Pair Spectrometer Hodoscope for Hall D at Jefferson Lab[☆]

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Abstract

We present the design of the pair spectrometer hodoscope fabricated at Jefferson Lab and installed in the experimental Hall D. The hodoscope consists of thin scintillator tiles; the light from each tile is collected using wave-length shifting fibers and detected using a Hamamatsu silicon photomultiplier. Light collection was measured using relativistic electrons produced in the tagger area of the experimental Hall B. *Keywords:* Scintillator detector, Pair spectrometer, Silicon photomultiplier

1 1. Introduction

The new detector GlueX has been constructed in the experimental Hall D at Thomas Jefferson National Accelerator Facility. The main physics goal of the GlueX experiment [1] is to search for hybrid mesons with exotic quantum numbers using a beam of linearly polarized photons incident on a liquid hydrogen target. The linearly polarized photon beam in Hall D will be produced via the coherent bremsstrahlung process by 12 GeV electrons in a thin diamond radiator. Coherent bremsstrahlung has been successfully used to produce linearly polarized photon beams at various experimental facilities [2, 3]. Coherent radiation from a diamond crystal lattice results in sharp monochromatic peaks in the photon energy spectrum. The peak energy depends on the relative orientation of the electron beam direction and the crystal lattice, and can

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be adjusted in the experiment by rotating the diamond. The GlueX photon energy re-12 gion of interest corresponding to the main coherent peak is between 8.4 GeV and 9.1 13 GeV. The fraction of linearly polarized photons can be increased by passing a photon beam through a 3.4 mm diameter collimator situated about 75 m downstream of 15 the radiator. The collimator preferentially filters out photons produced via incoherent 16 bremsstrahlung, which have a substantially larger mean emission angle. The colli-17 mated photons propagate toward the GlueX target. The energy of a beam photon can 18 be inferred by analyzing the momentum of the electron after it has radiated the pho-19 ton, given that the primary beam energy is known. This so-called "tagged" electron is 20 deflected in a 6 m long dipole magnet operated at a field of 1.8 T and is subsequently 21 registered in the tagging microscope or broad-band hodoscope scintillator counters. 22

One of the key components of the Hall D photon beam line is the magnetic pair 23 spectrometer, which is installed after the photon collimator in front of the GlueX de-24 tector. The spectrometer will reconstruct the energy of a beam photon by detecting 25 the e^{\pm} pair produced by the photon in a thin converter. The main purpose of the spec-26 trometer is to measure the spectrum of the collimated photon beam and determine the 27 fraction of linearly polarized photons in the coherent peak energy region. It will also 28 monitor the photon beam flux and can be used for the energy calibration of the tagging 29 hodoscope and microscope detectors. The description of the Hall D pair spectrometer 30 will be presented in Section 2. The design of the pair spectrometer high-granularity 31 hodoscope and light collection measurements from hodoscope scintillator tiles will be 32 described in Section 3 and Section 4, respectively. 33

34 2. Pair Spectrometer

Layout of the Hall D pair spectrometer is presented in Fig. 1. Electron-positron pairs are created by beam photons inside a thin converter with a typical thickness ranging between $10^{-4} - 10^{-2}$ radiation lengths. The choice of the converter thickness depends on the photon beam flux. The maximum flux for GlueX physics runs is expected to be $5 \cdot 10^7 \gamma$ /sec in the coherent peak energy region. Three converters with different thicknesses are installed in a movable fork, that can insert one of them into the photon



Figure 1: Schematic top view of the Hall D pair spectrometer. PS and PSC denote the hodoscope and coarse scintillator counters, respectively.

beam. Produced leptons are deflected in a 18D36 dipole magnet with an effective field 41 length of about 0.94 m. The magnet was brought from Brookhaven National Labora-42 tory and was modified at Jefferson Lab by reducing the pole gap from 6 inches to 3 43 inches. The magnet is operated at a nominal field of 1.8 T. A 1.5 meter long vacuum 44 chamber is installed after the magnet. Electrons and positrons are registered in two 45 layers of scintillator detectors: a high-granularity hodoscope and a set of coarse coun-46 ters, denoted in Fig. 1 as PS and PSC, respectively. The detectors are organized into 47 two arms positioned symmetrically with respect to the photon beam line. Each detector 48 arm covers a momentum range of e^{\pm} between 3.0 GeV/c to 6.2 GeV/c, corresponding 49 to reconstructed photon energies between 6 GeV and 12.4 GeV. Relatively large accep-50 tance of the hodoscope allows one to reconstruct photons with energies in the coherent 51 peak energy region and also in the range near the beam end-point energy of 12 GeV. 52 This can be used for the energy calibration of the hodoscope detectors. 53

The high-resolution hodoscope is used for precise measurements of e^{\pm} momenta. The hodoscope consists of a set of thin scintillator tiles, stacked together. Momentum of the reconstructed lepton is related to the *x*-coordinate of each tile, where *x* is an axis perpendicular to the beam line. Hodoscope tiles will be calibrated in units of lepton energy, the energy scale will be used hereafter in the text. The hodoscope momentum resolution was studied using a detailed Geant detector simulation. The hodoscope energy resolution varies between 12 MeV and 20 MeV for 6 GeV and 12 GeV reconstructed photons, respectively. Light from each scintillator tile is detected using a Hamamatsu silicon photomultiplier (SiPM). The detailed description of the hodoscope will be presented in Section 3.

Sixteen coarse scintillator counters, eight in each detector arm, are positioned about 64 40 cm behind the hodoscope. The counters are 4.4 cm wide in the direction perpen-65 dicular to the lepton trajectory, 2 cm thick, and 6 cm in height. Hamamatsu R6427-01 66 PMTs are used to detect the scintillation light. The counters are used to produce a pair 67 spectrometer trigger by requiring a coincidence of hits in the two detector arms. They 68 also help to reduce background originating from interactions of e^{\pm} inside the magnet 69 pole edges to the level below 1% by constraining e^{\pm} trajectories. Counter rates depend 70 on the converter thickness and photon beam flux. The maximum rate will not exceed 71 10 kHz per PSC counter during GlueX operation. 72

Signals from both the hodoscope and coarse counters are digitized using a twelvebit multi-channel flash ADC operated at a sampling rate of 250 MHz [5]. The PSC counters are also instrumented with TDCs, with a channel width of 60 ps [6]. The time resolution is expected to be better than 200 ps. This resolution will allow one to distinguish the electron beam bunch where the bremsstrahlung photon is emitted and therefore relate hits from the pair spectrometer and tagging detectors originating from the same event.

3. Pair Spectrometer Hodoscope

81 3.1. Design and Fabrication

The pair spectrometer hodoscope consists of 145 rectangular tiles made of EJ-212 scintillator [4], stacked together as shown in Fig. 2. The tile height is 3 cm and the length along the electron path in scintillator is 1 cm. Tiles with two different widths, in the direction perpendicular to the electron trajectory, of 2 mm and 1 mm are used. Forty 1 mm wide tiles are instrumented in the detector high-energy region around 6

GeV in order to provide better energy resolution in this energy range. The energy bin 87 size of the hodoscope tiles varies depending on the lepton energy and constitutes about 88 13 MeV and 24 MeV for 3 GeV and 6 GeV leptons, respectively. Tiles are optically 89 isolated using 10 µm aluminized Mylar foil. This reflective foil also covers the bottom 90 of the tile assembly. The deflection angle of leptons in the dipole magnet depends on 91 the lepton momentum; the angular spread is about 5 degrees between 3 GeV and 6.2 92 GeV electrons. To account for the angular dependence, hodoscope tiles are organized 93 into 18 groups. Each group is tilted by about 0.3 degrees. This was done by attaching a 94 shim, $a \sim 50 \mu m$ thick adhesive strip, to one side of the aluminum foil after each group. 95 Light from a tile is collected using two 20 cm long 1 mm x 1 mm square double-96 clad BCF-92 wave-length shifting (WLS) fibers, each of which is glued to the side of 97 the tile using BC 600 Optical Cement. A tile assembly with two WLS optical fibers 98 is shown in Fig. 3. The peak of the emission spectrum for EJ-212 scintillator occurs 99 at the wavelength of 423 nm, which couples well with the absorption spectrum of the 100 BCF-92 fiber. Light is subsequently reemitted inside the fiber in the green range with 10 an emission peak of 492 nm. 102

Collected light is transmitted to the end of the WLS fiber. A pair of fibers from each 103 tile is inserted into a hole in an aluminum mounting plate, as shown on the upper plot 104 of Fig. 2. An electronics circuit board containing 145 photo detectors is attached to 105 the mounting plate. Each photo sensor is coupled to two WLS fibers from a single tile. 106 The light detection is performed using Hamamatsu surface mount S10931-050P silicon 107 photomultipliers with an effective photosensitive area of 3 mm x 3 mm and a pixel size 108 of 50 µm. These sensors have a photon detection efficiency (PDE) larger than 20% at 109 a wavelength of 500 nm and a typical gain of about $7 \cdot 10^5$ [7]. The electronics board 110 with SiPMs is presented in Fig. 4. The SiPMs are arranged in two arrays of 3 x 35 and 111 5 x 8 sensors, which are connected to 2 mm and 1 mm tiles, respectively. SiPMs are 112 optically isolated using a plastic spacer. 113

114 3.2. Electronics

The hodoscope front end electronics consists of an amplifier and a SiPM bias voltage control circuit developed at Jefferson Lab. Signals from the SiPMs are amplified

using the amplifier with a gain of about a factor of 20. The amplifier is based on 117 commercially available devices (operational amplifiers) with 3 GHz bandwidth. Pulse 118 shaping is employed to compensate for the characteristically high SiPM capacitance 119 and package inductance. The impulse response shows rise and fall times of 3 ns and 120 with trans-impedance gain of 1 mV/ μ A. The SiPM operating bias voltage is about 73 121 V. The nominal bias setting is as specified by Hamamatsu (1 V over voltage) and fed 122 to the SiPM through a resistive network employing a thermistor and a linearizing resis-123 tor. The hodoscope control electronics supply individually adjusted voltages to groups 124 of 5 SiPM channels; inside the group the voltage is adjusted among channels using 125 resistors. The thermistor senses the average temperature of closely packed SiPMs in 126 thermal equilibrium via a heat spreader PCB layout, thus forming a well controlled 127 loop. The commercially available bias power supply has very low noise characteris-128 tics and is well regulated to less than 1 mV long term. The supply allows the user to 129 monitor and adjust the levels as needed and if required. The optimal bias setting will 130 be determined based on experimental conditions. Amplified SiPM signals are digitized 131 by flash ADCs positioned in two readout VXS crates. An example of the flash ADC 132 signal pulse obtained from a scintillator tile is shown in Fig. 5. The ADC sampling 133 time is 4 ns. 134

After fabrication, the performance of SiPMs and electronics were tested using a 135 picoseconds laser light pulser¹. The typical rise time of a SiPM pulse produced by a 136 laser pulser is about 6 ns and the corresponding fall time is 45 ns. The pulser trigger 137 output was used to initiate the flash ADC readout. SiPM amplitudes digitized by ADC 138 are integrated in a time window of 60 ns. A typical ADC spectrum is shown in Fig. 6. 139 Peaks on this plot correspond to the number of pixels fired in the SiPM. The position 140 of a single-pixel peak depends on the combined gain of a SiPM and electronics. The 141 single-pixel peak distribution measured for all hodoscope SiPMs is presented in Fig. 7. 142 The solid curve corresponds to the fit of this distribution to a Gaussian function. The 143 relative gain variation is found to be smaller than 1.5%. 144

¹Hamamatsu PLP-10 light pulser was used to test the hodoscope electronics.

145 4. Light Collection Studies

146 4.1. Beam Test

During the design phase of the hodoscope, we studied light collection from scin-147 tillator tiles using a detector prototype that consisted of fourteen 2 cm x 1 cm x 1 mm 148 and five 2 cm x 1 cm x 2 mm tiles. Tiles were made from the same scintillator type 149 EJ-212 but had smaller height of 2 cm. Light collection was performed using a sin-150 gle BCF-92 fiber, which was glued to one side of the tile using DYMAX UV optical 15 adhesive. Light was detected by Hamamatsu SiPMs S10362-11-025C and S10362-11-152 050C, which have the pixel size of 25 μ m and 50 μ m, respectively. These SiPMs have 153 a smaller sensitive area of 1 mm x 1 mm. Prototypes of the SiPM amplifier and the 154 bias control circuit were used. The trigger was provided by a small plastic scintillator 155 mounted behind the tiles. The trigger scintillator was connected to a Philips XP 1911 156 photomultiplier tube via two optical fibers. The setup used in the beam test is shown in 157 Fig. 8. SiPM pulses were digitized using a flash ADC in a VXS crate. 158

The hodoscope prototype was positioned in the tagger area of the experimental Hall 159 B at Jefferson Lab [3] and was operated concurrently with the CLAS detector. The 160 CLAS used a secondary beam of photons produced in the tagger area by the Jefferson 16 Lab electron beam in a thin radiator via the bremsstrahlung process. A beam electron 162 that radiates a photon is deflected in the dipole tagger magnet and sent to the floor of 163 the experimental hall. The prototype was placed inside a light-tight box and positioned 164 on the floor of the tagger area. The hodoscope was oriented in such a way that it 165 faced normal to the path of electrons, i.e., electrons travel 1 cm through the tile. A 166 schematic view of the tagger area and the location of the prototype is shown in Fig. 9. 167 Dashed curves on this plot denote trajectories of tagged electrons with the fractional 168 energy k/E_0 , where k is the tagged electron energy and E_0 is the electron beam energy. 169 Data was taken with an electron beam energy of 2.54 GeV, corresponding to a tagged 170 electron energy of about 1.4 GeV. 171

The amount of light collected from a tile was measured in units of pixels fired in the SiPM. The average number of pixels was estimated to be

$$N_{\rm pixels} = A_{\rm peak} / A_{\rm single},\tag{1}$$

where A_{peak} is the ADC value corresponding to the maximum of an ADC spectrum 172 measured in the beam test and A_{single} is the ADC value of a single pixel peak. The 173 position of the single pixel peak was calibrated using a LED pulser. Exact counting 174 of the number of photons that induce pixel-breakdowns is complicated due to the op-175 tical cross-talk and after-pulses in the SiPM. The cross-talk is produced by secondary 176 photons emitted in the pixel breakdown that trigger an additional avalanche in a neigh-177 boring pixel. The after-pulsing is believed to be produced by charge carriers that are 178 released from the substrate after some delay from the avalanche resulting in consec-179 utive SiPM pulses. The cross-talk and after-pulse probabilities for the SiPMs used in 180 the hodoscope are expected to be relatively small, less than 10% [8]. For the prototype 18 operational bias voltage (about 1 V over) we expect on average about 0.2 pixels fire 182 per incident photon. SiPM operational voltages for the final detector will be adjusted 183 according to the experimental conditions and detector performance. 184

The beam test results are listed in Table 1. The amount of light collected from the 185 1 mm and 2 mm wide tiles is found to be almost the same and corresponds to about 186 29 pixels fired in the SiPM with the pixel size of 50 µm. This indicates that the light 187 collection mainly depends on the surface of the WLS fiber glued to the tile, which is 188 similar for 1 mm and 2 mm tiles. The fewer number of fired pixels measured with a 189 25 µm SiPM is consistent with about a factor of two smaller photon detection efficiency 190 of these SiPMs compared with 50 µm SiPMs. The smaller pixel fill factor (the ratio of 191 the active area of a pixel to the entire area) accounts for the smaller PDE for the 25 µm 192 S10362-11-025C SiPMs. We estimate that errors on light collection measurements are 193 dominated by alignment of the WLS fiber with respect to the SiPM. The sensitive area 194 of SiPMs used in the prototype is the same as the fiber cross section. 195

Light collection from the final hodoscope detector will be performed using two WLS fibers and SiPMs with a factor of 9 larger effective area. Therefore, we expect about a factor of two more light collected from the tile and a SiPM response corresponding to about 60 fired pixels. This amount of light should be sufficient to provide a high detection efficiency for e^{\pm} pairs in the hodoscope. The hodoscope can be operated at relatively small ADC thresholds because a coincidence of hits in two detector arms

Tile width (mm)	1		2
SiPM pixel size (µm)	25	50	50
Light detected (pixels)	16.1 ± 2.1	28.7 ± 3.4	29.2 ± 3.6

Table 1: The amount of light in units of fired pixels seen by SiPM for various tiles used in the beam test.

202 is required.

203 4.2. Cosmic Setup

During the fabrication of scintillator tiles for the final hodoscope, we checked light 204 collection using a cosmic ray setup presented in Fig. 10. Muon candidates were se-205 lected using two small trigger scintillator counters with an effective area of 6 mm x 206 3 cm and thickness of 2 mm. Counters were positioned perpendicular to each other 207 and formed an overlap region of 6 mm x 6 mm. The distance between scintillators 208 was about 20 cm. A hodoscope tile was placed between these counters, in such a way 209 that cosmic particles go perpendicular to the tile, in the middle. The average muon 210 path length in the scintillator corresponds to the tile width. A lead brick about 10 cm 211 thick was position above the bottom scintillator to filter low-energy particles. Light 212 detection was performed using a Hamamatsu S10931-050P SiPM instrumented with a 213 single-channel amplifier. The amplifier gain was about a factor of ten smaller than that 214 used in the hodoscope. A typical ADC spectrum obtained from the SiPM is shown on 215 the top plot of Fig. 11. The average amount of light collected from 2 mm wide tiles 216 corresponds to about 18 pixels fired in the SiPM. A Sr90 source was used to check 217 quality of some selected fabricated tile assemblies. The SiPM spectrum obtained from 218 a 2 mm wide tile using the radioactive source is shown on the bottom plot of Fig. 11 219 for comparison. 220

221 5. Summary

We have described the design and fabrication details of the pair spectrometer hodoscope, an array of thin scintillator tiles. Light from each tile is detected using a 3

mm x 3 mm Hamamatsu SiPM. A detector prototype was built to perform light collec-224 tion studies using relativistic electrons produced in the experimental Hall B at Jefferson 225 Lab. The amount of light for 1 mm and 2 mm tiles measured with the prototype corre-226 sponds to a SiPM response of about 29 fired pixels. At least a factor of two more light 227 is expected for the final hodoscope for which the light collection is performed using 228 two WLS fibers and tiles with larger size, i.e., larger light collection surface, are used. 229 This amount of light is sufficient to provide a high detection efficiency of leptons by 230 the hodoscope. Two arms of the hodoscope detector were commissioned and installed 23 in the experimental Hall D. 232

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Figure 2: Schematic view of the pair spectrometer hodoscope (top). Light from scintillator tiles is collected using WLS fibers, which are inserted into holes in the mounting plate. Fabricated detector is shown in the bottom picture.



Figure 3: Scintillator tiles with two WLS fibers glued to sides of each tile: 3 cm x 1 cm x 1 mm tile (left) and 3 cm x 1 cm x 2 mm tile (right).



Figure 4: Electronics board with 145 silicon photomultipliers.



Figure 5: A typical flash ADC signal pulse obtained from a scintillator tile.



Figure 6: An example of the ADC spectrum obtained from SiPMs during electronics tests using a laser pulser. The first peak on the spectrum around ADC count zero corresponds to the ADC pedestal.



Figure 7: Distribution of the single-pixel peak measured by flash ADC for 290 hodoscope channels.



Figure 8: Top view of the hodoscope prototype with the trigger scintillator used in the beam test in experimental Hall B.



Figure 9: Hall B tagger system. Location of the hodoscope prototype under the Hall B tagger. Dashed curves denote trajectories of tagged electrons.



Figure 10: A cosmic ray setup.



Figure 11: ADC spectra for a 2 mm wide tile measured with cosmic muons (a) and a ${}^{90}Sr$ radioactive source (b).