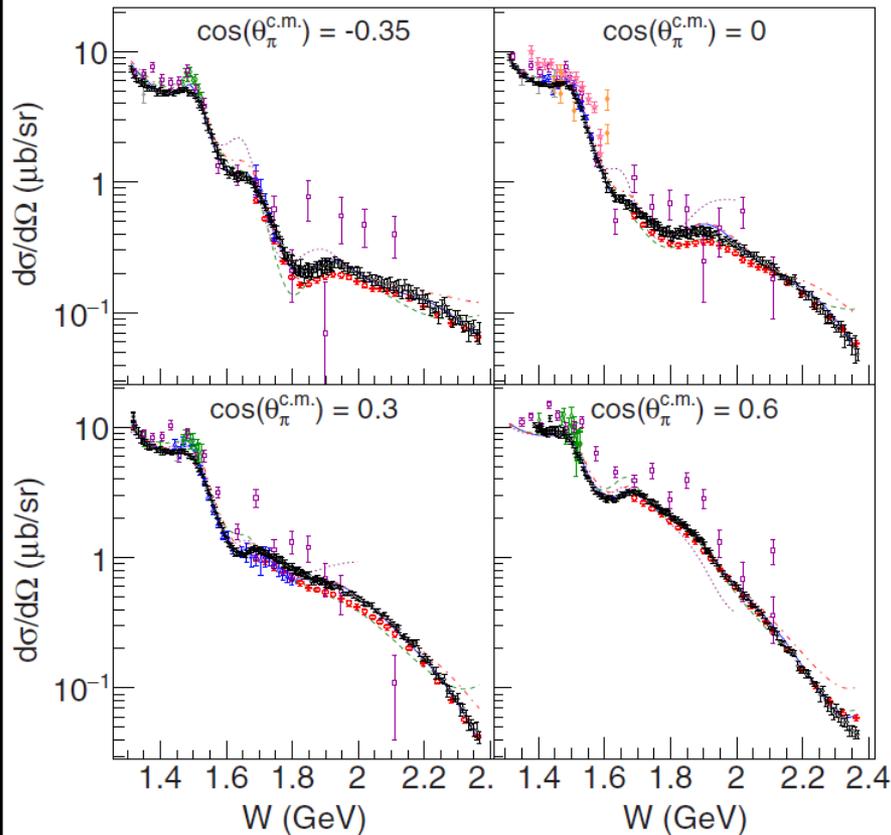
$$\gamma p \rightarrow n \pi^+ : \quad \pm \sqrt{\frac{2}{3}} \left[ C^0 \ominus \sqrt{\frac{1}{3}} C^1 \right] N^* + \frac{\sqrt{2}}{3} C \Delta^*$$

$$\gamma n \rightarrow p \pi^- : \quad \mp \sqrt{\frac{2}{3}} \left[ C^0 \oplus \sqrt{\frac{1}{3}} C^1 \right] N^* + \frac{\sqrt{2}}{3} C \Delta^*$$

# Observable: $\sigma$

G13

## Reaction: $\gamma n \rightarrow p \pi$



- First-ever determination of the excited neutron multipoles for:  $N(1440)1/2^+$ ,  $N(1535)1/2^-$ ,  $N(1650)1/2^-$ , and  $N(1720)3/2^+$

# Observable: $\Sigma$

Reactions:  $\gamma p \rightarrow p \pi^0$  and  $\gamma p \rightarrow n \pi^+$

Configuration:

- **Linear photon polarization**
- No target polarization
- No recoil polarization

Experiments:

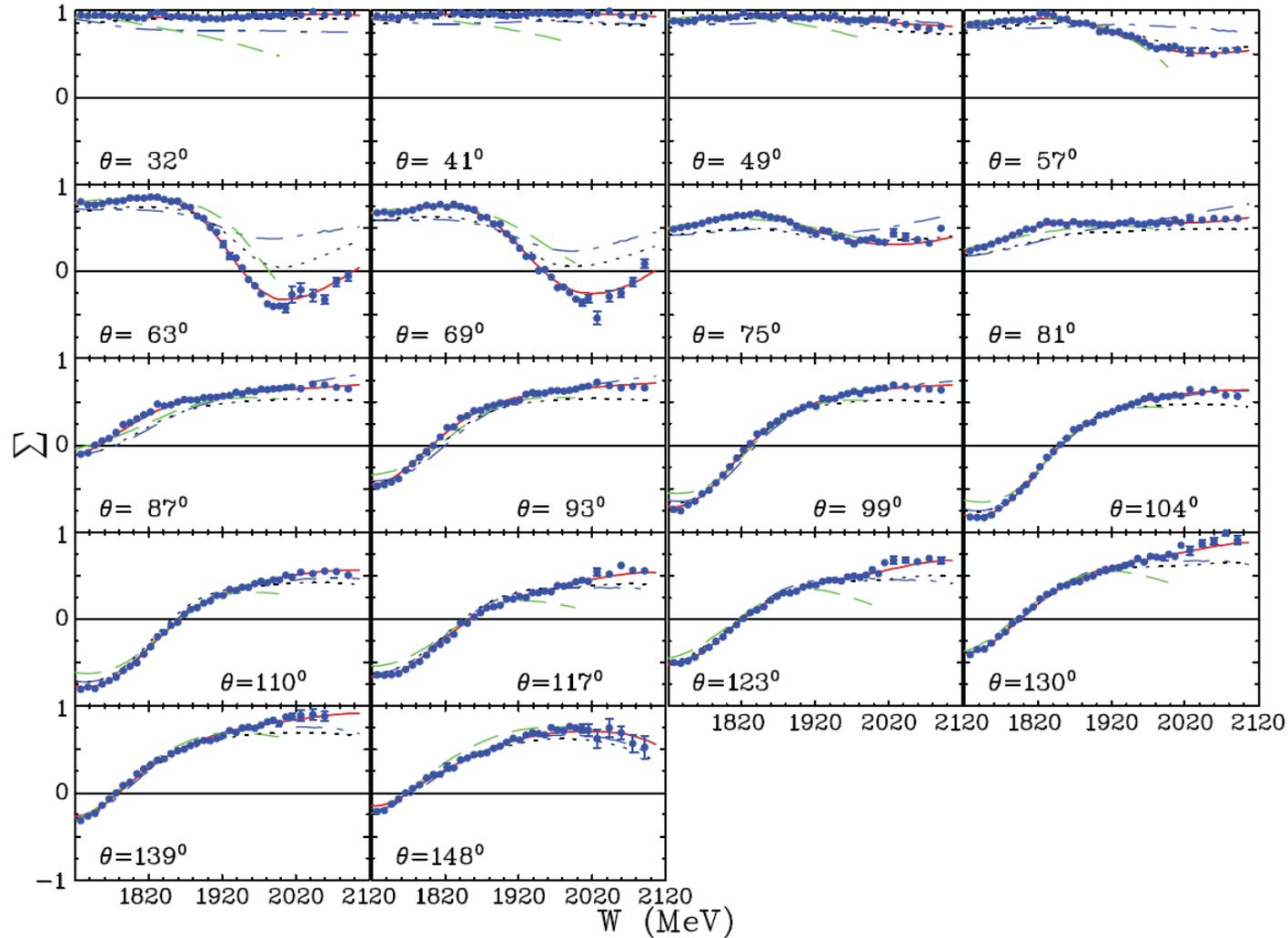
- g8b  $\rightarrow$  proton reactions
- g13  $\rightarrow$  neutron reactions

Photon		Target			Recoil			Target + Recoil			
	–	–	–	–	$x'$	$y'$	$z'$	$x'$	$x'$	$z'$	$z'$
	–	$x$	$y$	$z$	–	–	–	$x$	$z$	$x$	$z$
unpolarized	$\sigma_0$	0	$T$	0	0	$P$	0	$T_{x'}$	$-L_{x'}$	$T_{z'}$	$L_{z'}$
linear pol.	$-\Sigma$	$H$	$(-P)$	$-G$	$O_{x'}$	$(-T)$	$O_{z'}$	$(-L_{z'})$	$(T_{z'})$	$(-L_{x'})$	$(-T_{x'})$
circular pol.	0	$F$	0	$-E$	$-C_{x'}$	0	$-C_{z'}$	0	0	0	0

# $\Sigma$ for $\gamma p \rightarrow p \pi^0$

G8b

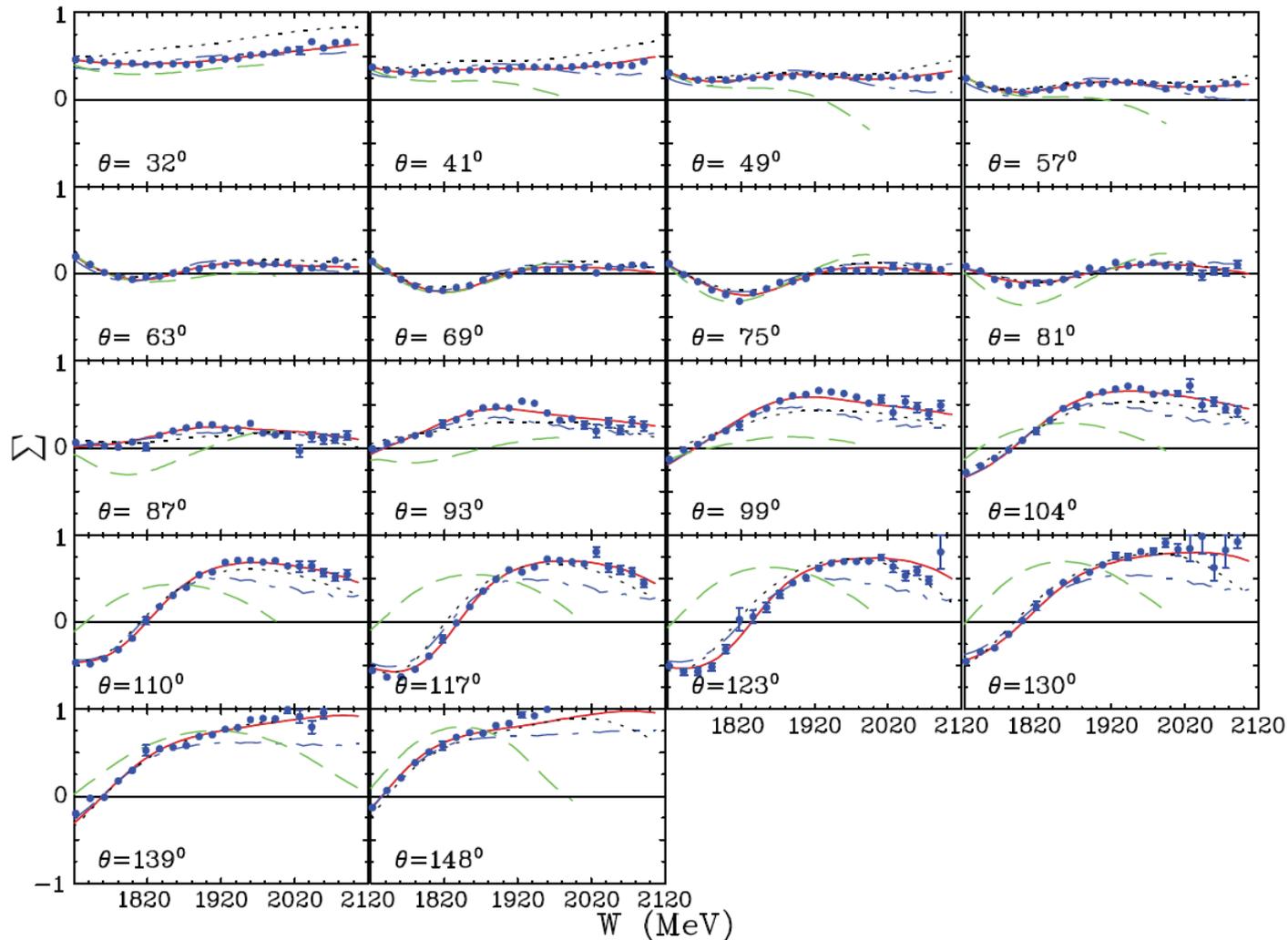
RED: SAID fit



# $\Sigma$ for $\gamma p \rightarrow n \pi^+$

G8b

RED: SAID fit

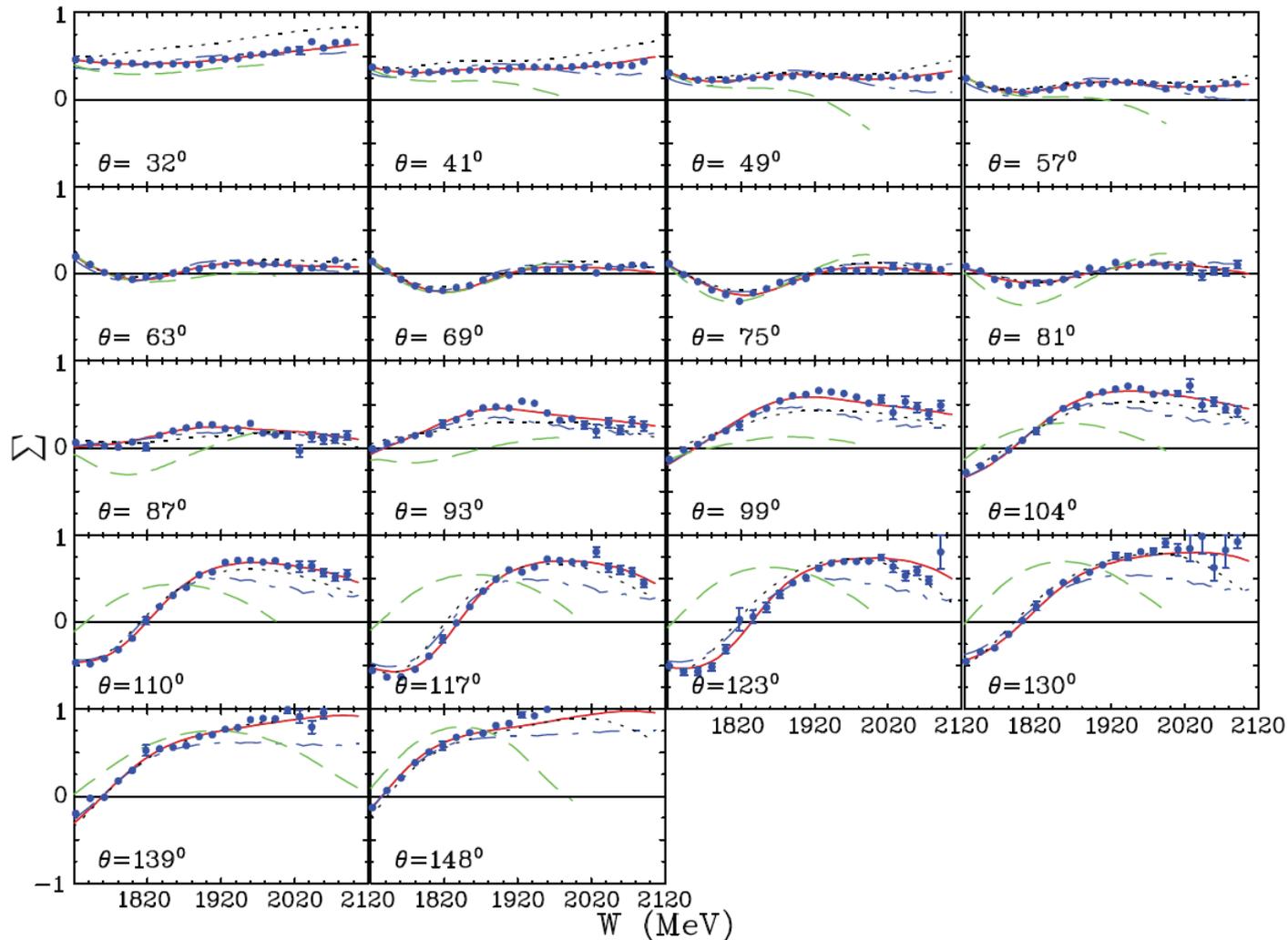


- Data for both reactions more than doubled the world database

# $\Sigma$ for $\gamma p \rightarrow n \pi^+$

G8b

RED: SAID fit



- Largest change from fits to prior  $\Sigma$  data for pions found in resonance couplings of  $\Delta(1700)3/2^-$  and  $\Delta(1905)5/2^+$

# Observable: $G$

Reactions:  $\gamma p \rightarrow p \pi^0$  and  $\gamma p \rightarrow n \pi^+$

Configuration:

- **Linear photon polarization**
- **Longitudinal target polarization**
- No recoil polarization

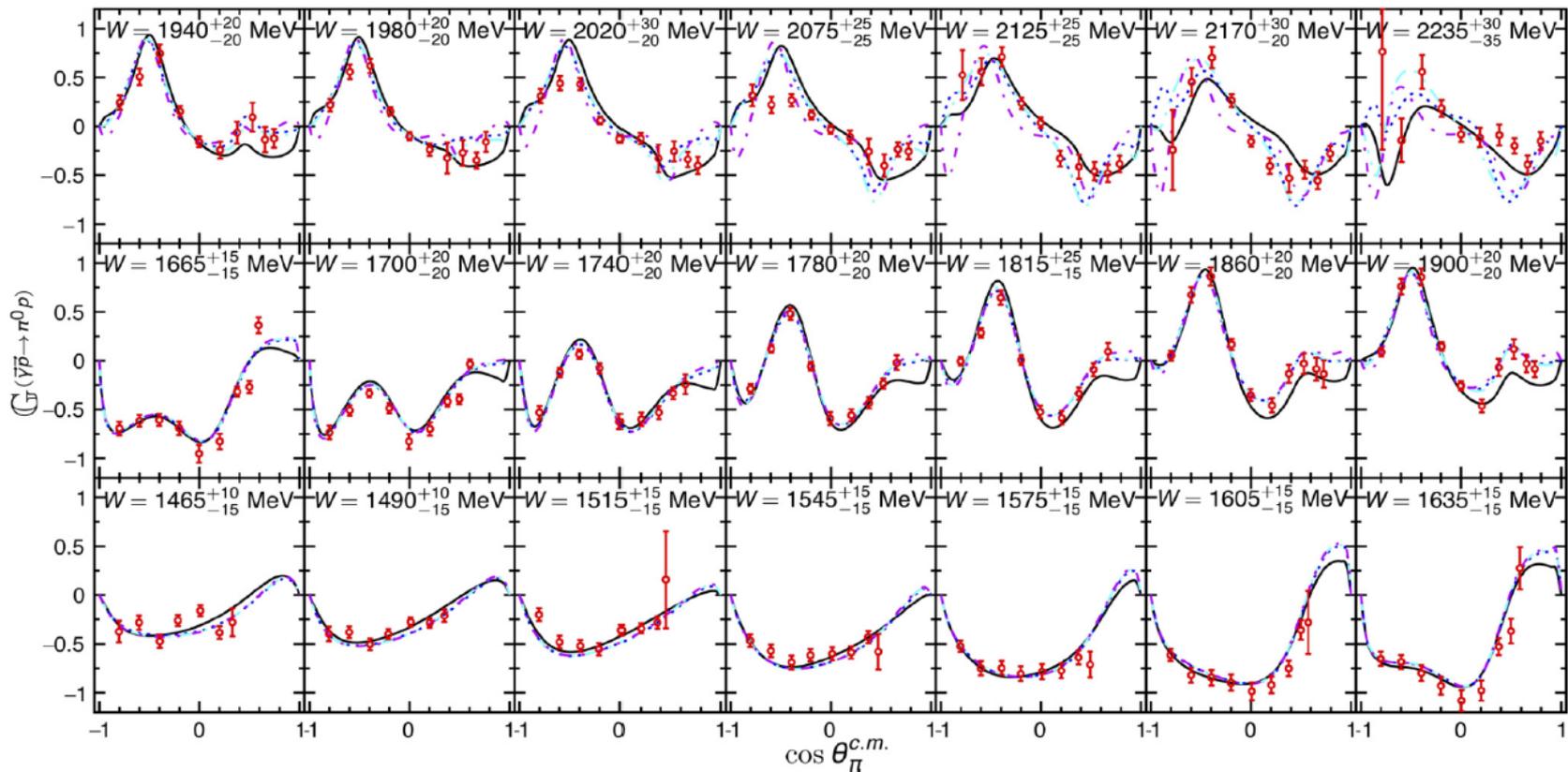
Experiment:

- g9b: FROST

Photon		Target			Recoil			Target + Recoil			
	–	–	–	↓	$x'$	$y'$	$z'$	$x'$	$x'$	$z'$	$z'$
	–	$x$	$y$	$z$	–	–	–	$x$	$z$	$x$	$z$
unpolarized	$\sigma_0$	0	$T$	0	0	$P$	0	$T_{x'}$	$-L_{x'}$	$T_{z'}$	$L_{z'}$
linear pol.	$-\Sigma$	$H$	$(-P)$	$-G$	$O_{x'}$	$(-T)$	$O_{z'}$	$(-L_{z'})$	$(T_{z'})$	$(-L_{x'})$	$(-T_{x'})$
circular pol.	0	$F$	0	$-E$	$-C_{x'}$	0	$-C_{z'}$	0	0	0	0

# $G$ for $\gamma p \rightarrow p \pi^0$

G9b: FROST

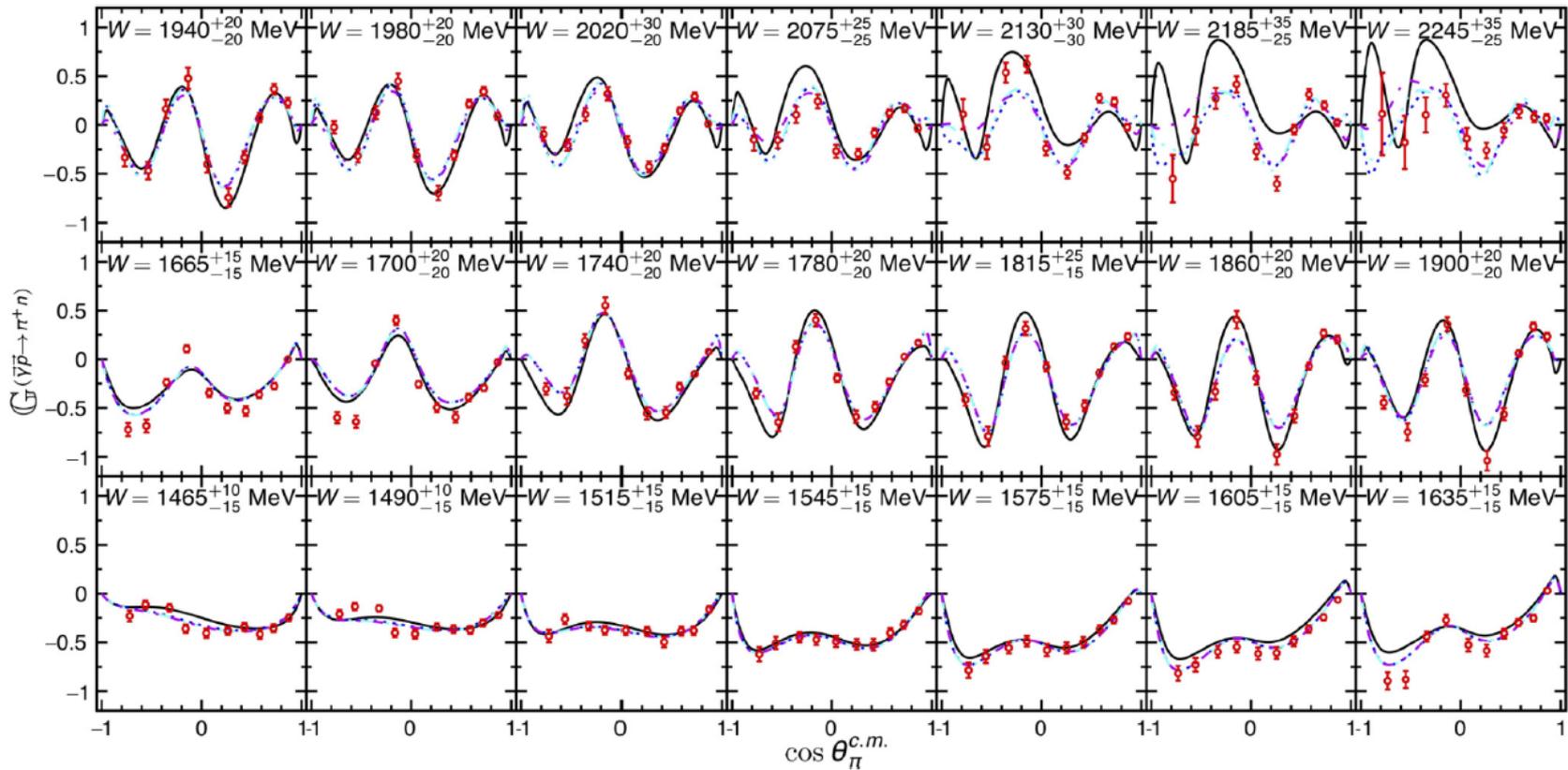


SAID: Solid black lines

Bonn-Gatchina: Dotted lines of various colors

# $G$ for $\gamma p \rightarrow n \pi^+$

G9b: FROST



Bonn-Gatchina analysis (dotted) sees important contribution from  $N(2190)7/2^-$  and  $\Delta(2200)7/2^-$

# Observables: $T$ and $F$



Configuration:

- **Circular photon polarization**
- **Transverse target polarization**
- Unpolarized photon (by adding circular beams)
- No recoil polarization

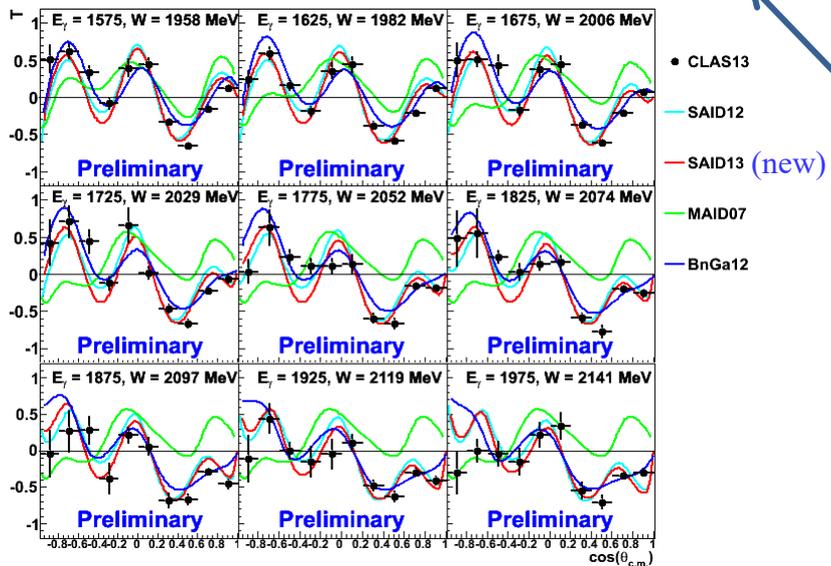
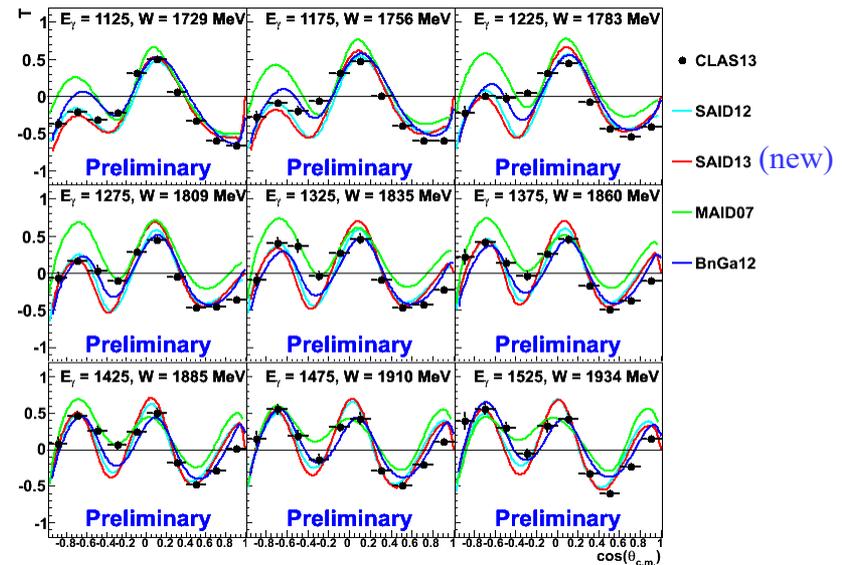
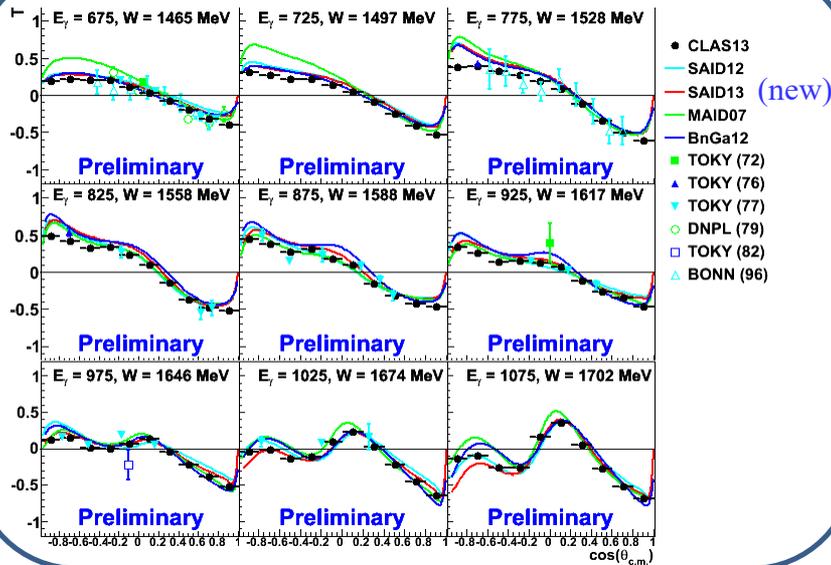
Experiment:

- g9b: FROST

Photon		Target			Recoil			Target + Recoil			
	–	↓	↓	–	$x'$	$y'$	$z'$	$x'$	$x'$	$z'$	$z'$
	–	$x$	$y$	$z$	–	–	–	$x$	$z$	$x$	$z$
→ unpolarized	$\sigma_0$	0	$T$	0	0	$P$	0	$T_{x'}$	$-L_{x'}$	$T_{z'}$	$L_{z'}$
linear pol.	$-\Sigma$	$H$	$(-P)$	$-G$	$O_{x'}$	$(-T)$	$O_{z'}$	$(-L_{z'})$	$(T_{z'})$	$(-L_{x'})$	$(-T_{x'})$
→ circular pol.	0	$F$	0	$-E$	$-C_{x'}$	0	$-C_{z'}$	0	0	0	0

# $T$ for $\gamma p \rightarrow n \pi^+$

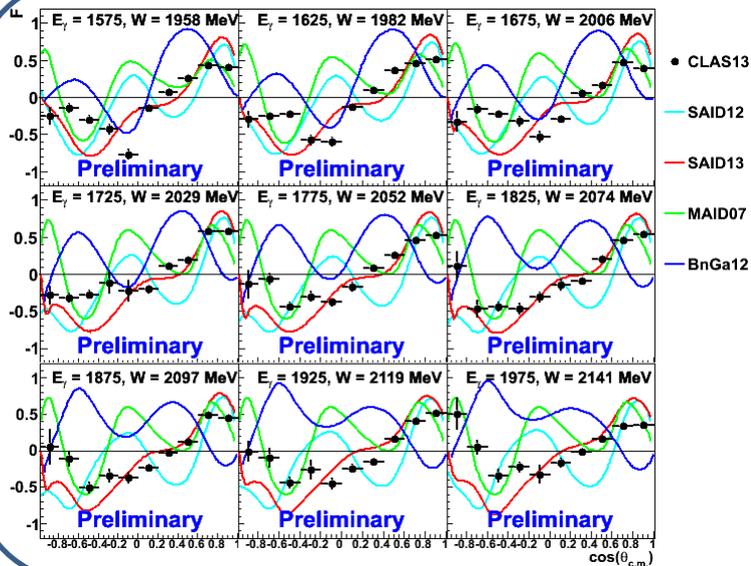
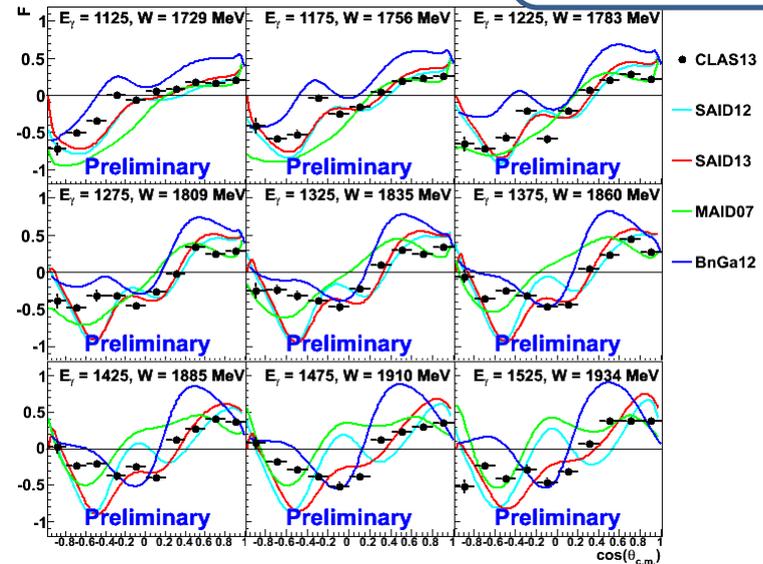
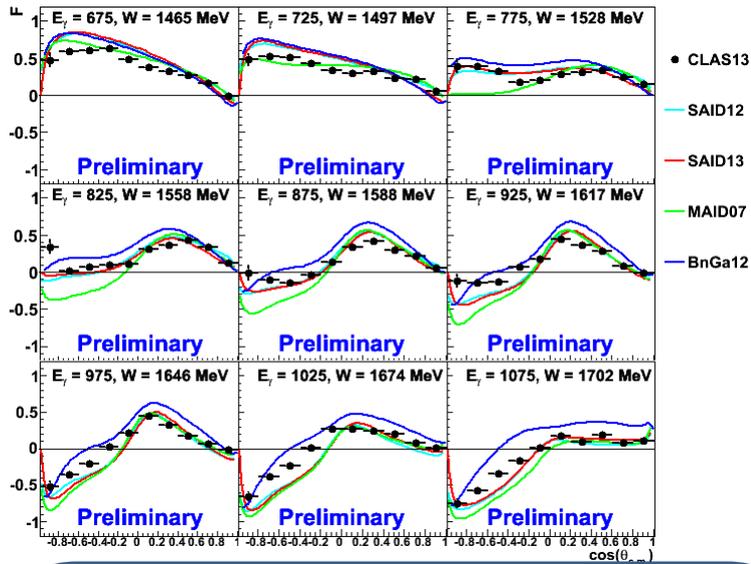
G9b: FROST



- Early stage results
- CLAS results agree well with previous data

# $F$ for $\gamma p \rightarrow n \pi^+$

G9b: FROST



- Early stage results
- Predictions get worse at higher energies

# Observable: $E$

Reactions:  $\gamma p \rightarrow n \pi^+, p \pi^0$  and  $\gamma n \rightarrow p \pi$

Configuration:

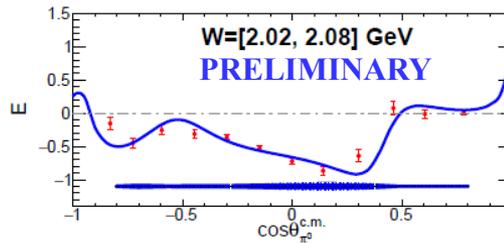
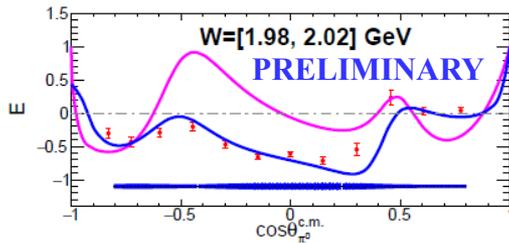
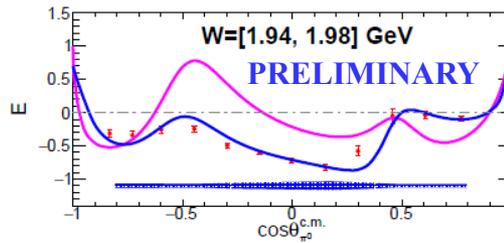
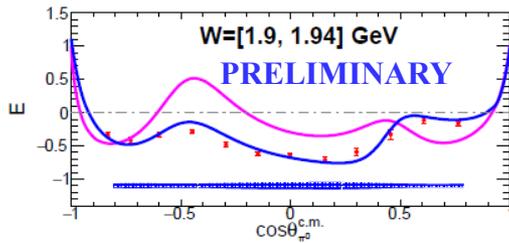
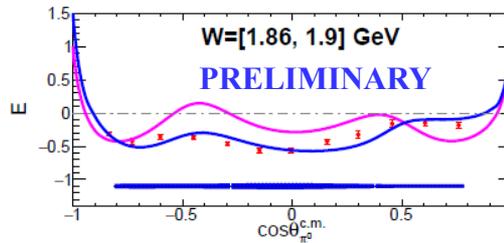
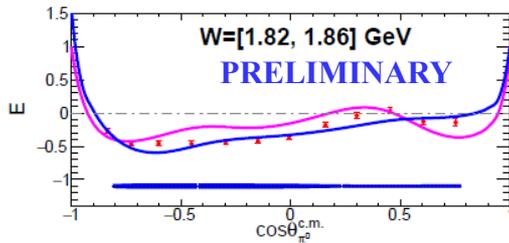
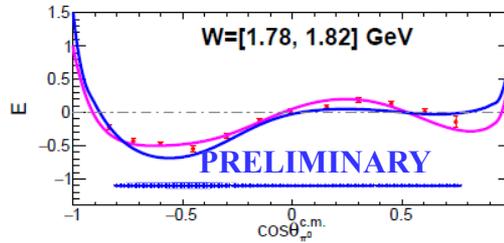
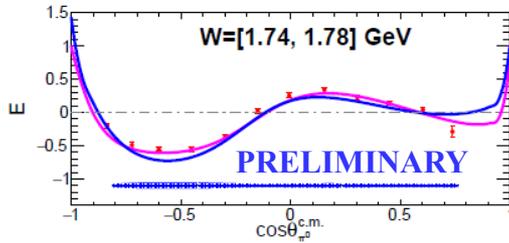
- **Circular photon polarization**
- **Longitudinal Target polarization**
- No recoil polarization

Experiments:

- g9a: FROST  $\rightarrow$  proton reactions
- g14: HDICE  $\rightarrow$  neutron reactions

Photon	Target			Recoil			Target + Recoil				
	$-$	$x$	$y$	$z$	$x'$	$y'$	$z'$	$x'$	$x'$	$z'$	$z'$
unpolarized	$\sigma_0$	0	$T$	0	0	$P$	0	$T_{x'}$	$-L_{x'}$	$T_{z'}$	$L_{z'}$
linear pol.	$-\Sigma$	$H$	$(-P)$	$-G$	$O_{x'}$	$(-T)$	$O_{z'}$	$(-L_{z'})$	$(T_{z'})$	$(-L_{x'})$	$(-T_{x'})$
$\rightarrow$ circular pol.	0	$F$	0	$-E$	$-C_{x'}$	0	$-C_{z'}$	0	0	0	0

# $E$ for $\gamma p \rightarrow p \pi^0$

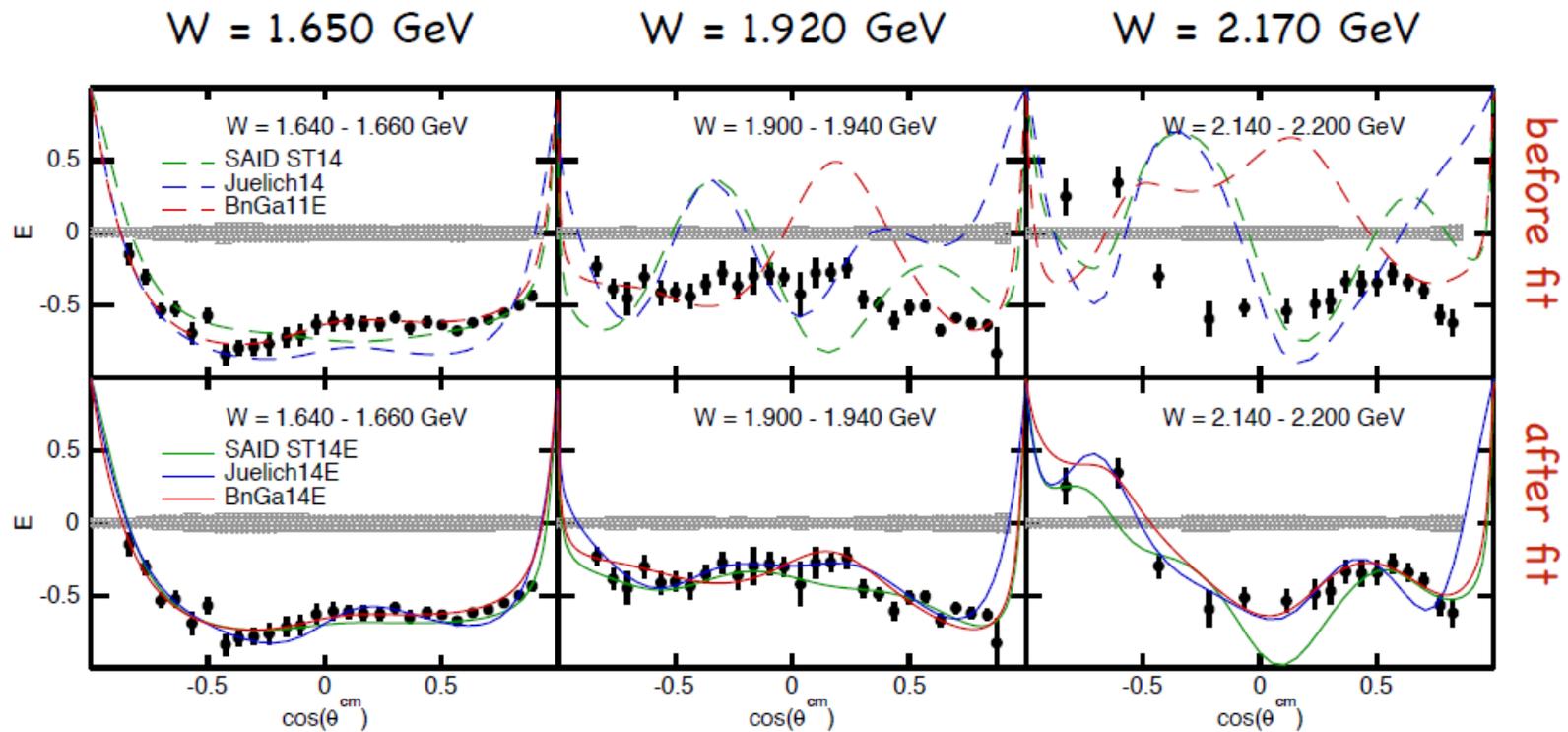


- Sample of results taken from analysis note

- **Blue** lines: SAID

- **Magenta** lines: MAID

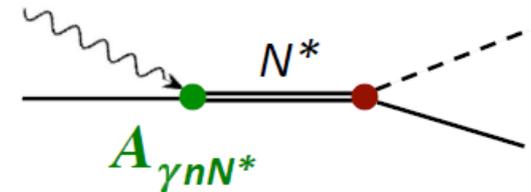
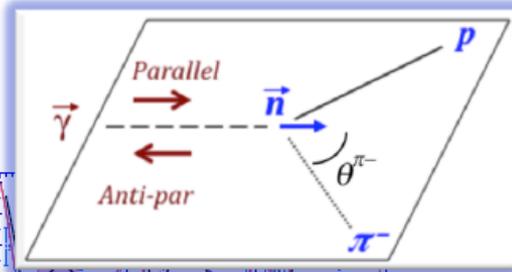
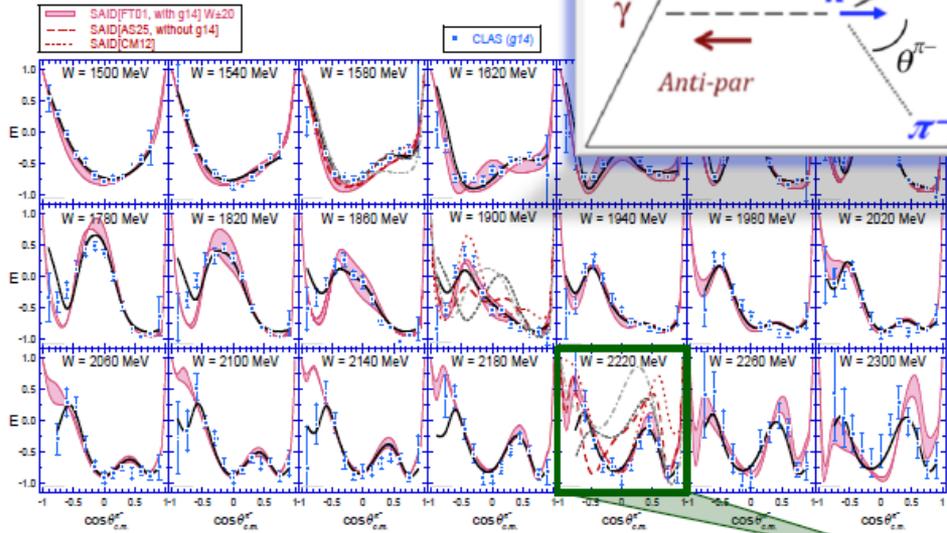
# Selected results of FROST Experiment $\vec{\gamma}\vec{p} \rightarrow \pi^+n$



- FROST experiment produced 900 data points of the **double-polarization observable  $E$**  in  $\pi^+$  photoproduction with circularly polarized beam on longitudinally polarized protons for  $W = 1240 - 2260$  MeV.
- Significant improvements of the description of the data in SAID, Jülich, and BnGa partial-wave analyses after fitting.
- **New evidence found in this data for a  $\Delta(2200)7/2^-$  resonance (BnGa analysis).**

# $g_{14}$ beam-target helicity asymmetries for $\gamma n \rightarrow \pi^- p$ and $N^*$ states excited from the neutron

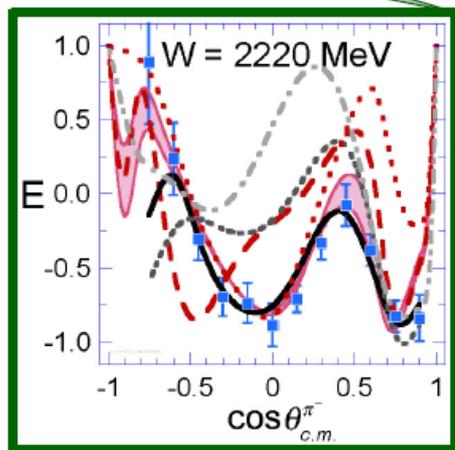
- 1<sup>st</sup> double-polarized  $\vec{n}$  data  
PRL **118** (2017) 242002



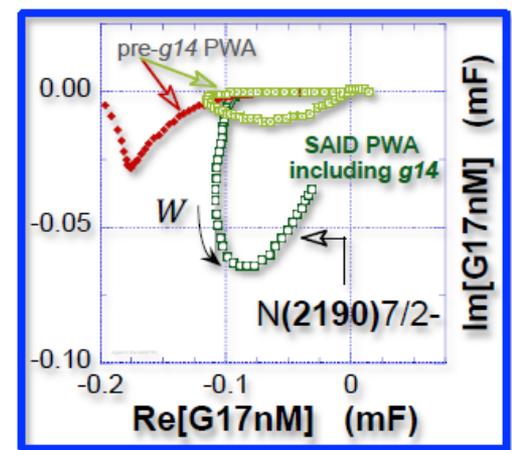
- E&M interaction is not isospin symmetric
- $\gamma n N^*$  and  $\gamma p N^*$  couplings are different  
 $\Leftrightarrow$  probes of dynamics in  $N^*$  excitation

- eg. SAID Partial Wave Analysis (PWA):  
 $A_{\gamma n}^{1/2} [N(2190)7/2^-] \rightarrow -16 \pm 5 (10^{-3} \text{ GeV}^{-1/2})$   
 $A_{\gamma n}^{3/2} [N(2190)7/2^-] \rightarrow -35 \pm 5 (10^{-3} \text{ GeV}^{-1/2})$

- very little previous spin-dependent  $\gamma n$  data exists
- for invariant masses ( $W$ ) over 1800 MeV, predictions from previous Partial Wave Analyses (PWA) fail badly
- $\vec{\gamma} \vec{n}$  data probes  $N^*$  states

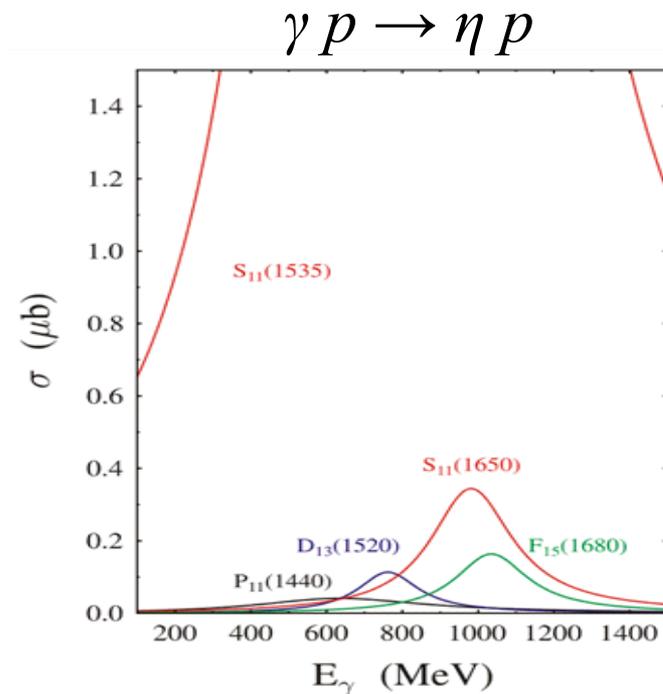
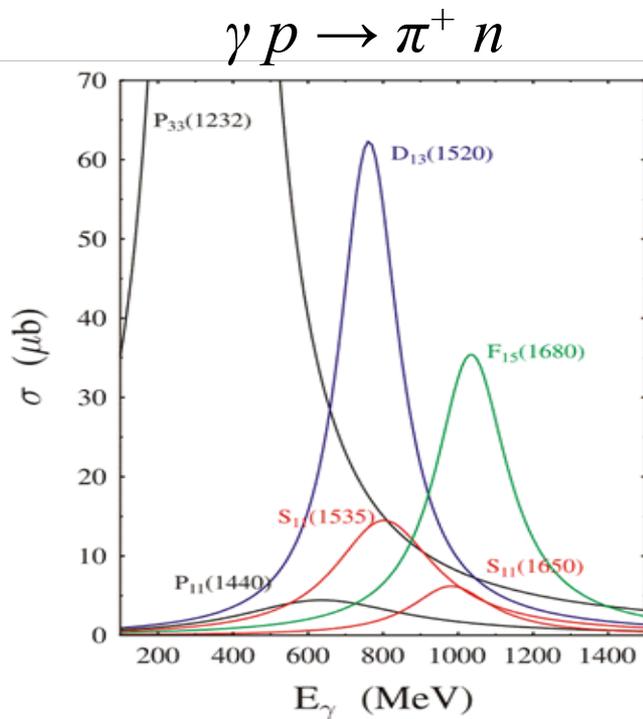


$\rightarrow$  PWA  $\rightarrow$



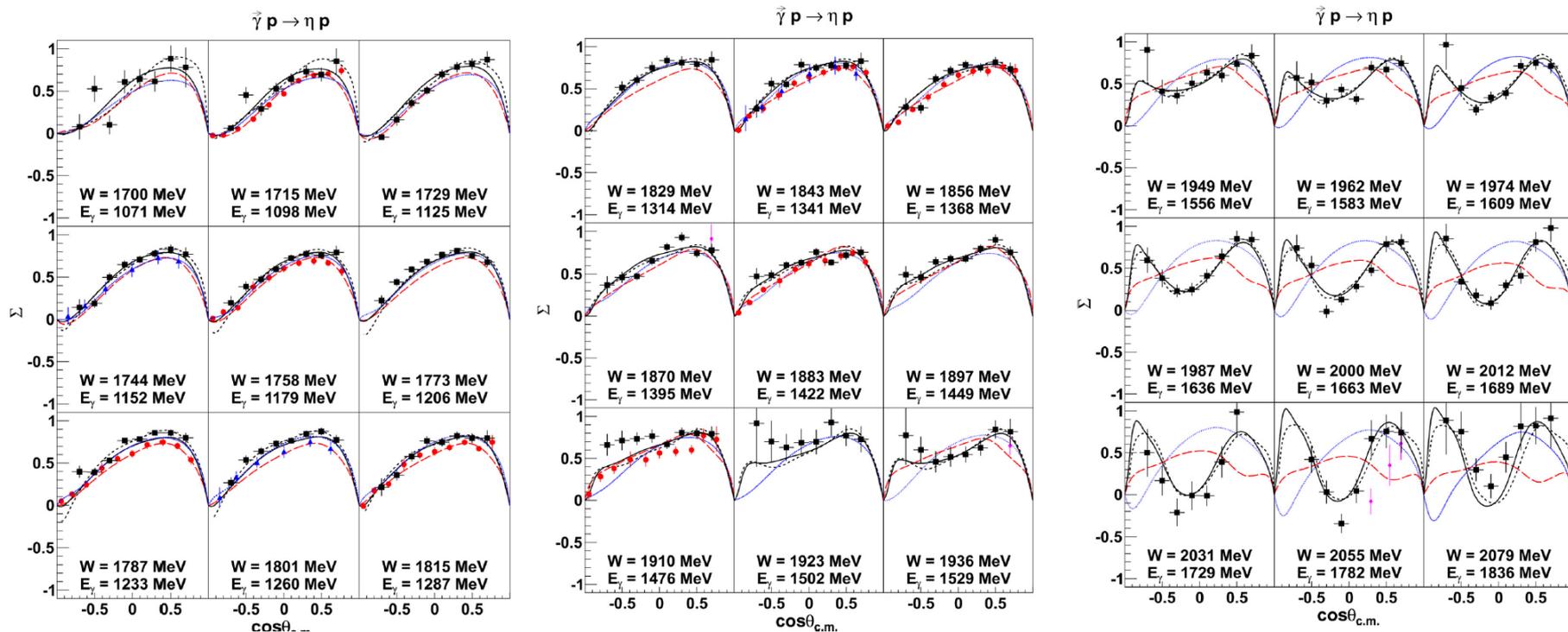
# “Isospin filters”

- The  $\eta p$ ,  $\omega p$  and  $K^+ \Lambda$  systems have isospin  $\frac{1}{2}$  and limit one-step excited states of the proton to be isospin  $\frac{1}{2}$ . The final states  $\eta p$ ,  $\omega p$ , and  $K^+ \Lambda$  act as **isospin filters** to the resonance spectrum.



# $\Sigma$ for $\eta$

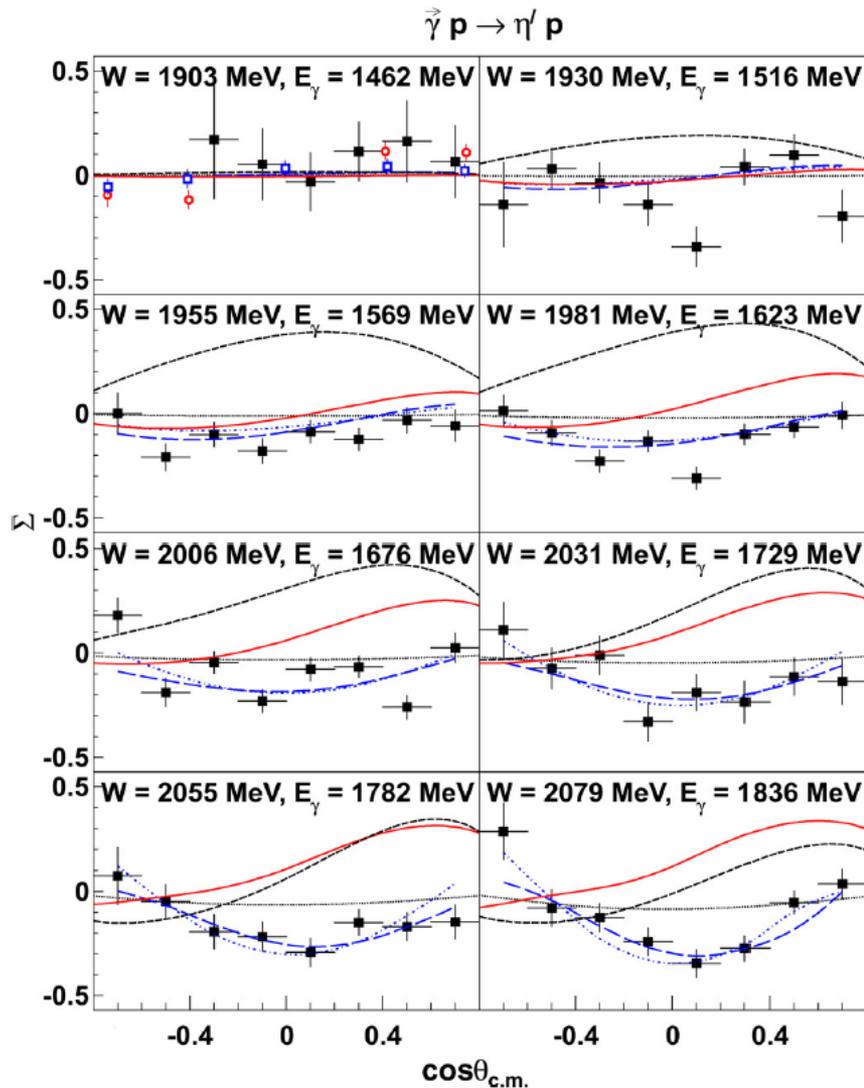
G8b



- Fit to Julich Bonn model (black line) with presence of  $N(1900)3/2^-$  (solid) and without (dashed)
- The inclusion of the  $N(1900)3/2^+$  was found to be important by Bonn-Gatchina for  $K\Lambda$  and  $K\Sigma$  photoproduction

# $\Sigma$ for $\eta'$

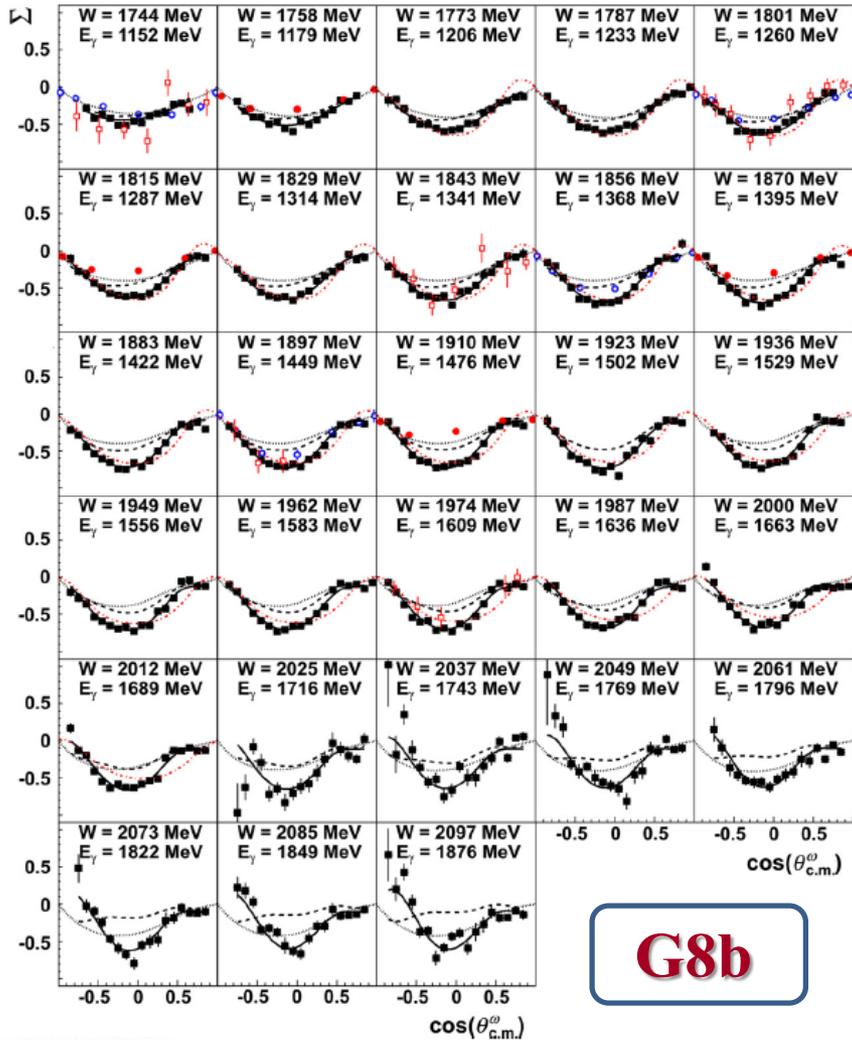
G8b



- Fit to Bonn-Gatchina model (blue lines) indicates presence of  $N(1895)1/2^-$ ,  $N(2100)1/2^+$ ,  $N(2120)3/2^-$  and strong presence of  $N(1900)1/2^-$

# $\Sigma$ for $\omega$

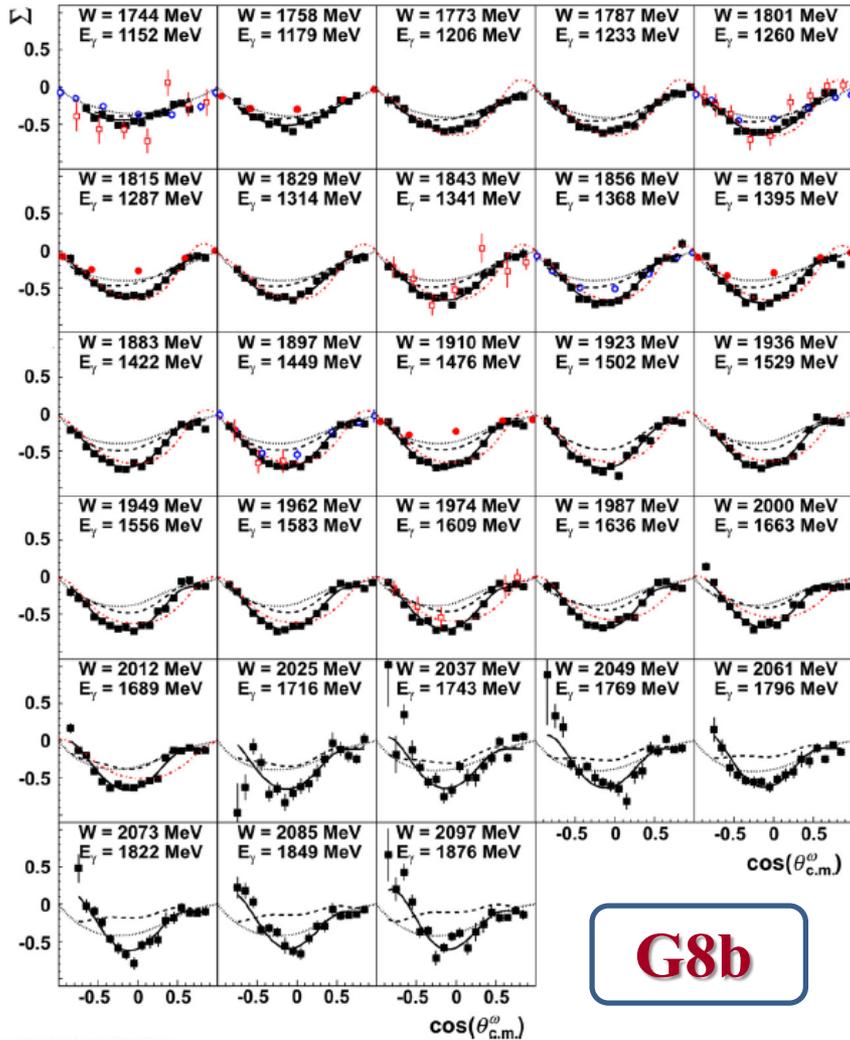
$$\vec{\gamma} p \rightarrow p \omega$$



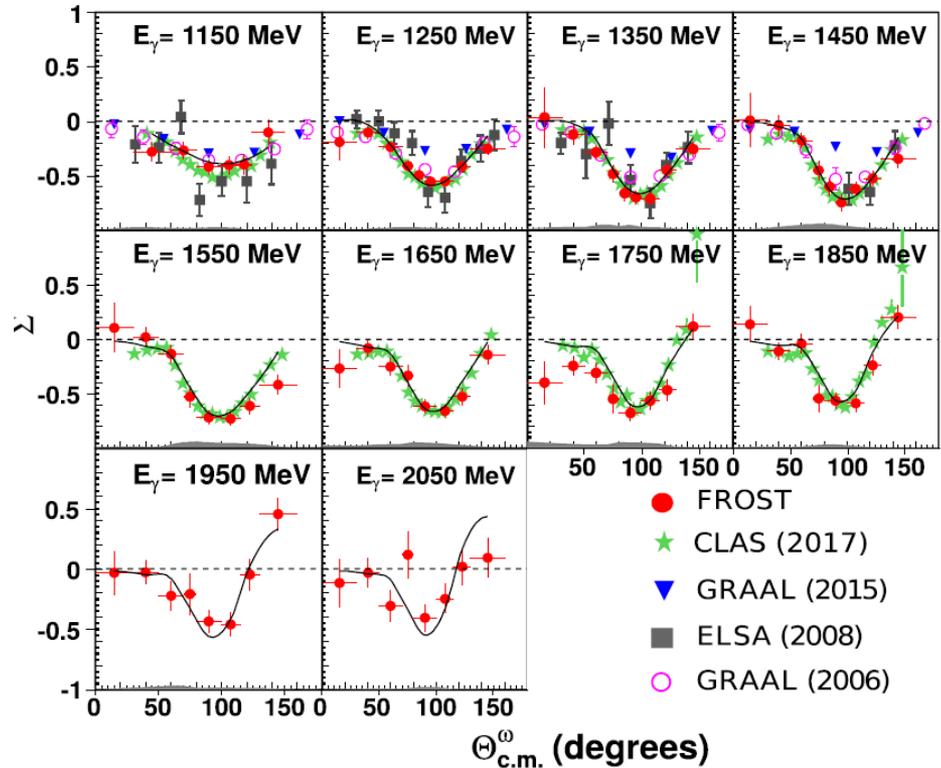
G8b

# $\Sigma$ for $\omega$

$$\vec{\gamma} p \rightarrow p \omega$$



**G8b**

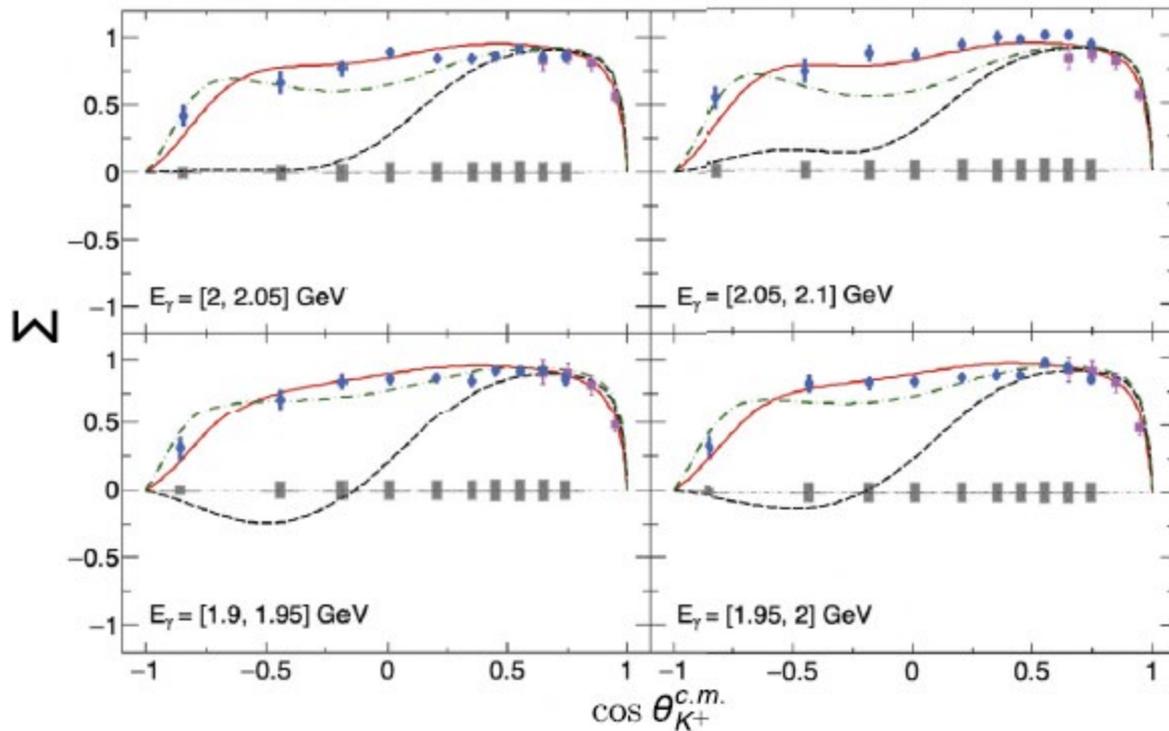


**G9a: FROST**

P. Roy, *et al.*, (CLAS Collaboration), Phys. Rev. C 97, 055202 (2018)

# Beam asymmetries for $\gamma n \rightarrow K^+ \Sigma^-$

G13



**Red:** Full solution (Bonn-Gatchina)

**Black:** Contribution of  $N(1720)3/2^+$  removed

**Green:** Contribution of  $\Delta(1900)1/2^-$  removed

# Observable: $T, F, P$ and $H$



Configuration:

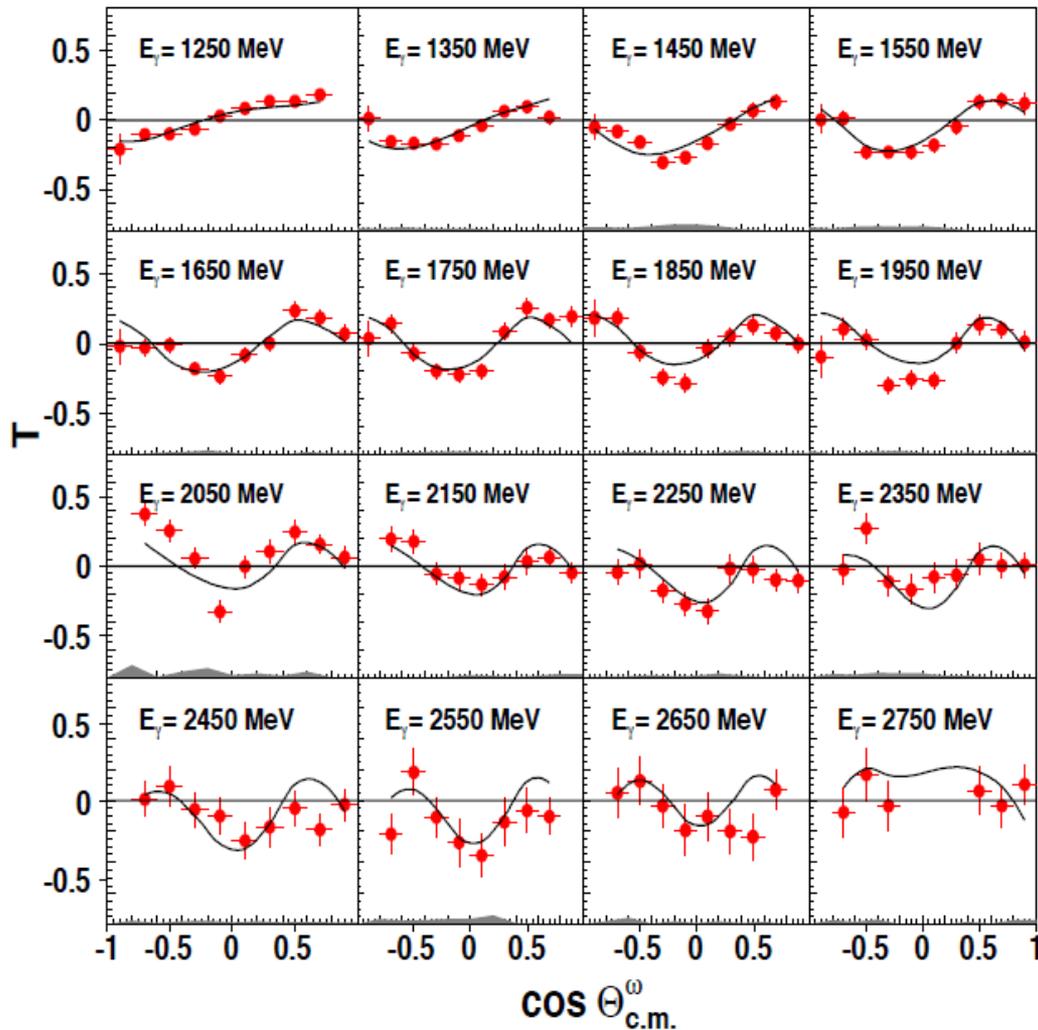
- **Circular photon polarization**
- **Transverse target polarization**
- **Unpolarized photon** (by adding circular beams)
- No recoil polarization

Experiment:

- g9b: FROST

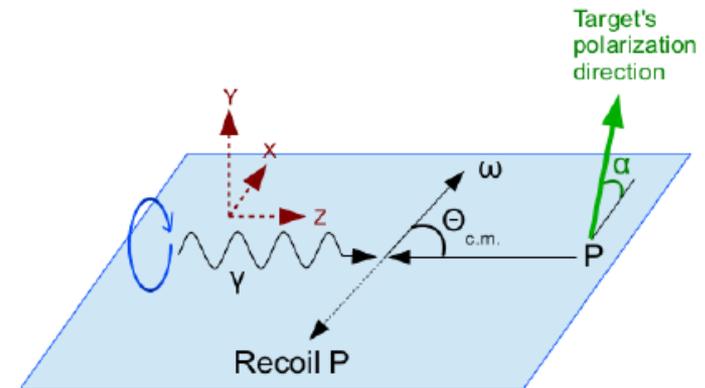
	Photon	Target			Recoil			Target + Recoil				
		$x$	$y$	$z$	$x'$	$y'$	$z'$	$x'$	$x'$	$z'$	$z'$	
	—	$\downarrow$	$\downarrow$	—	$x'$	$y'$	$z'$	$x'$	$x'$	$z'$	$z'$	
	—	$x$	$y$	$z$	—	—	—	$x$	$z$	$x$	$z$	
→	unpolarized	$\sigma_0$	0	$T$	0	$P$	0	$T_{x'}$	$-L_{x'}$	$T_{z'}$	$L_{z'}$	
→	linear pol.	$-\Sigma$	$H$	$(-P)$	$-G$	$O_{x'}$	$(-T)$	$O_{z'}$	$(-L_{z'})$	$(T_{z'})$	$(-L_{x'})$	$(-T_{x'})$
→	circular pol.	0	$F$	0	$-E$	$-C_{x'}$	0	$-C_{z'}$	0	0	0	0

# Target Asymmetry $T$ in $\gamma \vec{p} \rightarrow p \omega$ (CLAS g9b)



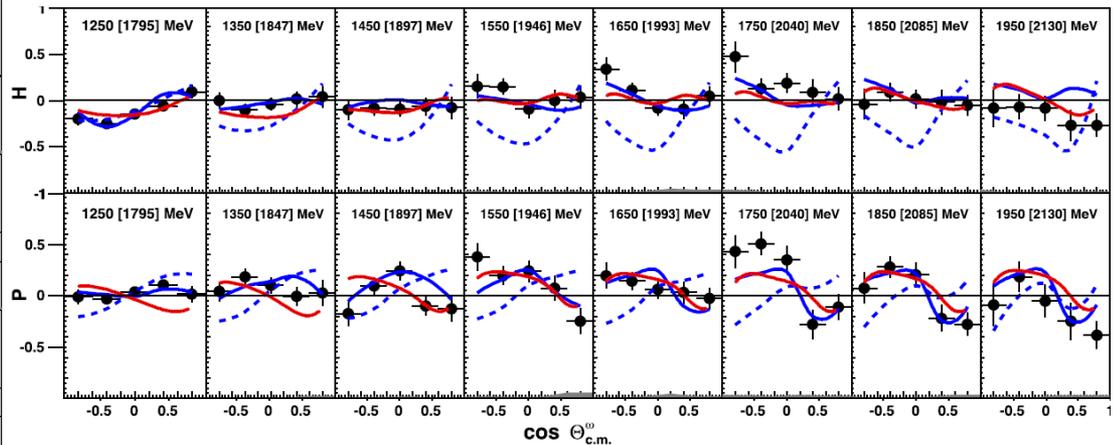
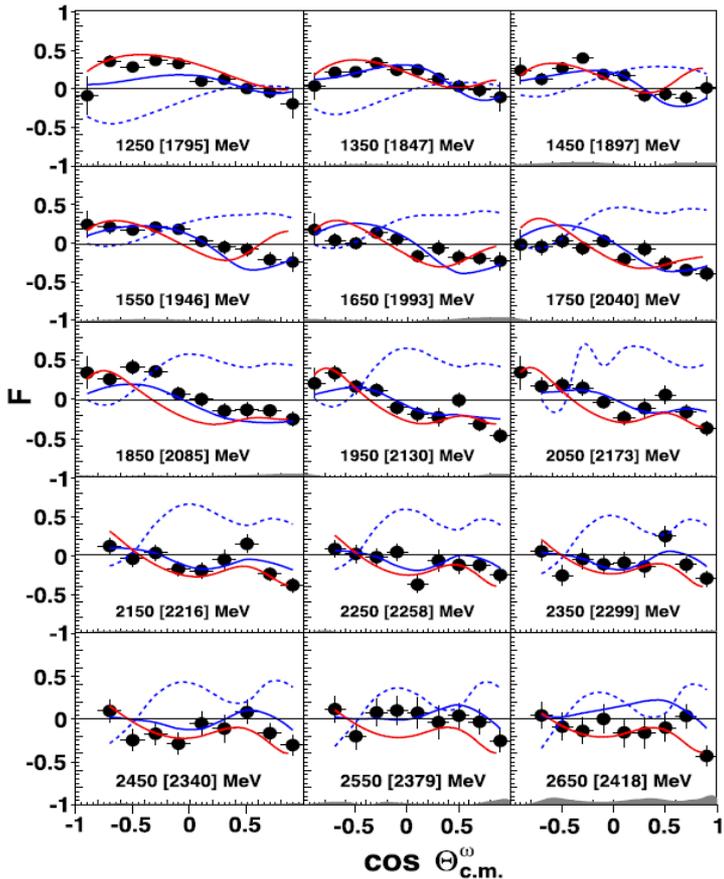
## Polarized Cross Section

$$\frac{d\sigma}{d\Omega} = \sigma_0 \left\{ 1 - \delta_I \Sigma \cos 2\phi \right. \\ \left. + \Lambda_x (-\delta_I H \sin 2\phi + \delta_{\odot} F) \right. \\ \left. - \Lambda_y (-T + \delta_I P \cos 2\phi) \right. \\ \left. - \Lambda_z (-\delta_I G \sin 2\phi + \delta_{\odot} E) \right\}$$



P. Roy *et al.* [CLAS Collaboration], Phys. Rev. C **97**, no. 5, 055202 (2018)

# $F$ , $P$ and $H$ for $\omega$



- Red : Wei
- Blue : Bon-Gatchina, where dashed = old
- Indicates notable contributions from  $N(1875)3/2^-$ ,  $N(2120)3/2^-$  and  $N(1880)1/2^+$

# Observable: $E$

Reactions:  $\gamma p \rightarrow p \omega, p \eta$  and  $\gamma n \rightarrow K^+ \Sigma^-$

Configuration:

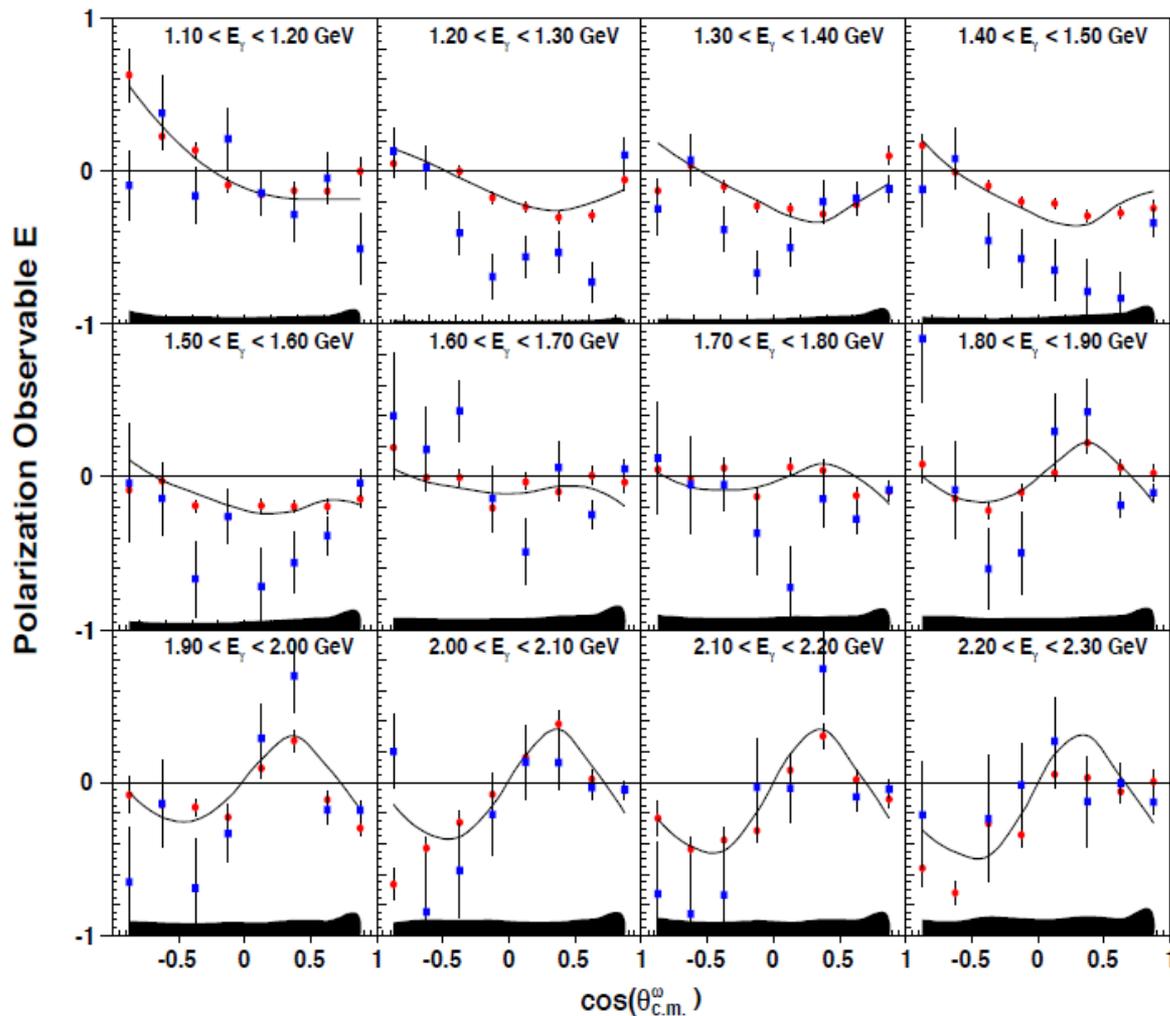
- **Circular photon polarization**
- **Longitudinal Target polarization**
- No recoil polarization

Experiment:

- g9b: FROST
- g14: HD-ICE

Photon		Target			Recoil			Target + Recoil			
	–	–	–	↓	$x'$	$y'$	$z'$	$x'$	$x'$	$z'$	$z'$
	–	$x$	$y$	$z$	–	–	–	$x$	$z$	$x$	$z$
unpolarized	$\sigma_0$	0	$T$	0	0	$P$	0	$T_{x'}$	$-L_{x'}$	$T_{z'}$	$L_{z'}$
linear pol.	$-\Sigma$	$H$	$(-P)$	$-G$	$O_{x'}$	$(-T)$	$O_{z'}$	$(-L_{z'})$	$(T_{z'})$	$(-L_{x'})$	$(-T_{x'})$
→ circular pol.	0	$F$	0	$-E$	$-C_{x'}$	0	$-C_{z'}$	0	0	0	0

# Helicity Asymmetry in $\vec{\gamma} \vec{p} \rightarrow p \omega$ (CLAS g9a)



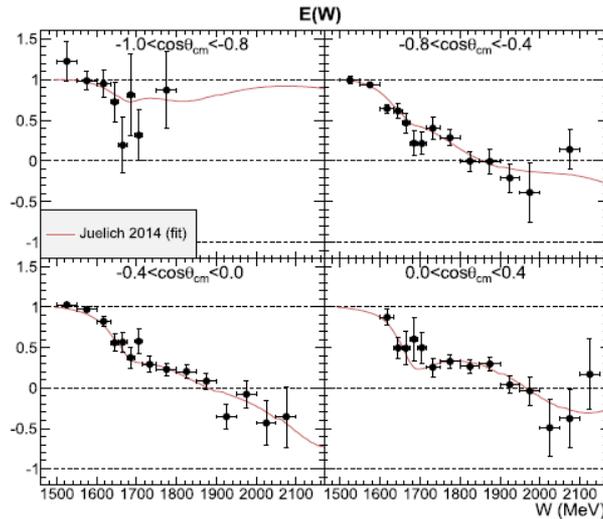
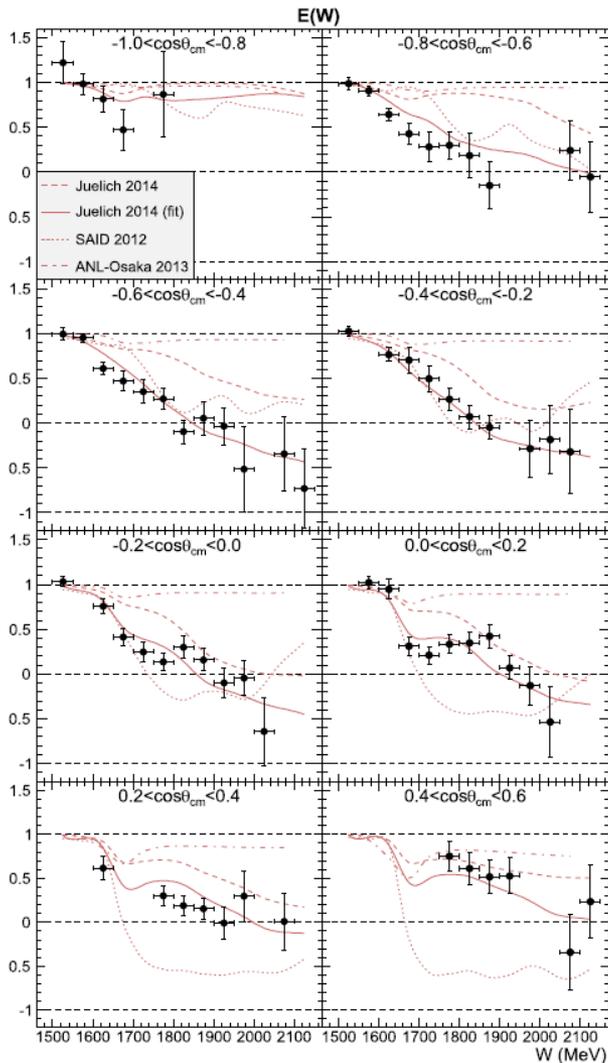
BnGa (coupled-channels) PWA

- Dominant **P** exchange
- Complex  $3/2^+$  wave
  - ①  $N(1720)$
  - ②  $W \approx 1.9$  GeV
- $N(1895) 1/2^-$  (new state)
- $N(1680), N(2000) 5/2^+$
- $7/2$  wave  $> 2.1$  GeV

- CLAS-g9a
  - CBELSA/TAPS
- Phys. Lett. B **750**, 453 (2015)

Z. Akbar *et al.* [CLAS Collaboration], Phys. Rev. C **96**, no. 6, 065209 (2017)

# $E$ for $\eta$

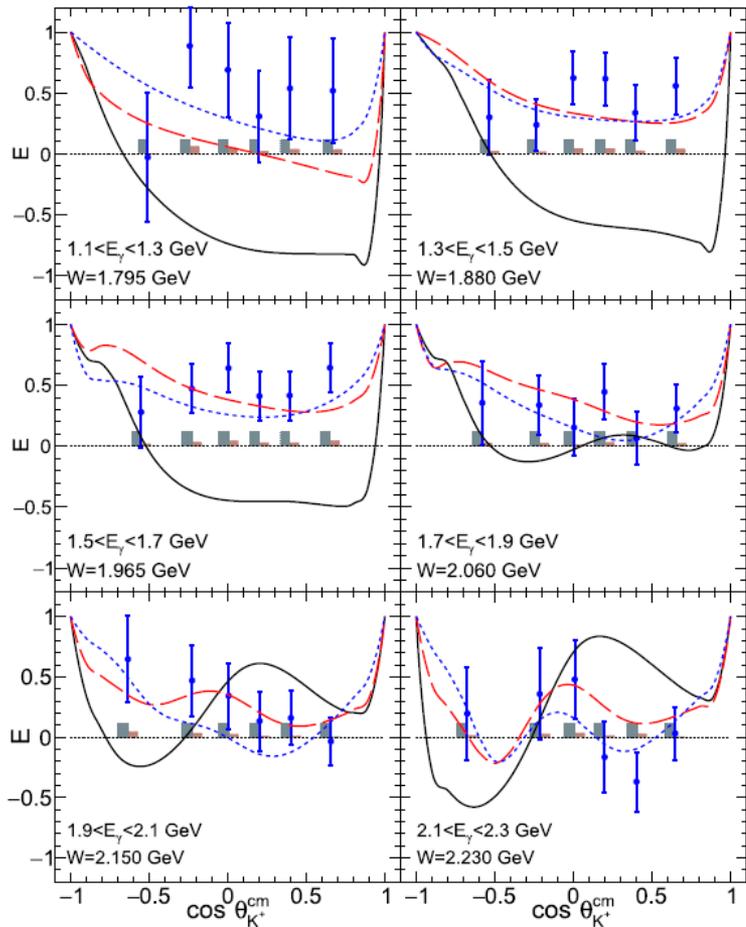


**G9a: FROST**

- Fit to Julich-Bonn model (red lines) does not indicate the need for a narrow resonance  $\sim 1.7$  GeV
- Structure near  $\sim 1.7$  GeV appears to be interference of  $E_0^+$  and  $M_2^+$  multipoles

# $E$ for $\gamma n \rightarrow K^+ \Sigma^-$

G14: HD-ICE



**Red:** Bonn-Gatchina prior to fit  
**Blue:** Full fit including “missing”  $D_{13}$   
**Black:** Full fit without  $D_{13}$

# Self-analyzing reaction $K^+ Y$ (hyperon)

- The weak decay of the hyperon allows the extraction of the hyperon polarization by looking at the decay distribution of the baryon in the hyperon center of mass system:

$$I(\cos \theta) = \frac{1}{2} (1 + \alpha P_Y \cos \theta)$$

where  $I$  is the decay distribution of the baryon,  $\alpha$  is the weak decay asymmetry ( $\alpha_{\Lambda} = 0.642$  and  $\alpha_{\Sigma^0} = -\frac{1}{3} \alpha_{\Lambda}$ ), and  $P_Y$  is the hyperon polarization.

- We can obtain recoil polarization information without a recoil polarimeter and the reaction is said to be **“self-analyzing”**

# Observables: $\Sigma$ , $T$ , $O_x$ , $O_z$

Reaction:  $\gamma p \rightarrow K^+ \Lambda, K^+ \Sigma$

Configuration:

- **Linear photon polarization**
- **Recoil polarization self analyzed**
- No target polarization

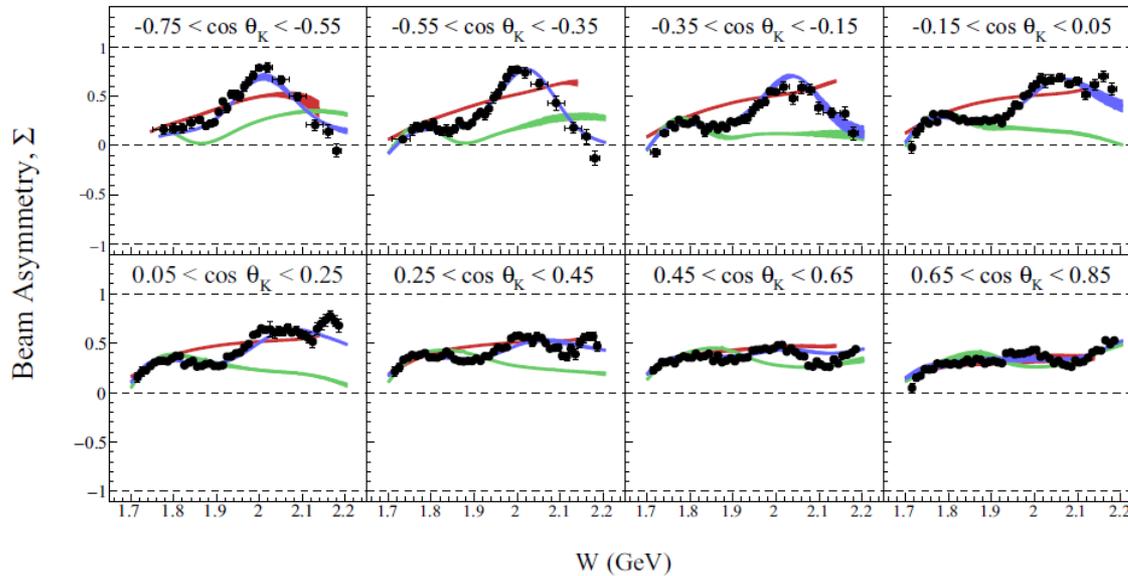
Experiments:

- g8b  $\rightarrow$  proton reactions
- g13  $\rightarrow$  neutron reactions

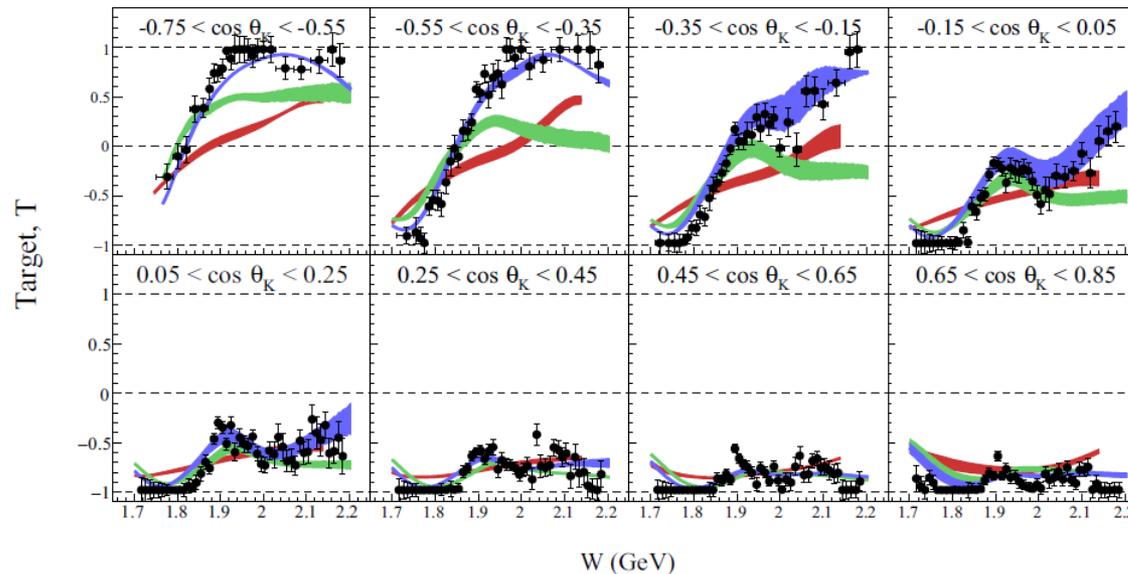
Photon		Target			Recoil			Target + Recoil			
	–	–	–	–	$x'$	$y'$	$z'$	$x'$	$x'$	$z'$	$z'$
	–	$x$	$y$	$z$	–	–	–	$x$	$z$	$x$	$z$
unpolarized	$\sigma_0$	0	$T$	0	0	$P$	0	$T_{x'}$	$-L_{x'}$	$T_{z'}$	$L_{z'}$
→ linear pol.	$-\Sigma$	$H$	$(-P)$	$-G$	$O_{x'}$	$(-T)$	$O_{z'}$	$(-L_{z'})$	$(T_{z'})$	$(-L_{x'})$	$(-T_{x'})$
circular pol.	0	$F$	0	$-E$	$-C_{x'}$	0	$-C_{z'}$	0	0	0	0

# $\Sigma, T$ for $\gamma p \rightarrow K^+ \Lambda$

G8b

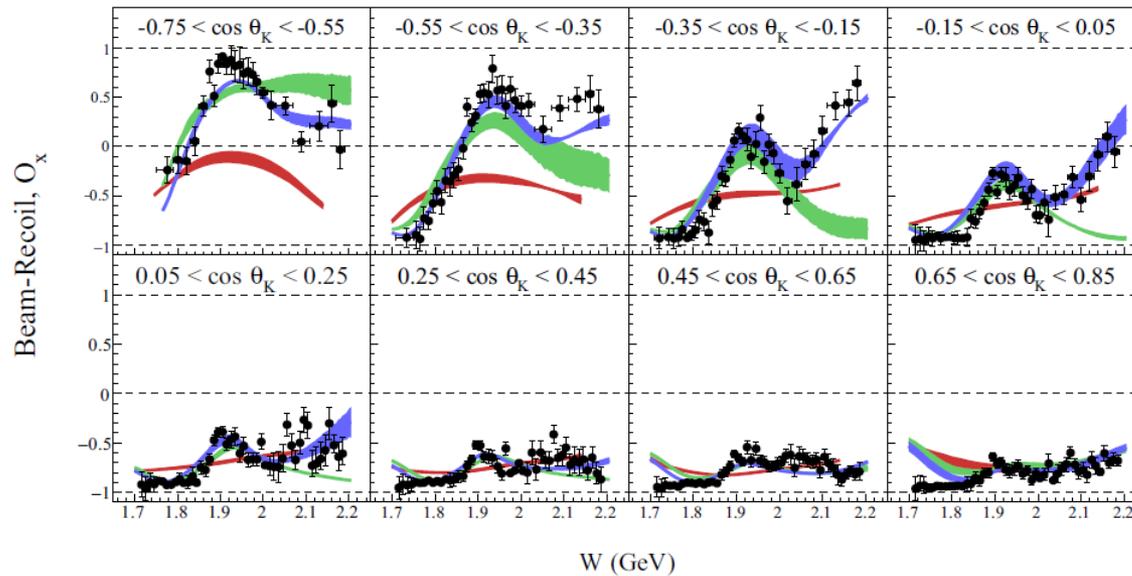


- Blue lines represent fits to Bonn-Gatchina model
- Other lines represent various predictions

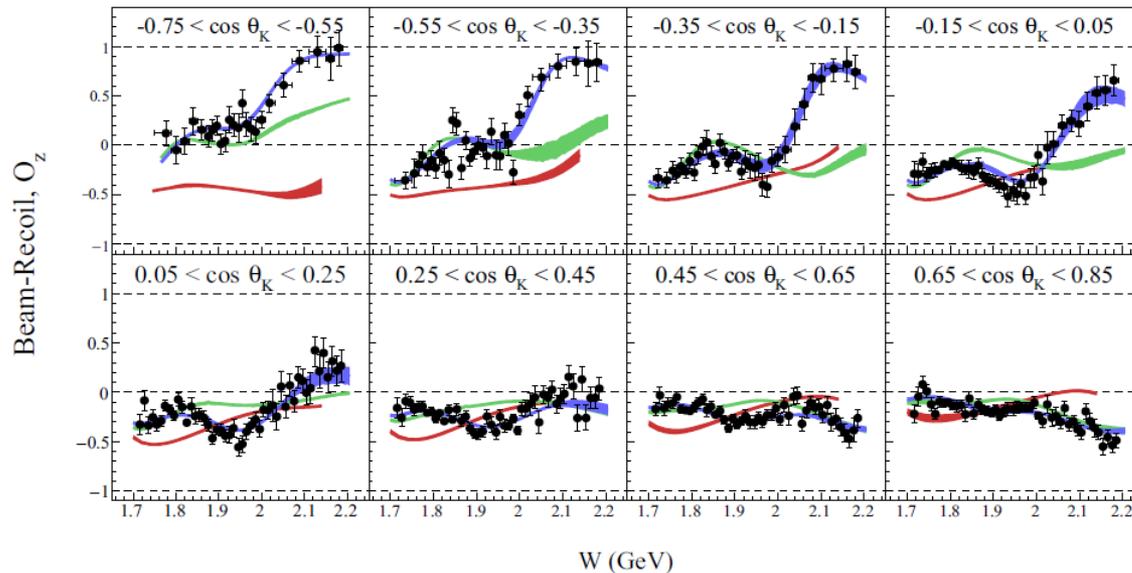


# $O_x, O_z$ for $\gamma p \rightarrow K^+ \Lambda$

G8b

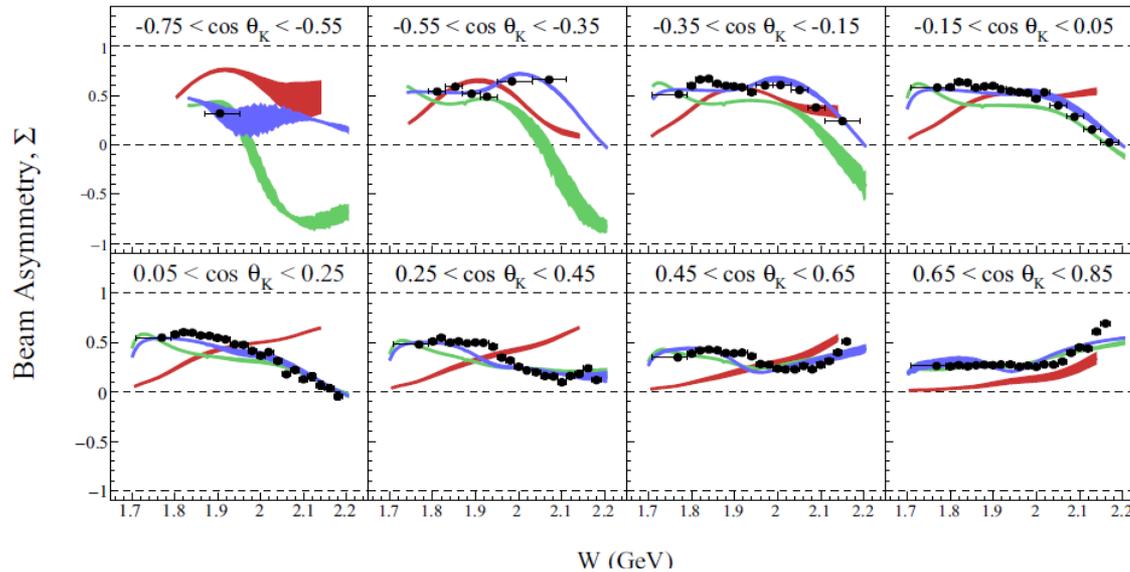


- Blue lines represent fits to Bonn-Gatchina model
- Other lines represent various predictions

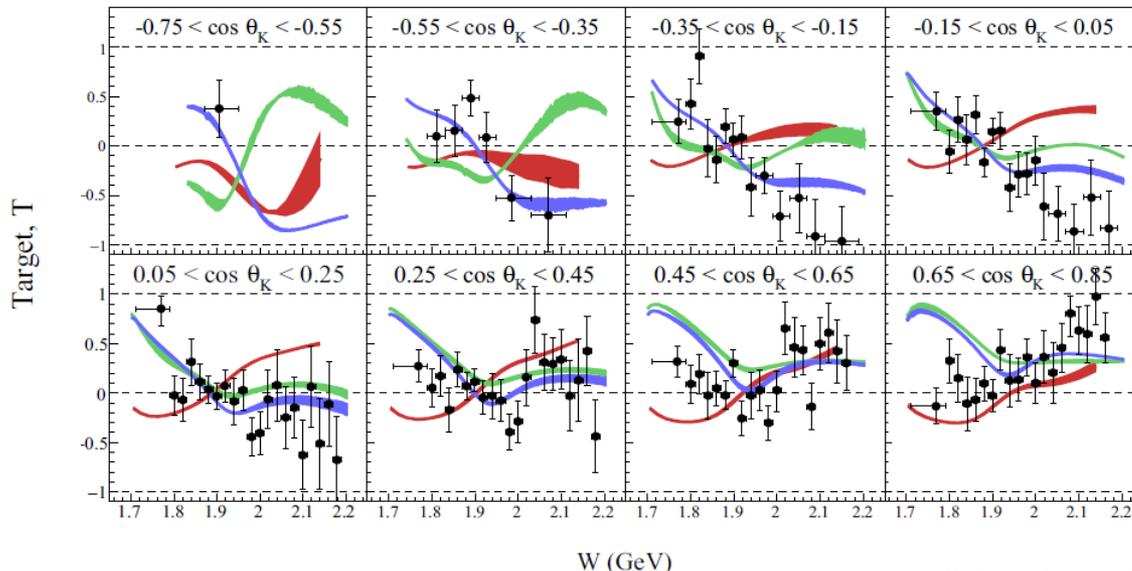


# $\Sigma, T$ for $\gamma p \rightarrow K^+ \Sigma^0$

G8b

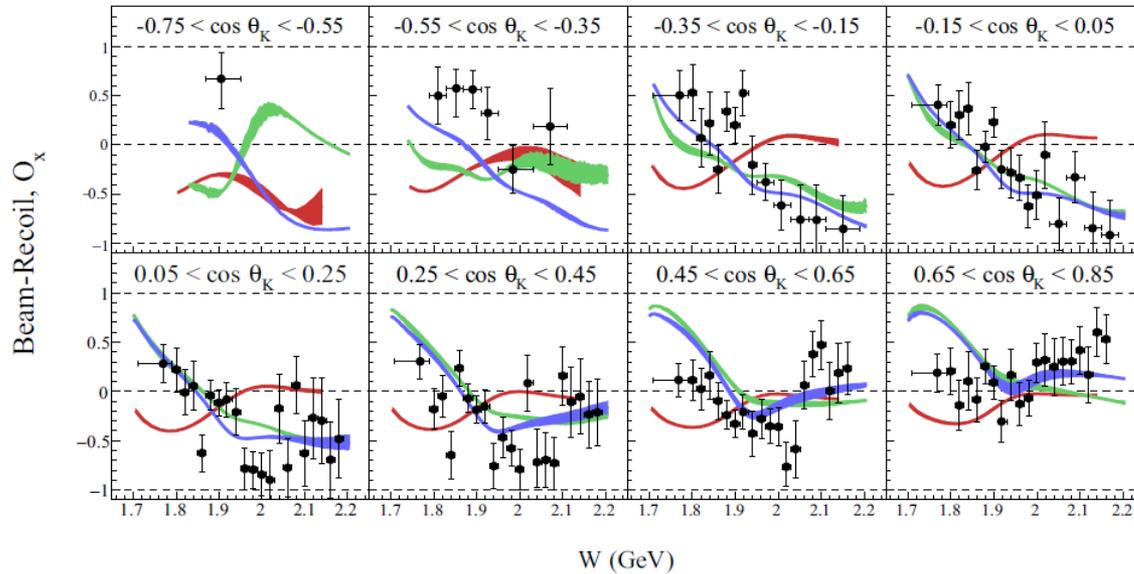


- Blue lines represent fits to Bonn-Gatchina model
- Other lines represent various predictions

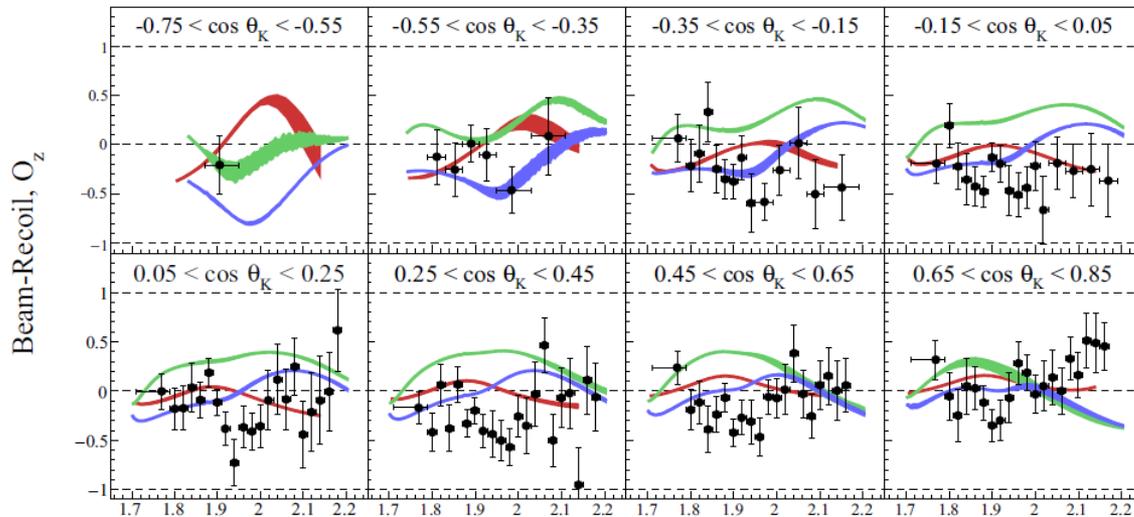


# $O_x, O_z$ for $\gamma p \rightarrow K^+ \Sigma^0$

G8b

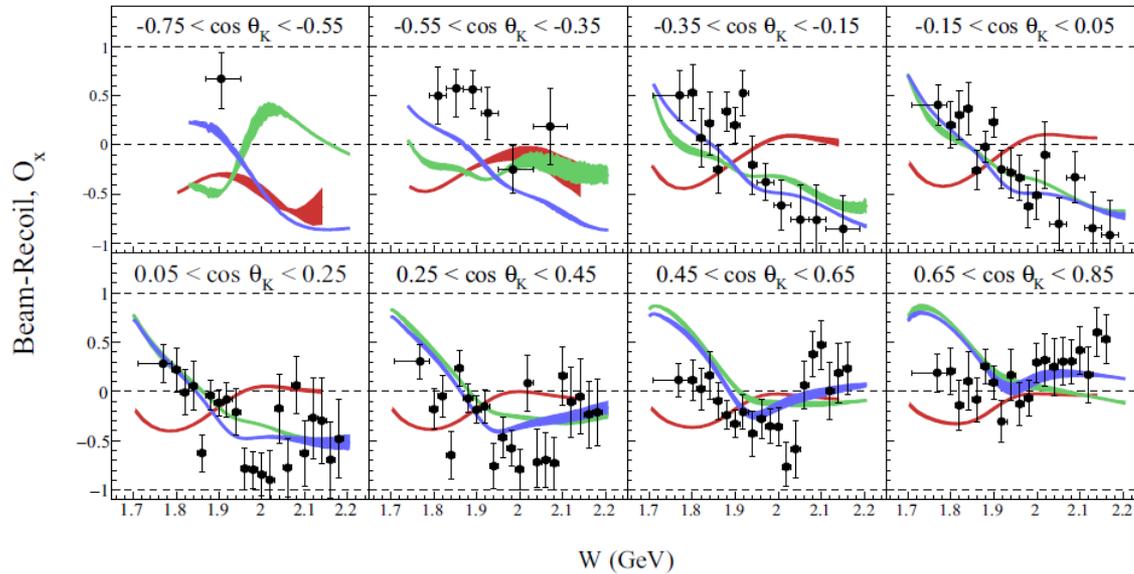


- Blue lines represent fits to Bonn-Gatchina model
- Other lines represent various predictions

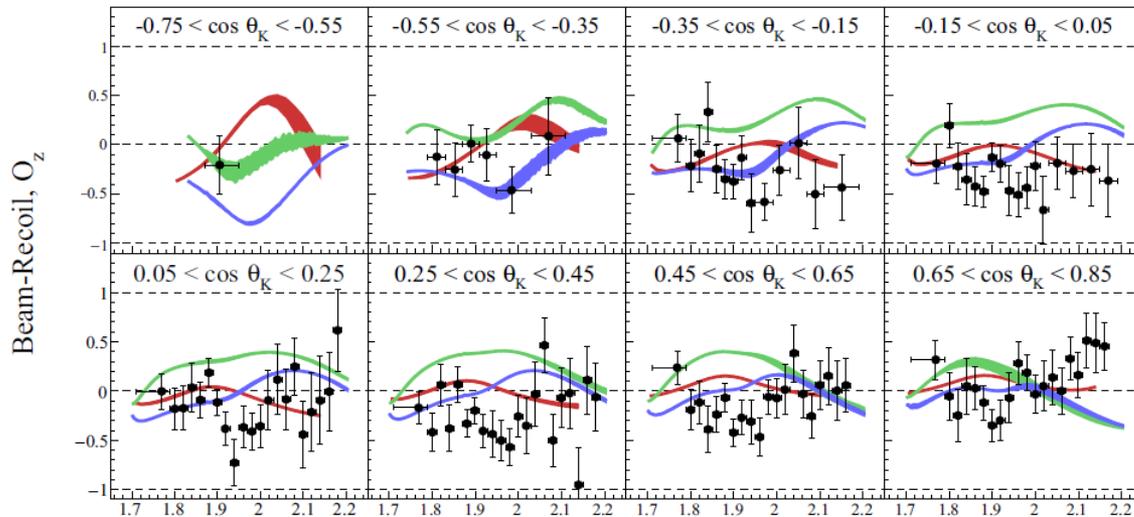


# $O_x, O_z$ for $\gamma p \rightarrow K^+ \Sigma^0$

G8b

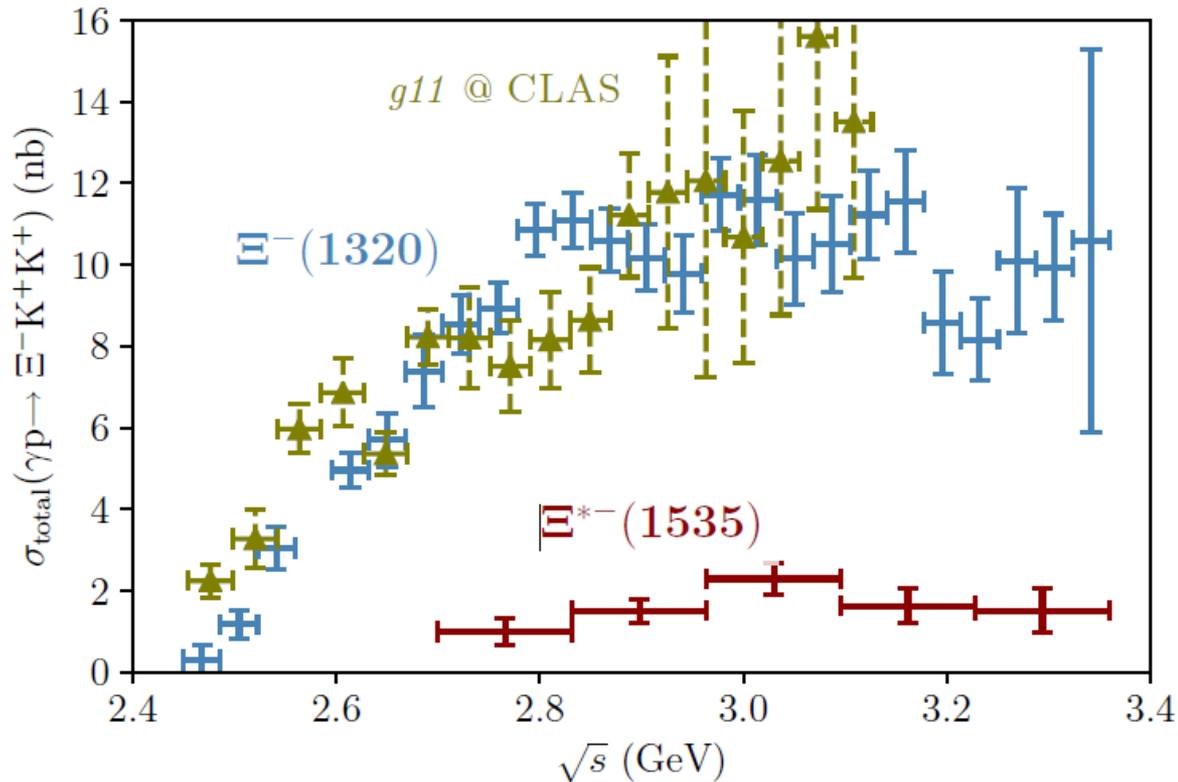


- Indicates some evidence for additional  $N^*(3/2^+)$  and  $N^*(5/2^+)$  resonances of undetermined mass



# $E$ photoproduction

# $\sigma$ for $\gamma p \rightarrow K^+ K^+ \Xi^-$



- All data from CLAS (G11, and G12)
- First total cross sections or photoproduction of these states above  $W=2.8$  GeV

# Observables: $P, C_x, C_z$

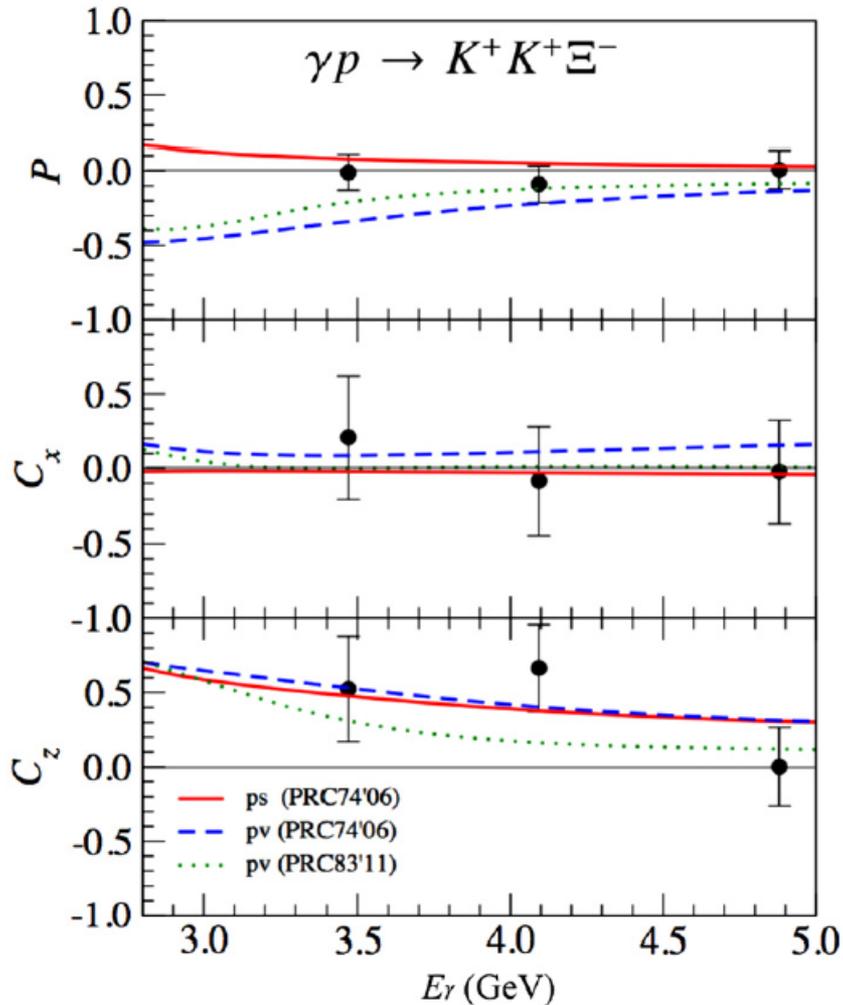


Configuration:

- **Circular photon polarization**
- **Recoil polarization self analyzed**
- No target polarization

Photon		Target			Recoil			Target + Recoil			
	–	–	–	–	$x'$	$y'$	$z'$	$x'$	$x'$	$z'$	$z'$
	–	$x$	$y$	$z$	–	–	–	$x$	$z$	$x$	$z$
unpolarized	$\sigma_0$	0	$T$	0	0	$P$	0	$T_{x'}$	$-L_{x'}$	$T_{z'}$	$L_{z'}$
linear pol.	$-\Sigma$	$H$	$(-P)$	$-G$	$O_{x'}$	$(-T)$	$O_{z'}$	$(-L_{z'})$	$(T_{z'})$	$(-L_{x'})$	$(-T_{x'})$
→ circular pol.	0	$F$	0	$-E$	$-C_{x'}$	0	$-C_{z'}$	0	0	0	0

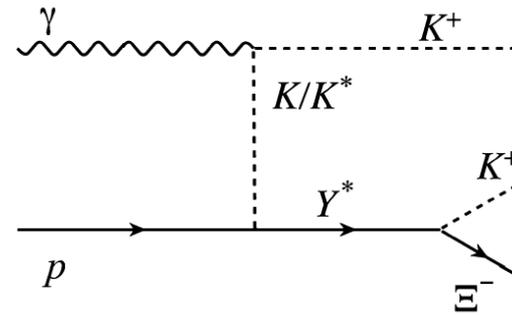
# $P, C_x, C_z$ for $\gamma p \rightarrow K^+ K^+ \Xi^-$



- First-time measurement

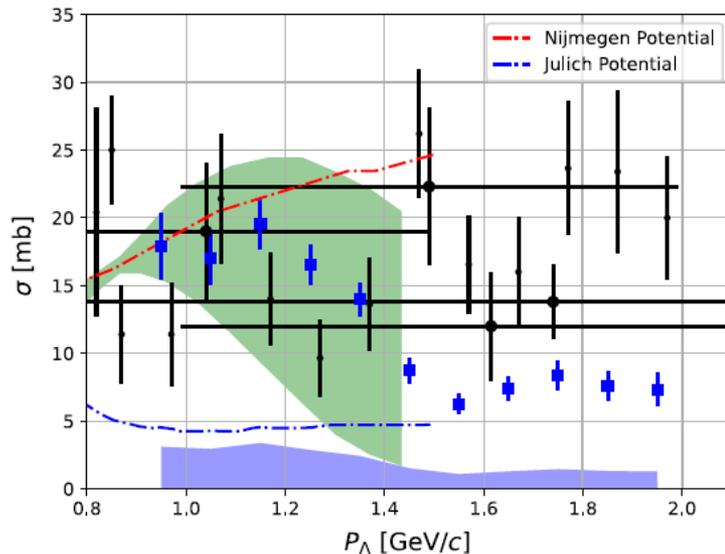
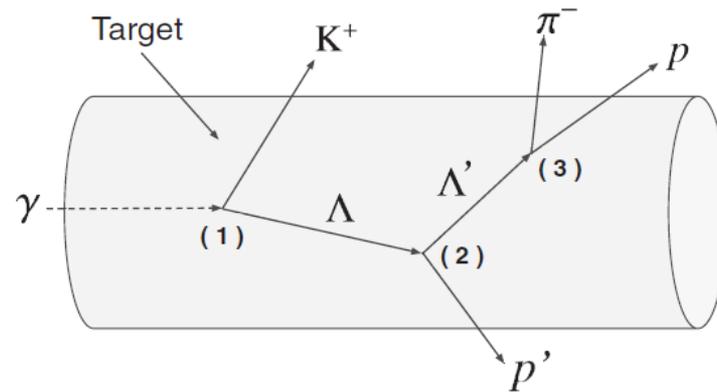
- Coupling:

- ps = pseudoscalar
- pv = pseudovector



- Green dotted includes  $\Sigma(2030)$  contribution

# *$p\Lambda$ elastic scattering: $p\Lambda \rightarrow p\Lambda$*



- **Black** circles: previous world data (bubble chambers)
- **Blue** squares: CLAS results
- Momentum range important to neutron star physics

# Status of meson photoproduction

	$\sigma$	$\Sigma$	$T$	$P$	$E$	$F$	$G$	$H$	$T_x$	$T_z$	$L_x$	$L_z$	$O_x$	$O_z$	$C_x$	$C_z$
Proton target																
$\rho\pi^0$	✓	✓	✓	✓	✓	✓	✓	✓								
$n\pi^+$	✓	✓	✓	✓	✓	✓	✓	✓								
$\rho\eta$	✓	✓	✓	✓	✓	✓	✓	✓								
$\rho\eta'$	✓	✓	✓	✓	✓	✓	✓	✓								
$\rho\omega$	✓	✓	✓	✓	✓	✓	✓	✓								
$K^+\Lambda$	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓
$K^+\Sigma^0$	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓
$K^0\Sigma^+$	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓
"Neutron" target																
$\rho\pi^-$	✓	✓	✓	✓	✓	✓	✓	✓								
$K^+\Sigma^-$	✓	✓	✓	✓	✓	✓	✓	✓								
$K^0\Lambda$	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓
$K^0\Sigma^0$	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓	✓

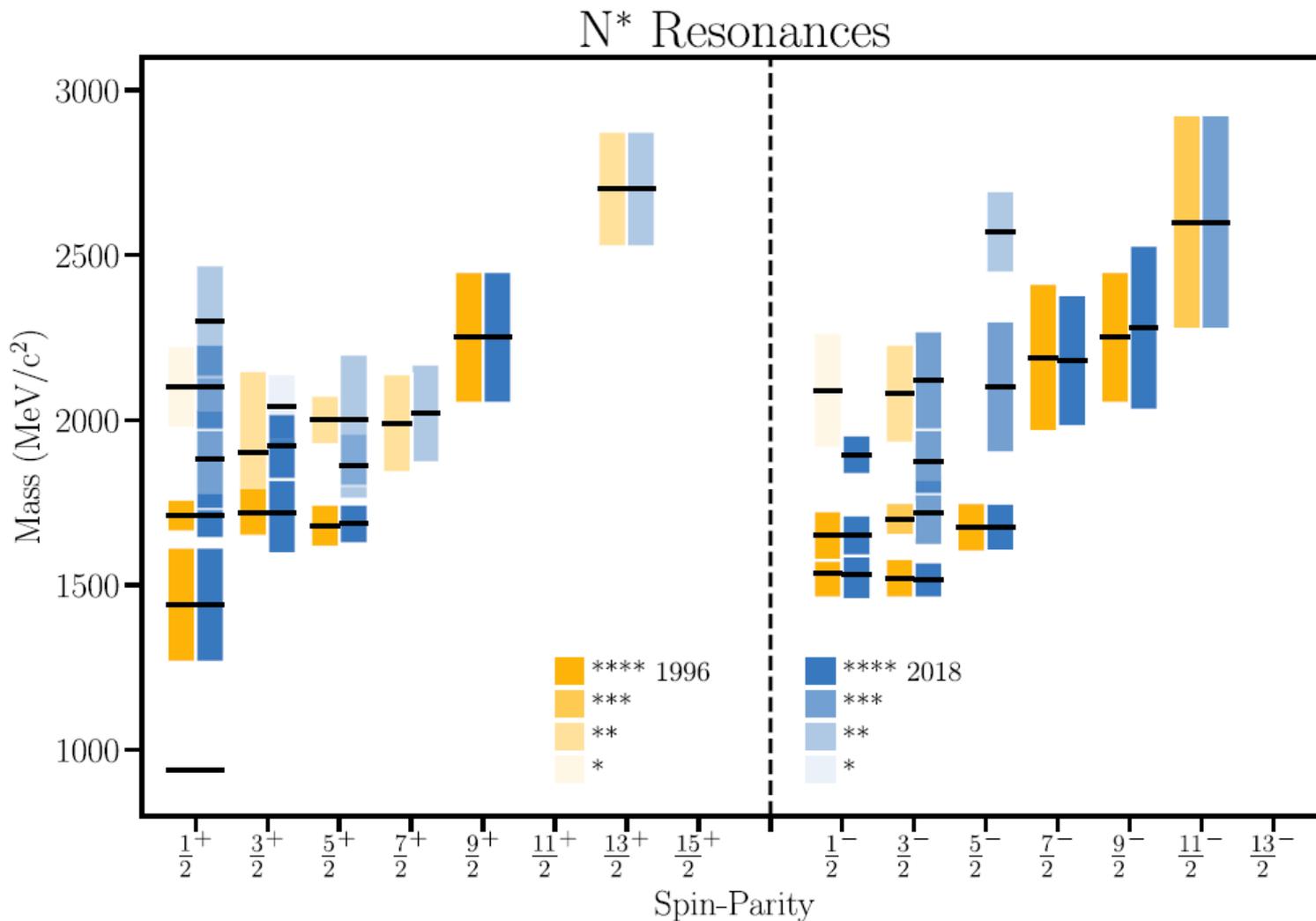
Not shown in table:

- $\pi\pi$  photoproduction observables or
- $E$  states
- $p\Lambda$  scattering

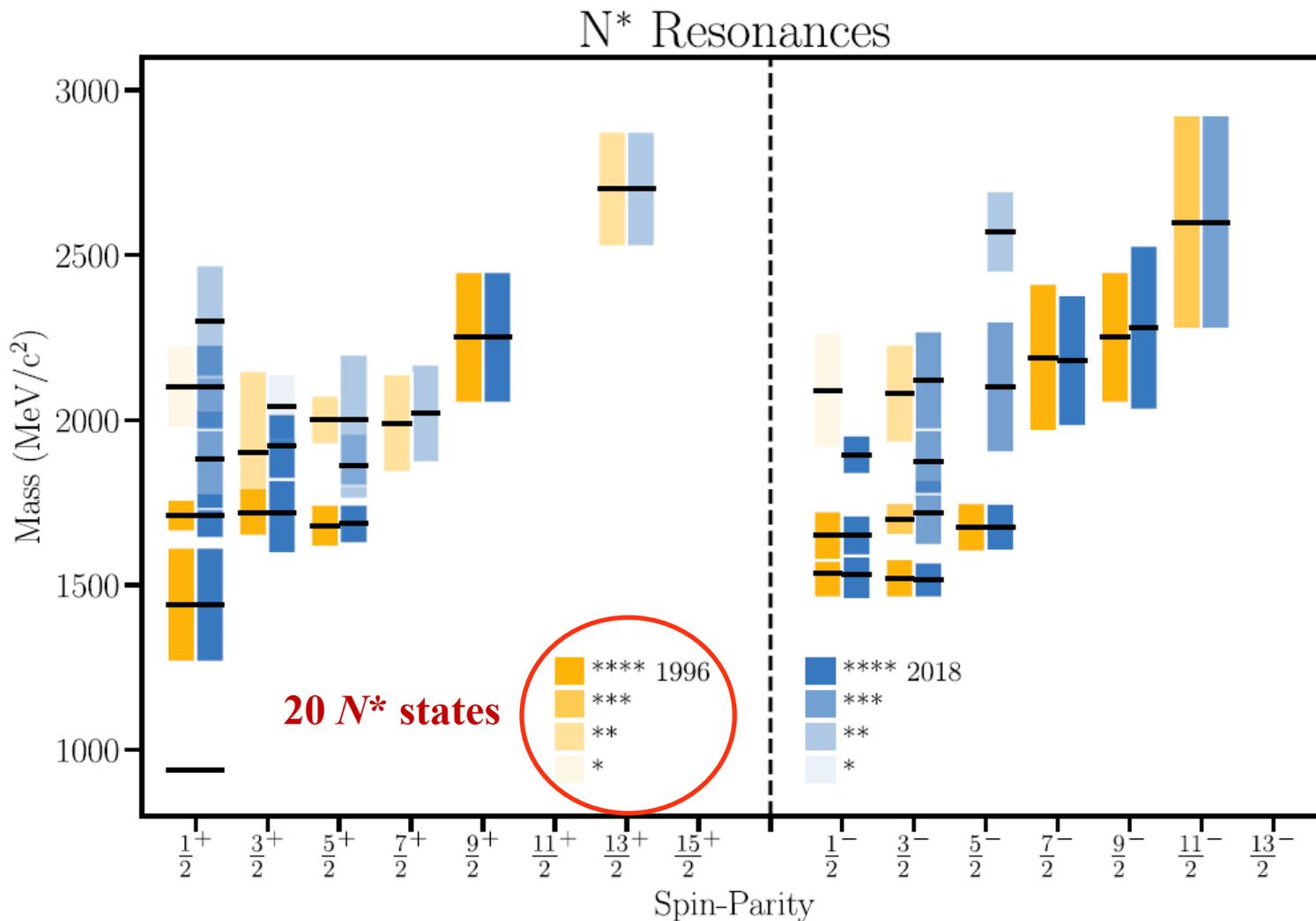
✓ - published    ✓ - acquired



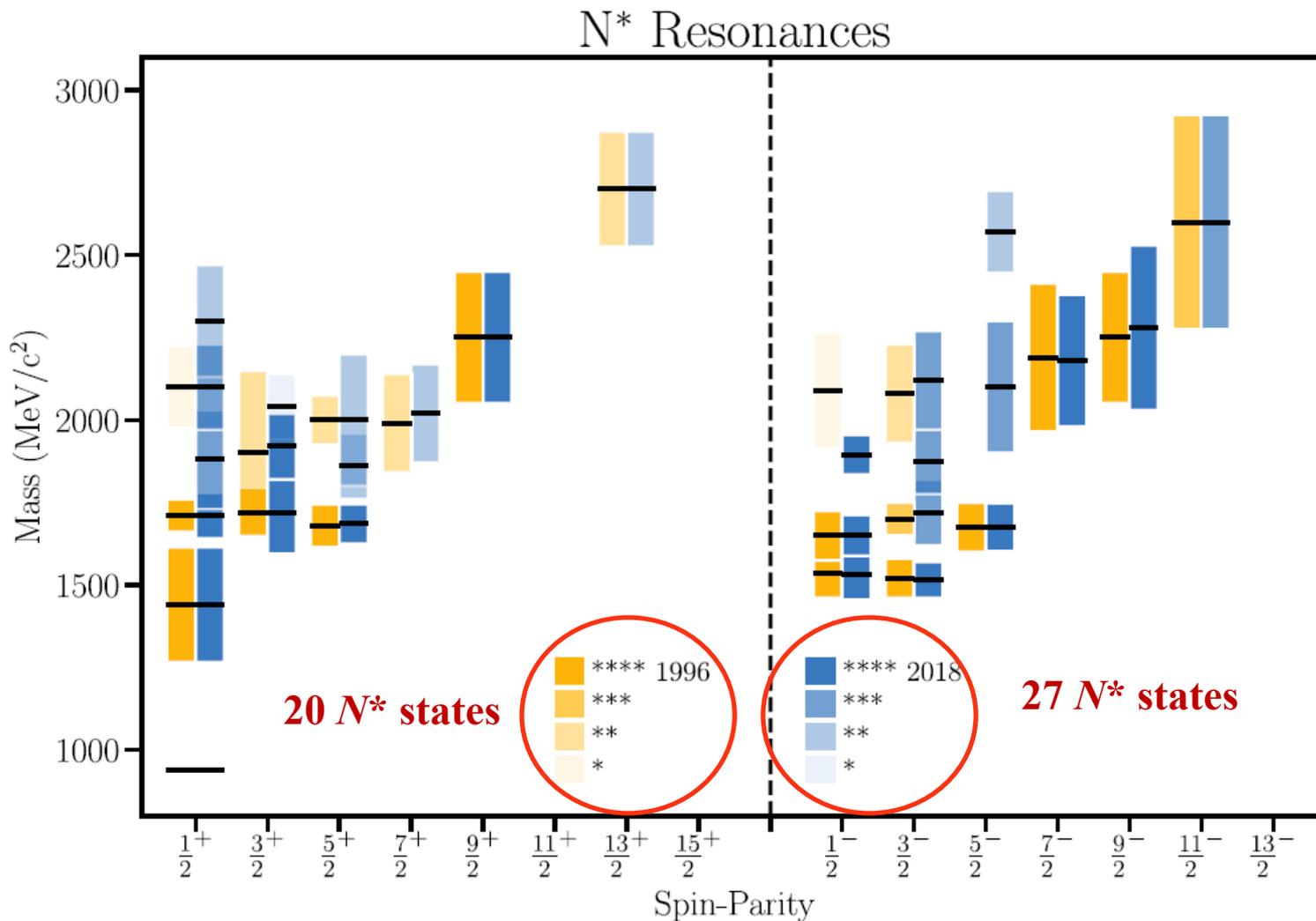
# Changes to PDG from 1996 to 2018



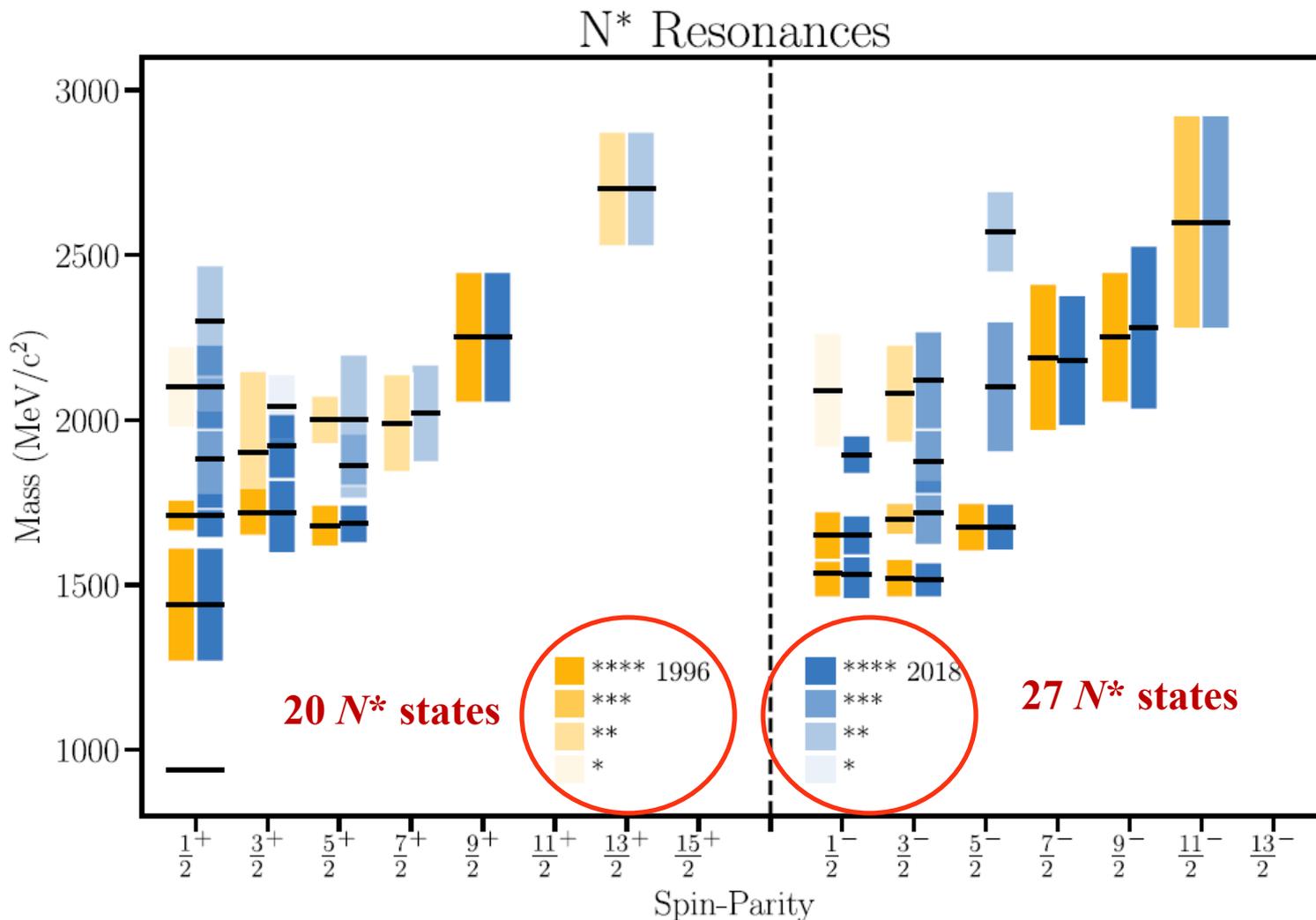
# Changes to PDG from 1996 to 2018



# Changes to PDG from 1996 to 2018

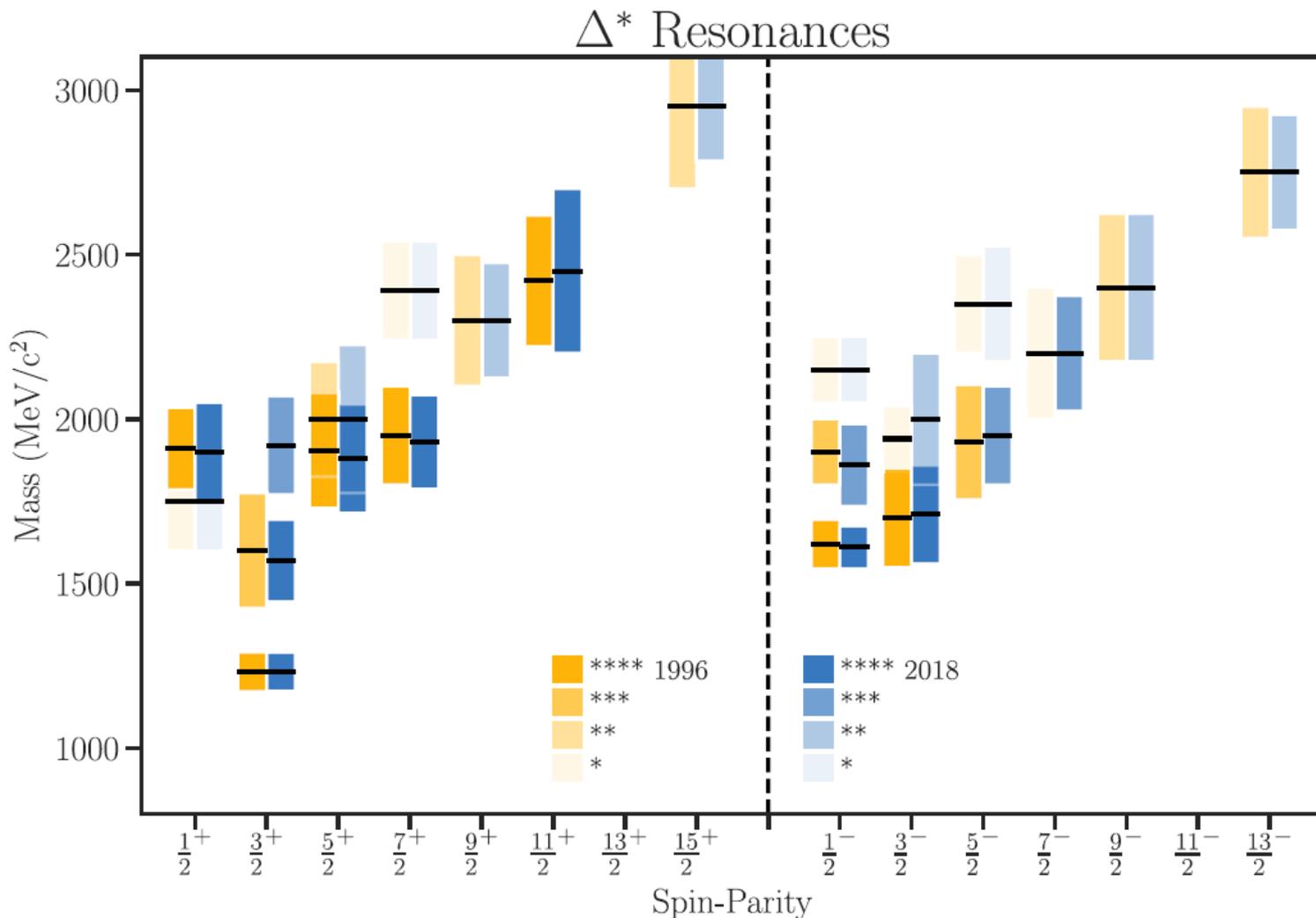


# Changes to PDG from 1996 to 2018



Along with additional new states, “old” states have been measured better and PDG properties have changed

# Changes to PDG from 1996 to 2018



States have been measured better and PDG properties have changed





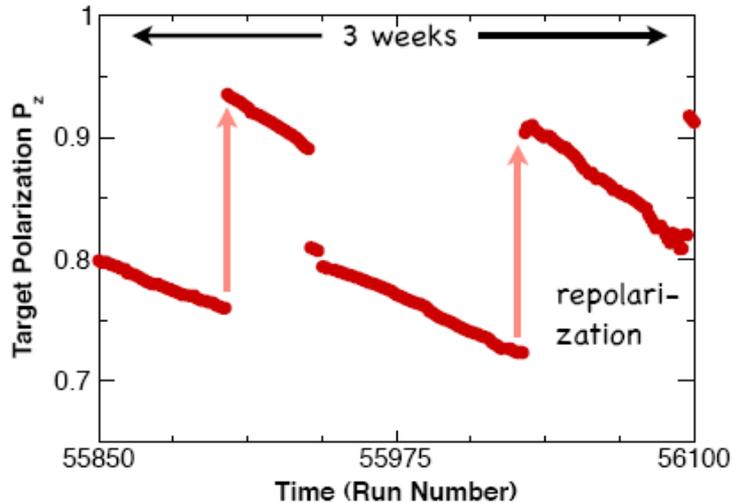




# Frost target

- Brute force polarization requires large magnet
- Instead use “trick” (Dynamic Nuclear Polarization):
  - Dope butanol with paramagnetic radical TEMPO
  - Polarize unpaired TEMPO electrons to 99.999% with  $B = 5 \text{ T}$  and  $T = 0.3 \text{ K}$
  - Transfer electron polarization to free protons with microwaves at  $\sim 140 \text{ GHz}$
  - Remove microwaves
  - Cool to  $T = 3 \text{ mK}$  and use  $B=0.5\text{T}$  holding field
  - Put target in CLAS and run experiment

# Performance: target polarization



- Frozen spin butanol ( $C_4H_9OH$ )
- $P_z \approx 80\%$
- Target depolarization:  $\tau \approx 100$  days

- For g9a (longitudinal orientation) 10% of allocated time was used polarizing target

- For g9b (transverse orientation) 5% of allocated time was used polarizing target

# Frost target

## Brute Force Polarization

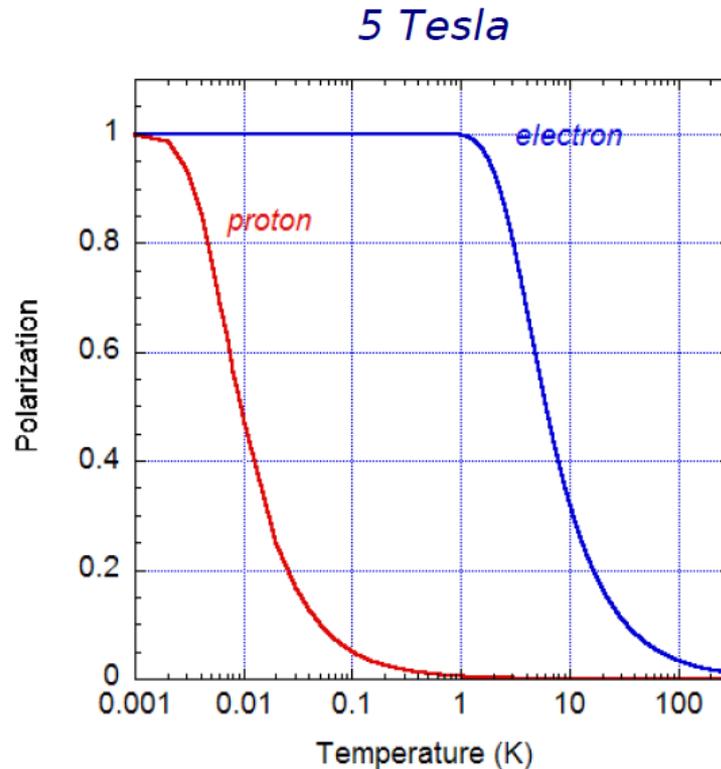
$$P = \tanh\left(\frac{\vec{\mu} \cdot \vec{B}}{kT}\right) \longrightarrow \begin{array}{l} \text{maximize } B, \\ \text{minimize } T \end{array}$$

### Disadvantages:

1. Requires very large magnet
2. Low temperatures mean low luminosity
3. Polarization can take a very long time

We need a trick!

Slide from Chris Keith



# Frost target

## The Trick -- Dynamic Nuclear Polarization

Use brute force to polarize free electrons in the target material. Use microwaves to “transfer” this polarization to nuclei. Mutual electron-nucleus spin flips re-arrange the nuclear Zeeman populations to favor one spin state over the other.

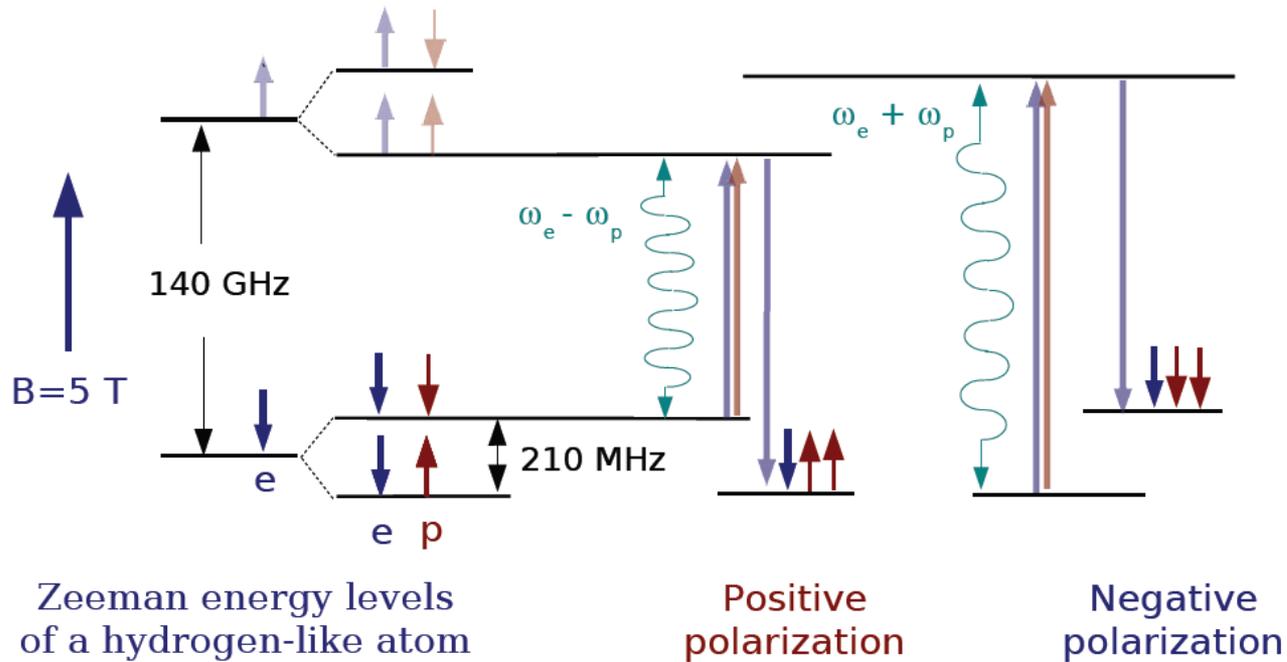
For best results, DNP is performed at  $B/T$  conditions where electron  $t_1$  is short (ms) and nuclear  $t_1$  is long (minutes)

JLab:  $B = 5 \text{ Tesla}$   
 $T = 1 \text{ Kelvin}$

Slide from Chris Keith

# Frost target

## The Resolved Solid Effect



Slide from Chris Keith

# Frost target

## Materials for DNP Targets

- Choice of material dictated by 4 factors:

1. Maximum polarization
2. Resistance to ionizing radiation
3. Presence of unpolarized nuclei
4. Presence of unwanted, polarized nuclei

—————> quality factor,  $f \equiv \frac{\vec{N}}{N_{total}}$

- Free electrons must be embedded into target material:

1. Chemical doping with paramagnetic radicals
2. Paramagnetic radicals created by ionizing radiation

- Typically 1 free electron can “service”  $\sim 10^3$  free protons

Slide from Chris Keith

## Materials for DNP Targets, examples

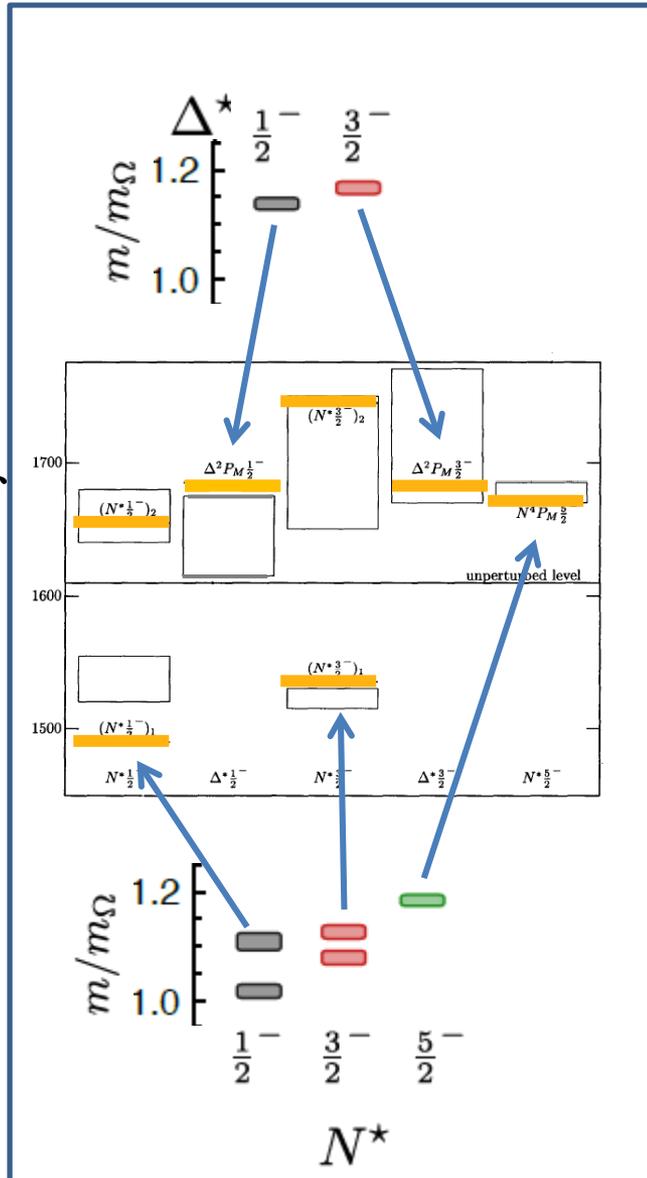
Name	Dopant	f	Rad. Resistance
Polyethelyne, C <sub>2</sub> H <sub>4</sub>	chemical	0.12	low
Polystyrene, C <sub>8</sub> H <sub>8</sub>	chemical	0.07	low
Propandiol, C <sub>3</sub> H <sub>6</sub> (OH) <sub>2</sub>	chemical	0.11	moderate
Butanol, C <sub>4</sub> H <sub>9</sub> OH	chemical	0.13	moderate
Ammonia, <sup>15</sup> NH <sub>3</sub>	radiation	0.17	high
Lithium Hydride, <sup>7</sup> LiH	radiation	0.12	very high

Slide from Chris Keith

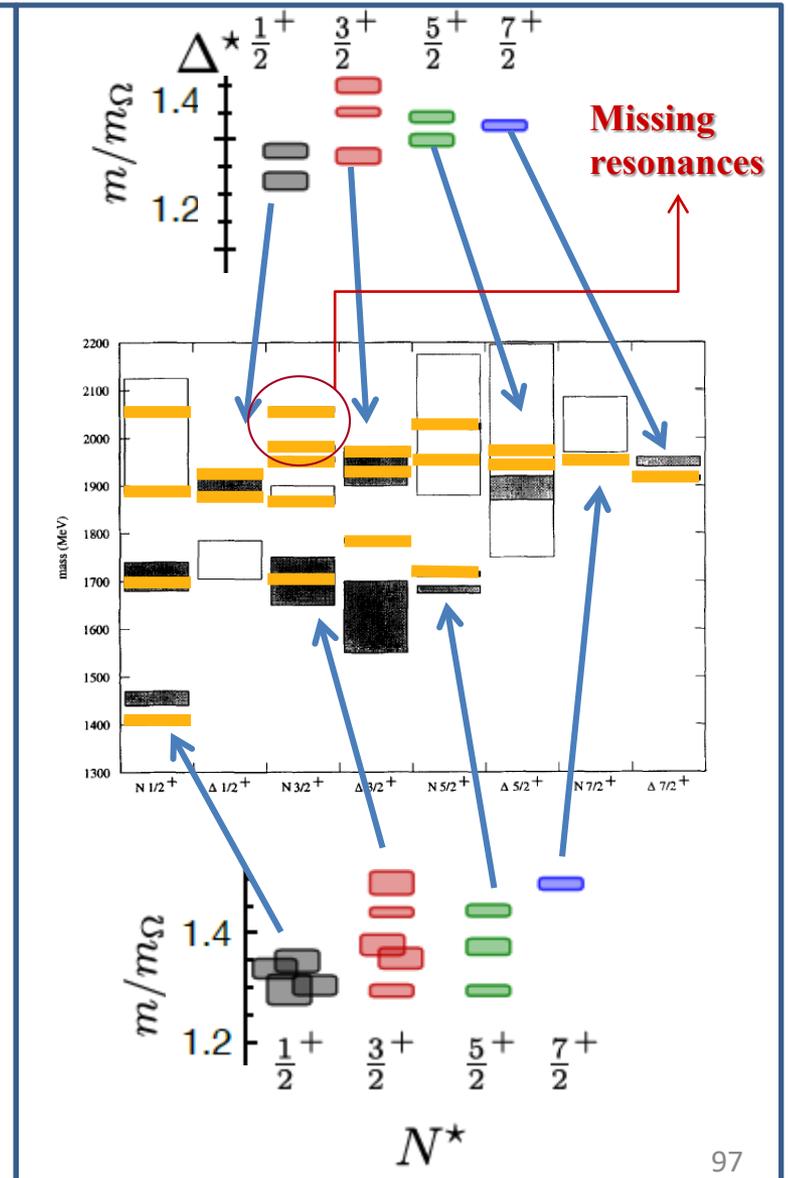
# Low-lying Resonance States

Lattice QCD is consistent with non-relativistic quark model for number of low-lying states

## Negative parity



## Positive parity





The differential cross section for  $\gamma p \rightarrow p \pi^+ \pi^-$

(without measuring the polarization of the recoiling nucleon)

$$\frac{d\sigma}{d\mathbf{x}_i} = \sigma_0 \left\{ (1 + \vec{\Lambda}_i \cdot \vec{\mathbf{P}}) + \delta_{\odot} (\mathbf{I}^{\odot} + \vec{\Lambda}_i \cdot \vec{\mathbf{P}}^{\odot}) \right.$$

Next slides

$$\left. + \delta_I [\sin 2\beta (\mathbf{I}^s + \vec{\Lambda}_i \cdot \vec{\mathbf{P}}^s) + \cos 2\beta (\mathbf{I}^c + \vec{\Lambda}_i \cdot \vec{\mathbf{P}}^c)] \right\}$$

Circular beam and longitudinal target:

$$\delta_I = \Lambda_x = \Lambda_y = 0$$

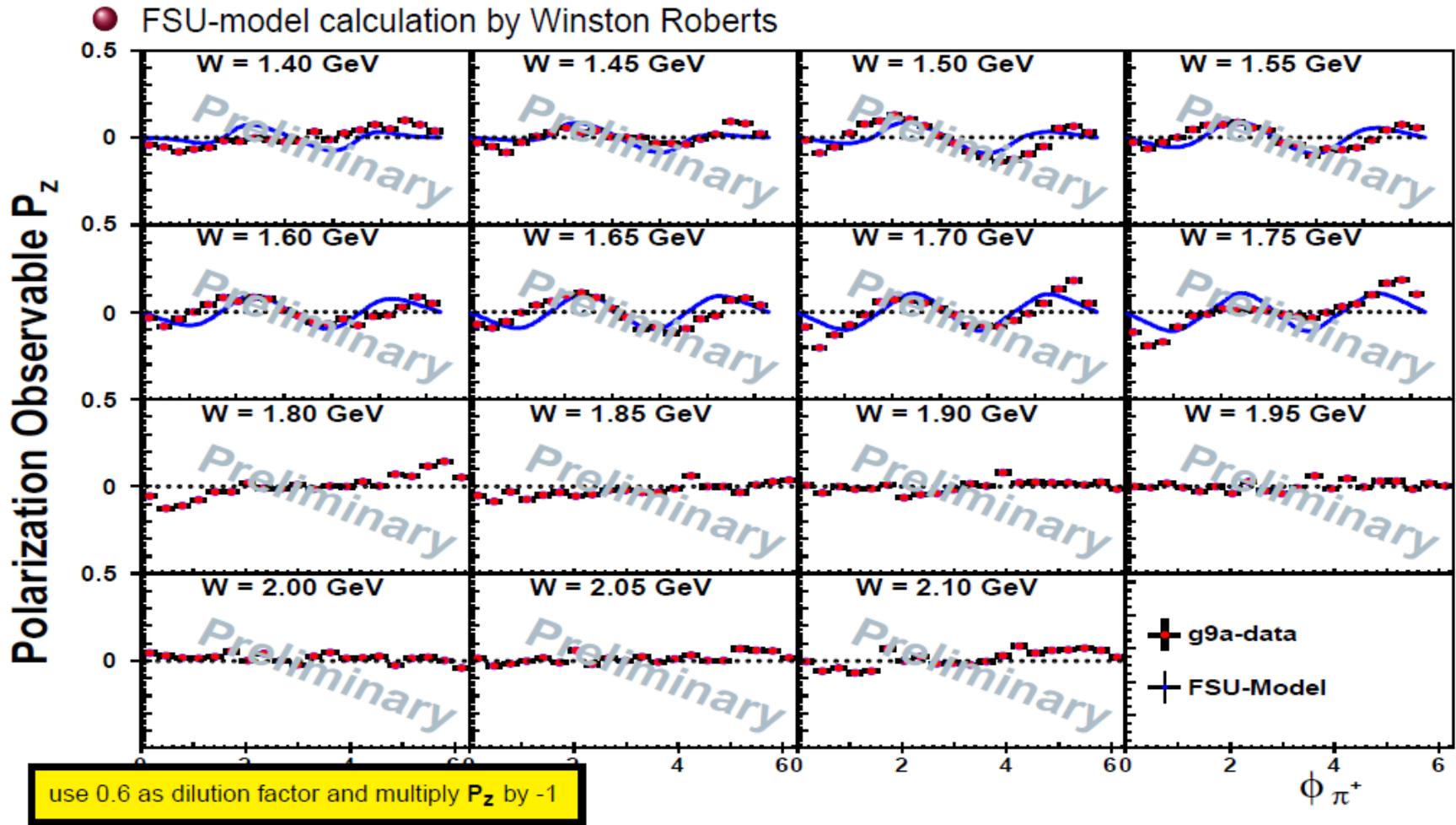
G9a: FROST

- $\sigma_0$ : The unpolarized cross section
- $\beta$ : The angle between the direction of polarization and the x-axis
- $\delta_{\odot, I}$ : The degree of polarization of the photon beam  $\Rightarrow \delta_{\odot}$ , and  $\delta_I$
- $\vec{\Lambda}_i$ : The polarization of the initial nucleon  $\Rightarrow (\Lambda_x, \Lambda_y, \Lambda_z)$
- $\mathbf{I}^{\odot, s, c}$ : The observable arising from use of polarized photons  $\Rightarrow \mathbf{I}^{\odot}, \mathbf{I}^s, \mathbf{I}^c$
- $\vec{\mathbf{P}}$ : The polarization observable  $\Rightarrow (\mathbf{P}_x, \mathbf{P}_y, \mathbf{P}_z) (\mathbf{P}_x^{\odot}, \mathbf{P}_y^{\odot}, \mathbf{P}_z^{\odot}) (\mathbf{P}_x^s, \mathbf{P}_y^s, \mathbf{P}_z^s) (\mathbf{P}_x^c, \mathbf{P}_y^c, \mathbf{P}_z^c)$

15 Observables

# $P_z$ for $p \pi^+ \pi^-$

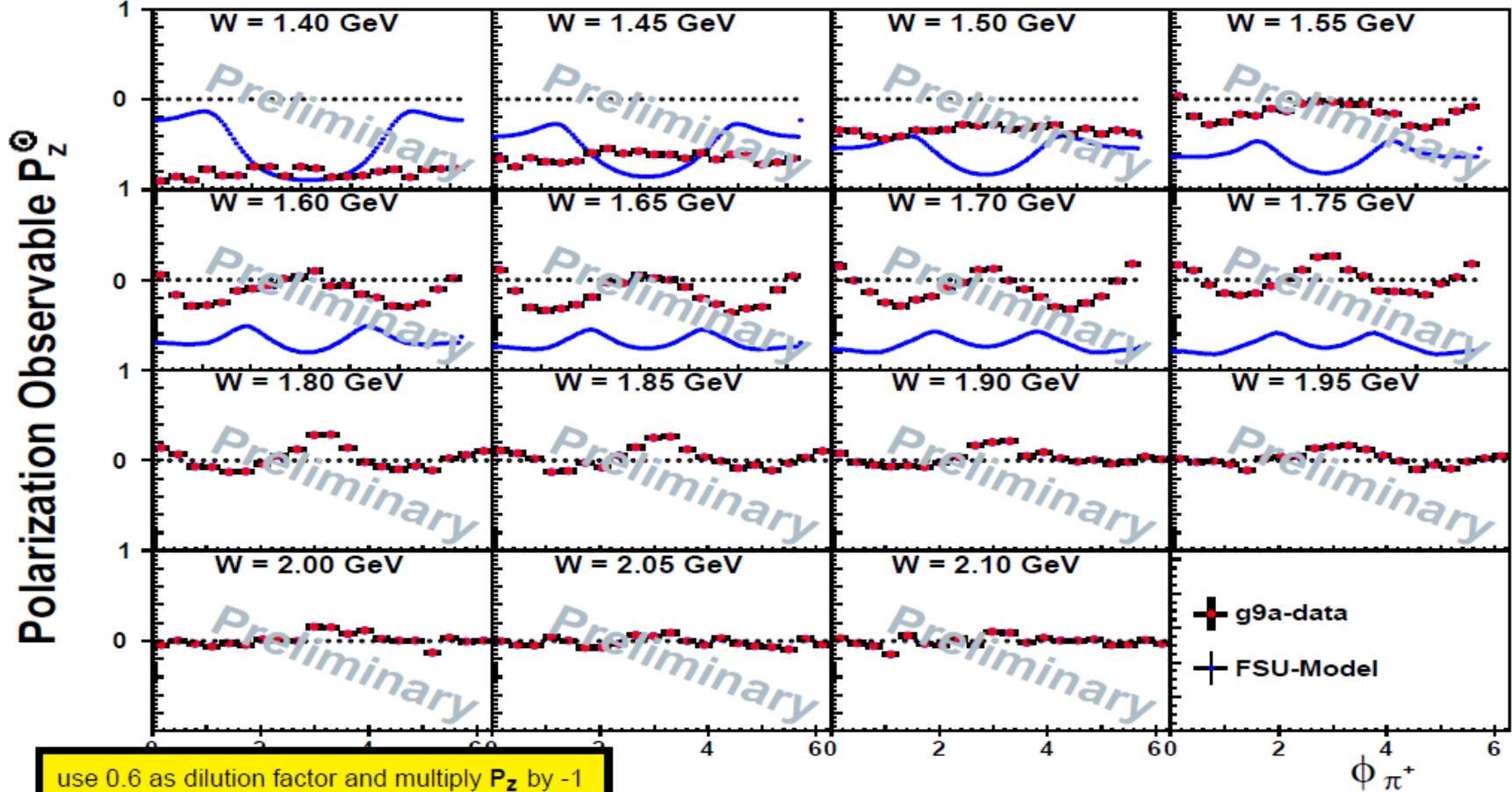
G9a: FROST



# $P^\ominus$ for $p \pi^+ \pi^-$

G9a: FROST

● FSU-model calculation by Winston Roberts



# Observable

Configuration:

- **Linear photon polarization**
- **Longitudinal Target polarization**
- No recoil polarization

Experiment:

- g9a: FROST

Photon		Target			Recoil			Target + Recoil			
	–	–	–	↓ z	x'	y'	z'	x'	x'	z'	z'
	–	x	y	z	–	–	–	x	z	x	z
unpolarized	$\sigma_0$	0	T	0	0	P	0	$T_{x'}$	$-L_{x'}$	$T_{z'}$	$L_{z'}$
linear pol.	$-\Sigma$	H	(-P)	<b>-G</b>	$O_{x'}$	(-T)	$O_{z'}$	$(-L_{z'})$	$(T_{z'})$	$(-L_{x'})$	$(-T_{x'})$
circular pol.	0	F	0	-E	$-C_{x'}$	0	$-C_{z'}$	0	0	0	0



# Isospin photo-couplings for $\gamma p \rightarrow n\pi^+$ and $\gamma n \rightarrow p\pi$

$$\begin{array}{l}
 \gamma p \begin{cases} \rightarrow \text{Iso-singlet} & A^0 |I=0, I_3=0\rangle |I=\frac{1}{2}, I_3=\frac{1}{2}\rangle = A^0 |I=\frac{1}{2}, I_3=\frac{1}{2}\rangle \\
 \rightarrow \text{Iso-vector} & A^1 |I=1, I_3=0\rangle |I=\frac{1}{2}, I_3=\frac{1}{2}\rangle = A^1 \left[ \sqrt{\frac{2}{3}} |I=\frac{3}{2}, I_3=\frac{1}{2}\rangle \ominus \sqrt{\frac{1}{3}} |I=\frac{1}{2}, I_3=\frac{1}{2}\rangle \right] \end{cases} \\
 \gamma n \begin{cases} \rightarrow \text{Iso-singlet} & A^0 |I=0, I_3=0\rangle |I=\frac{1}{2}, I_3=-\frac{1}{2}\rangle = A^0 |I=\frac{1}{2}, I_3=-\frac{1}{2}\rangle \\
 \rightarrow \text{Iso-vector} & A^1 |I=1, I_3=0\rangle |I=\frac{1}{2}, I_3=-\frac{1}{2}\rangle = A^1 \left[ \sqrt{\frac{2}{3}} |I=\frac{3}{2}, I_3=-\frac{1}{2}\rangle \oplus \sqrt{\frac{1}{3}} |I=\frac{1}{2}, I_3=-\frac{1}{2}\rangle \right] \end{cases}
 \end{array}$$

$$\gamma p \rightarrow n\pi^+: \quad \oplus \sqrt{\frac{2}{3}} \left[ A^0 \ominus \sqrt{\frac{1}{3}} A^1 \right] N^* + \frac{\sqrt{2}}{3} A^1 \Delta^*$$

$$\gamma n \rightarrow p\pi: \quad \oplus \sqrt{\frac{2}{3}} \left[ A^0 \oplus \sqrt{\frac{1}{3}} A^1 \right] N^* + \frac{\sqrt{2}}{3} A^1 \Delta^*$$

- Using both proton and neutron targets allows decomposition of iso-singlet and iso-vector photo-couplings
- The sings in  $\ominus$   $\oplus$  will give interference terms

